

# Tropical field theory

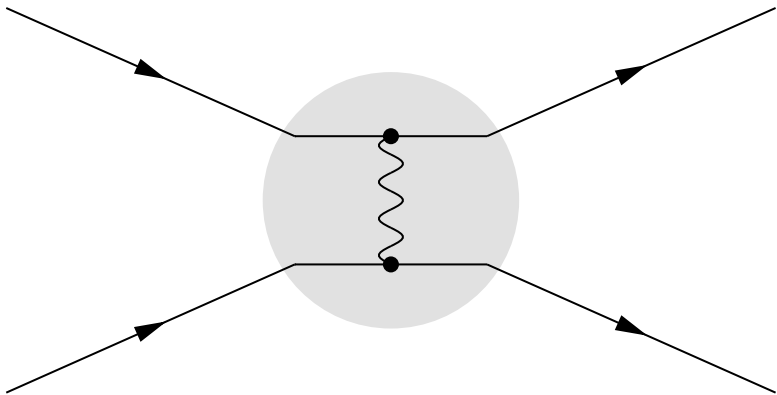
Erik Panzer

Royal Society University Research Fellow (Oxford)

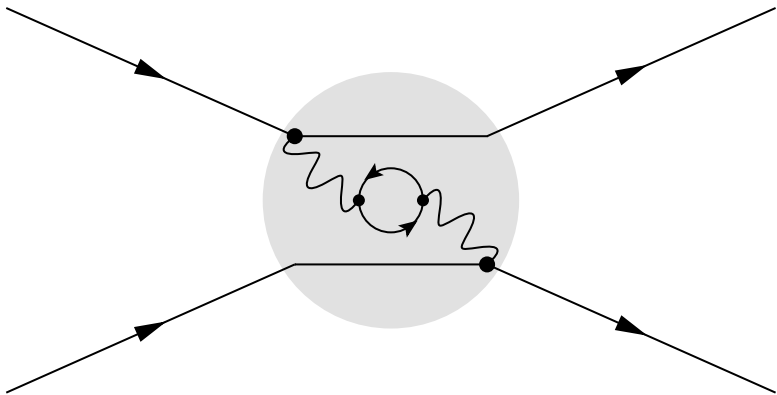
15th November 2021

ESI Vienna

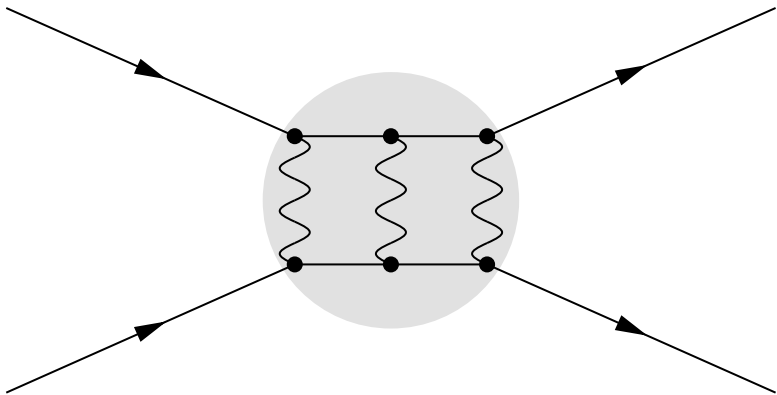
## Perturbative Quantum Field Theory



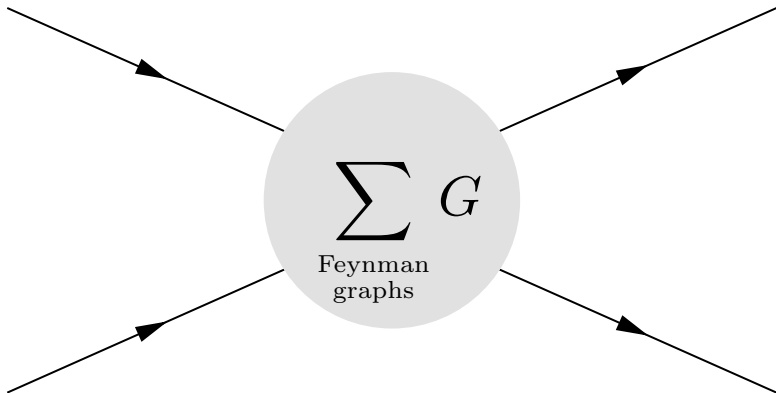
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- Feynman graph  $G \mapsto$  Feynman integral  $\Phi(G, \{m_i^2, \vec{p}_i \cdot \vec{p}_j\})$
- compute more graphs  $\sum_G \Phi(G) \Rightarrow$  higher precision

