## The Fuzzy Onion

# Samuel Kováčik ${ }^{1,2}$ with Juraj Tekel ${ }^{1}$ <br> Based on 2309.00576 

${ }^{1}$ Comenius University, Bratislava
${ }^{2}$ Masaryk University, Brno
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## The goal

We want(ed) to make a matrix formulation of three-dimensional quantum space. Sort of like the fuzzy sphere but with one more (radial) dimension.

## The Fuzzy Sphere

A finite-dimensional representation of $s u(2)$ expressed in terms of $N \times N$ Hermitian matrices with a natural cut-off, $I \leq N-1$.

$$
\begin{gathered}
{\left[L_{i}^{(N)}, L_{j}^{(N)}\right]=i \varepsilon_{i j k} L_{k}^{(N)}} \\
{\left[L_{i}^{(N)},\left[L_{i}^{(N)}, Y_{l m}^{(N)}\right]\right]=I(I+1) Y_{l m}^{(N)},\left[L_{3}^{(N)}, Y_{l m}^{(N)}\right]=m Y_{l m}^{(N)}} \\
\Phi^{(N)}=\sum_{l=0}^{N-1} \sum_{m=-l}^{I} c_{l m}^{(N)} Y_{l m}^{(N)}
\end{gathered}
$$

This comes without the mildly annoying ${ }^{(N)}$ and usually with physical scales $x_{i}=\lambda L_{i}, x^{2}=r^{2}$. Note that $N \sim r / \lambda+\ldots$

## The Fuzzy Sphere

$$
\begin{gathered}
\Phi^{(N)}=\sum_{l=0}^{N-1} \sum_{m=-I}^{l} c_{l m}^{(N)} Y_{I m}^{(N)} \\
\Phi^{(\infty)}(\theta, \phi)=\sum_{l=0}^{N-1} \sum_{m=-I}^{I} c_{l m}^{(\infty)} Y_{l m}^{(\infty)}
\end{gathered}
$$

$-0.182241 \quad-0.356949+0.0169752$ i $0.0260558+0.055678 \mathrm{i}-0.0418167-0.358403 \mathrm{i}$ $-0.356949-0.0169752 \mathrm{i} \quad 0.723061 \quad-0.266625-0.323709 \mathrm{i}-0.209613-0.250825 \mathrm{i}$ $\begin{array}{lllll}0.0260558-0.055678 i & -0.266625+0.323709 i & 0.93628 & 0.115833+0.0969497 i\end{array}$ $\begin{array}{lllll}-0.0418167+0.358403 i & -0.209613+0.250825 i & 0.115833-0.0969497 i & 0.30945\end{array}$

## The Fuzzy Sphere

Also usually one is interested in some kind of field theory on the fuzzy sphere

$$
S_{N}\left[\Phi^{(N)}\right]=\frac{4 \pi}{N} \operatorname{tr}_{N}\left(a \Phi^{(N)} \mathcal{K}^{(N)} \Phi^{(N)}+b\left(\Phi^{(N)}\right)^{2}+c\left(\Phi^{(N)}\right)^{4}\right)
$$

where

$$
\mathcal{K}^{(N)} \Phi^{(N)}=\left[L_{i}^{(N)},\left[L_{i}^{(N)}, \Phi^{(N)}\right]\right] .
$$

With this, one can compute mean values of observables:

$$
\langle\mathcal{O}(\Psi)\rangle=\frac{1}{Z} \int d \Psi e^{-S(\Psi)} \mathcal{O}(\Psi), d \Psi=\prod_{N=1}^{M} d \Phi^{(N)}
$$

This is nice!

## Bosonic contruction

A similar starting point

$$
\left[x_{i}, x_{j}\right]=2 \lambda i \varepsilon_{i j k} x_{k}
$$

We can now invoke an auxiliary Fock space and two sets of $c / a$ bosonic operators:

$$
\begin{gathered}
{\left[\mathrm{a}_{\alpha}, \mathrm{a}_{\beta}^{\dagger}\right]=\delta_{\alpha \beta}, \quad\left[\mathrm{a}_{\alpha}, \mathrm{a}_{\beta}\right]=\left[\mathrm{a}_{\alpha}^{\dagger}, \mathrm{a}_{\beta}^{\dagger}\right]=0} \\
\frac{\left(\mathrm{a}_{1}^{\dagger}\right)^{n_{1}}\left(\mathrm{a}_{2}^{\dagger}\right)^{n_{2}}}{\sqrt{n_{1}!n_{2}!}}|0\rangle=\left|n_{1}, n_{2}\right\rangle
\end{gathered}
$$

Now one can take

$$
x_{i}=\lambda a^{\dagger} \sigma_{i} a
$$

to satisfy the commutation relation and

$$
r=\lambda\left(\mathrm{a}^{\dagger} \mathrm{a}+1\right)
$$

## Bosonic contruction

One can define physical content using this construction; for example take $\Psi=\Psi(x)$ and

$$
H_{0} \Psi=\frac{1}{2 \lambda r}\left[\mathrm{a}_{\alpha}^{\dagger},\left[\mathrm{a}_{\alpha}, \Psi\right]\right] .
$$

This was done thoroughly for the Coulomb problem and the spectrum was found (exactly)

$$
E_{\lambda n}^{\prime}=\frac{\hbar}{m_{e} \lambda^{2}}\left(1-\sqrt{1+\frac{m_{e} q \lambda^{2}}{\hbar^{2} n}}\right)
$$

This is also nice but differently.

## Two similar models of quantum space

The fuzzy sphere: Convenient to work with, but it is a two-dimensional space.

Bosonic construction of $\mathbb{R}_{\lambda}^{3}$ : Is three-dimensional (like our space). Working with it is slightly cumbersome.