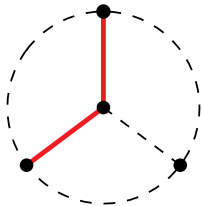
## Problems:

- 1  $\Phi(G)$  extremely complicated  
➔ *polylogarithms, iterated elliptic integrals, modular forms, K3 surfaces, Calabi-Yau manifolds, ...*
- 2  $\sum_G \Phi(G) = \infty$   
➔ *factorial growth  $A \cdot n! \cdot c^n \cdot n^\alpha$ , resummation, Borel transformation, resurgence, ...*

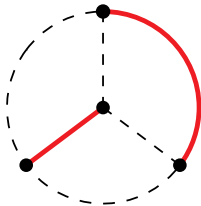
perturbation series are very poorly understood in QCD, QED,  $\phi^4$

## Simplifications:

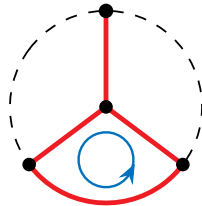
- 1 integrable models  
➔  $\Phi(G)$  *simplify somewhat, full  $\sum_G \Phi(G)$ , **convergent expansion***
- 2 truncated Dyson-Schwinger equations  
➔ *factorial growth possible, **sums very restricted class of diagrams***
- 3 tropical limit  
➔ *all  $\Phi(G)$  **simplify drastically, asymptotics unchanged***



not spanning



not connected



has a loop

## Definition

A **spanning tree**  $T \subset G$  is a spanning, simply connected subgraph.

$$\text{ST} \left( \left( \begin{array}{c} \bullet \\ | \\ \bullet \\ / \quad \backslash \\ \bullet \quad \bullet \end{array} \right) \right) = \left\{ \begin{array}{c} \bullet \\ | \\ \bullet \\ / \quad \backslash \\ \bullet \quad \bullet \end{array} \right\}, \left\{ \begin{array}{c} \bullet \\ | \\ \bullet \\ / \quad \backslash \\ \bullet \quad \bullet \end{array} \right\}, \dots \left\} \right.$$

## Definition

The **graph polynomial**  $\mathcal{U}$  and **Feynman period** of  $G$  are

$$\mathcal{U} = \sum_{T \in \text{ST}(G)} \prod_{e \notin T} x_e \quad \text{and} \quad \mathcal{P}(G) = \left( \prod_{e > 1} \int_0^\infty dx_e \right) \frac{1}{\mathcal{U}^2|_{x_1=1}}$$

$$G = \text{circle with 2 vertices} \Rightarrow \mathcal{U} = x_1 + x_2 \quad \text{and} \quad \mathcal{P}(\text{circle with 2 vertices}) = \int_0^\infty \frac{dx_2}{(1+x_2)^2} = 1$$

- contribute to the  $\beta$ -function

➔ *renormalization constants, running coupling, critical exponents*

- very hard to compute, **even numerically**

## Example

$$\mathcal{P}(\text{circle with 3 vertices}) = \int_{\mathbb{R}_+^5} \frac{dx_2 dx_3 dx_4 dx_5 dx_6}{(x_1 x_2 x_3 + 15 \text{ more terms})^2|_{x_1=1}} = 6\zeta(3) = 6 \sum_{n=1}^{\infty} \frac{1}{n^3}$$

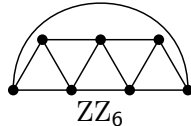
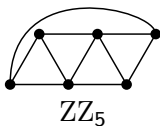
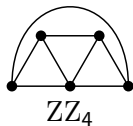
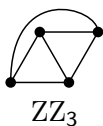


Only one *infinite* family of periods is known in  $\phi^4$ :

Theorem (Brown & Schnetz 2012)

$$\mathcal{P}(\mathbb{Z}\mathbb{Z}_n) = 4 \frac{(2n-2)!}{n!(n-1)!} \left(1 - \frac{1 - (-1)^n}{2^{2n-3}}\right) \zeta(2n-3)$$

conjectured by Broadhurst & Kreimer



- $> 1000$  periods are known in  $\phi^4$  [Broadhurst, Kreimer, Schnetz, Panzer]
- complete only up to 7 loops

## Hepp bound

$$\mathcal{H}(G) = \left( \prod_{e>1} \int_0^\infty dx_e \right) \frac{1}{\mathcal{U}_{\max|x_1=1}^2} \quad \text{where} \quad \mathcal{U}_{\max} = \max_{T \in \text{ST}} \prod_{e \notin T} x_e$$

### Example:

$$\mathcal{H} \left( \text{circle with two vertices} \right) = \int_0^\infty \frac{dx_2}{(\max\{1, x_2\})^2} = \int_0^1 dx_2 + \int_1^\infty \frac{dx_2}{x_2^2} = 2$$

### Properties:

- $\mathcal{H}(G) > \mathcal{P}(G) > \mathcal{H}(G)/|\text{ST}(G)|^2 \Rightarrow$  same asymptotics up to  $\mathcal{O}(c^\ell)$
- $\mathcal{H}(G) \in \mathbb{Q}_{>0}$
- computable for all  $G$
- correlates with  $\mathcal{P}(G)$
- respects symmetries of  $\mathcal{P}(G)$
- generalizes to  $\Phi(G, m_e^2, p_i \cdot p_j)$

$$\omega(\gamma_\ell) = |\gamma| - 2\ell$$

## Theorem

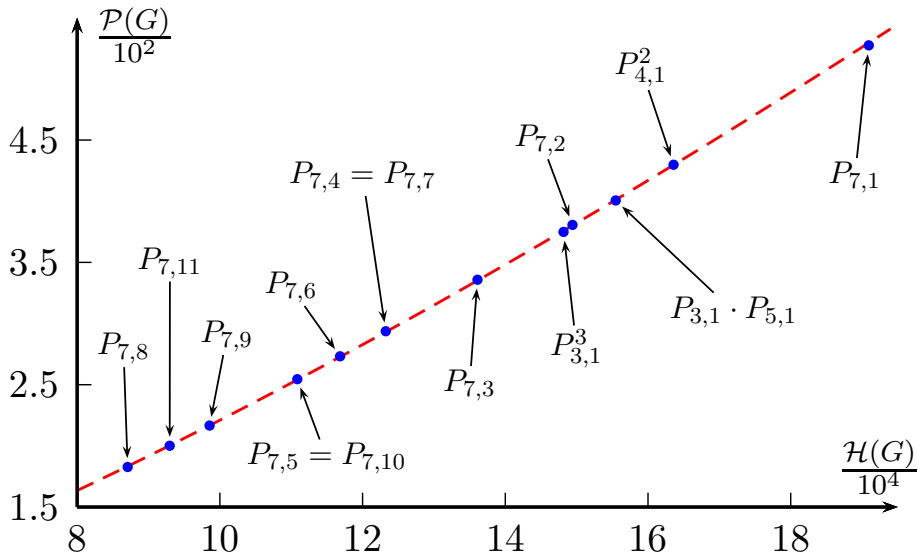
$$\mathcal{H}(G) = \sum_{\substack{\gamma_1 \subset \gamma_2 \subset \dots \subset \gamma_\ell = G \\ \text{each } \gamma_i \text{ is 1PI}}} \frac{|\gamma_1| \cdot |\gamma_2 \setminus \gamma_1| \cdots |G \setminus \gamma_{\ell-1}|}{\omega(\gamma_1) \cdots \omega(\gamma_{\ell-1})}$$

$\gamma_1$	$\subset$	$\gamma_2$	summand	#	$\Sigma$
	$\subset$		$\frac{3 \cdot 2 \cdot 1}{1 \cdot 1} = 6$	12	72
	$\subset$		$\frac{4 \cdot 1 \cdot 1}{2 \cdot 1} = 2$	6	12