Here we introduce the fuzzy onion: a three-dimensional quantum space with matrix realisation that mimics the bosonic construction.

The fuzzy onion


The fuzzy onion

$$
\Psi=\left(\begin{array}{lllll}
(.) & & & & \\
& (\cdots .) & & & \\
& & \left(\begin{array}{ll}
\ddots & \ddots
\end{array}\right) & \\
& & \ddots & .
\end{array}\right)
$$

## The kinetic term

The angular part of the kinetic term is easy to define

$$
\mathcal{K}_{L} \Psi=r^{-2}\left(\begin{array}{lllll}
\mathcal{K}^{(1)} \Phi^{(1)} & & & & \\
& \mathcal{K}^{(2)} \Phi^{(2)} & & & \\
& & \mathcal{K}^{(3)} \Phi^{(3)} & & \\
& & & \ddots & \\
& & & & \mathcal{K}^{(M)} \Phi^{(M)}
\end{array}\right)
$$

Considered before by Vitale, Jurić, Wallet, Poulain and others.

## The kinetic term

The radial part of the kinetic term is harder to define. We want to compare objects with different numbers of d.o.f.

$$
\begin{gathered}
\left(\begin{array}{ll}
\cdot & \cdot \\
\cdot & \cdot
\end{array}\right) \leftrightarrow\left(\begin{array}{lll}
\cdot & \cdot & \cdot \\
\cdot & \cdot & \cdot \\
\cdot & \cdot & \cdot
\end{array}\right) \\
\Phi^{(N)}=\sum_{l=0}^{N-1} \sum_{m=-1}^{l} c_{l m}^{(N)} Y_{l m}^{(N)}
\end{gathered}
$$

All expansion coefficients can be matched but the ones corresponding to the highest momentum. Solution: ignore them.

## The kinetic term

$$
\begin{gathered}
\mathcal{D}: \Phi^{(N+1)} \rightarrow \Phi^{(N)}=\sum_{l=0}^{N-1} \sum_{m=-1}^{l} c_{l m}^{(N)} Y_{l m}^{(N)}, c_{l m}^{(N)}=c_{l m}^{(N+1)} \\
\mathcal{U}: \Phi^{(N)} \rightarrow \Phi^{(N+1)}=\sum_{l=0}^{N} \sum_{m=-1}^{l} c_{l m}^{(N+1)} Y_{l m}^{(N+1)},\left\{\begin{array}{l}
c_{l m}^{(N+1)}=c_{l m}^{(N)} \\
c_{N m}^{(N+1)}=0
\end{array}\right. \\
\partial_{r}^{(N)} \Phi^{(N)}=\frac{\mathcal{D} \Phi^{(N+1)}-\mathcal{U} \Phi^{(N-1)}}{2 \lambda} \\
\mathcal{K}_{R} \Psi=\sum_{N, l, m} \frac{(N+1) c_{l m}^{(N+1)}+(N-1) c_{l m}^{(N-1)}-2 N c_{l m}^{(N)} Y_{l m}^{(N)}}{N \lambda^{2}}
\end{gathered}
$$

## The fuzzy onion

We have the field content $\Psi$ and the kinetic term $\mathcal{K}=\mathcal{K}_{L}+\mathcal{K}_{R}$.
Now we can do physics!

## Example I: Scalar field theory

$$
S[\Psi]=4 \pi \lambda^{2} \operatorname{Tr} r\left(\Psi \mathcal{K} \Psi+b \Psi^{2}+c \Psi^{4}\right)
$$



## Example II: Heat transfer

$$
\mathcal{K} \Psi(t)=\alpha \partial_{t} \Psi(t)
$$



## Example III: Coulomb problem

$$
H=-\frac{\hbar^{2}}{2 m_{e}}-\frac{q}{r}
$$

$$
\mathbf{H} \mathcal{C}_{l m}=E \mathcal{C}_{l m}
$$

| $n$ | 1 | 2 | 3 | 4 | 5 | 6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $E_{n}$ | -0.4142 | -0.1180 | -0.0541 | -0.0307 | -0.0179 | -0.0031 |
| $E_{\lambda n}^{l}$ | -0.4142 | -0.1180 | -0.0541 | -0.0307 | -0.0198 | -0.0138 |
| $E_{n}^{C Q M}$ | -0.5 | -0.125 | -0.0556 | -0.0313 | -0.02 | -0.0139 |

$$
E_{\lambda n}^{\prime}=\frac{\hbar}{m_{e} \lambda^{2}}\left(1-\sqrt{1+\frac{m_{e} q \lambda^{2}}{\hbar^{2} n}}\right)
$$

## Example III: Coulomb problem




- For $\lambda \ll a_{0}$ and $\lambda N \gg a_{0}$ we have great agreement with QM.
- For $\lambda \sim a_{0}$ and $\lambda N \gg a_{0}$ strong quantum-space effects.
- For $\lambda N \sim a_{0}$ space is not large enough to capture the physics.


## Outlook

## Scalar field theory

- What we have done so far is a very basic preliminary study.
- Study the phase diagram.
- Correlation functions, coexistence of phases, aligning stripes.


## Outlook

## Classical physics

- Behavior of granular materials.
- Pixelated physics (weather forecast, planetary weather, ...)
- Neutron stars, gravitational collapse, ...


## Outlook

## Expansion of the space

- One can think of the radiation direction as a temporal one, $\Delta_{R} \rightarrow \Delta_{t}$.
- The model would then describe an expanding fuzzy sphere with a growing number of degrees of freedom.
- Solution to the past hypothesis? (probably not)
- Definitely not a good model of expanding universe but a nice toy model to understand some ideas.

Thank you for your attention!