$\left. \vphantom{\begin{matrix} \text{row 1} \\ \text{row 2} \end{matrix}} \right\} \Rightarrow \mathcal{H} \left( \text{triangle} \right) = 84$

7 loops in  $\phi^4$ :

$G$	$\mathcal{P}(G \setminus v)$	$\mathcal{H}(G \setminus v)$
$P_{7,1}$	527.7	190952
$P_{4,1} \cdot P_{4,1}$	430.1	163592
$P_{3,1} \cdot P_{5,1}$	400.9	155484
$P_{7,2}$	380.9	149426
$P_{3,1} \cdot P_{3,1} \cdot P_{3,1}$	375.2	148176
$P_{7,3}$	336.1	136114
$\{P_{7,4}, P_{7,7}\}$	294.0	123260
$P_{7,6}$	273.5	116860
$\{P_{7,5}, P_{7,10}\}$	254.8	110864
$P_{7,9}$	216.9	98568
$P_{7,11}$	200.4	92984
$P_{7,8}$	183.0	87088



**Michael Borinsky:**

Tropical Monte Carlo quadrature for Feynman integrals

➔ numeric evaluation at large loop orders

## Symmetries

When is  $\mathcal{P}(G_1) = \mathcal{P}(G_2)$ ?

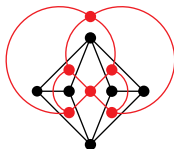
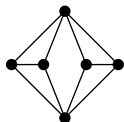
① Product:

$$\mathcal{P} \left( \begin{array}{c} \text{---} \text{---} \text{---} \text{---} \\ \diagup \quad \diagdown \quad \diagup \quad \diagdown \\ \bullet \quad \bullet \\ \vdots \\ \bullet \quad \bullet \\ \diagdown \quad \diagup \quad \diagdown \quad \diagup \\ \text{---} \text{---} \text{---} \end{array} \right) = \mathcal{P} \left( \begin{array}{c} \text{---} \text{---} \text{---} \\ \diagup \quad \diagdown \\ \bullet \\ \vdots \\ \bullet \\ \diagdown \quad \diagup \\ \text{---} \text{---} \end{array} \right) \cdot \mathcal{P} \left( \begin{array}{c} \bullet \\ \vdots \\ \bullet \\ \diagdown \quad \diagup \\ \text{---} \text{---} \end{array} \right)$$

Example

$$\mathcal{P} \left( \begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \quad \bullet \\ \vdots \\ \bullet \quad \bullet \\ \diagdown \quad \diagup \\ \bullet \end{array} \right) = \mathcal{P} \left( \begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \\ \vdots \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \end{array} \right)^2 = (6\zeta(3))^2$$

② Planar duality:  $\mathcal{P}(G) = \mathcal{P}(G^{\text{dual}})$



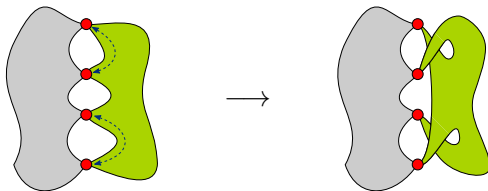


3 Completion:  $\mathcal{P}(G \setminus v) = \mathcal{P}(G \setminus w)$

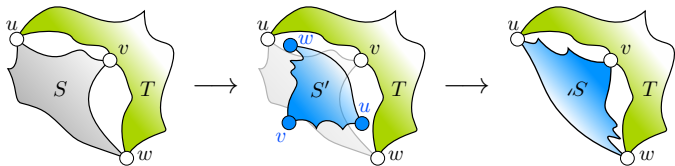
## Example

$$\mathcal{P}\left(\begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \end{array}\right) = \mathcal{P}\left(\begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \\ \text{---} w \text{---} \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \\ \text{---} v \text{---} \\ \bullet \end{array}\right) = \mathcal{P}\left(\begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \\ \text{---} w \text{---} \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \\ \text{---} v \text{---} \\ \bullet \end{array}\right) = \mathcal{P}\left(\begin{array}{c} \bullet \\ \diagup \quad \diagdown \\ \bullet \\ \diagdown \quad \diagup \\ \bullet \end{array}\right)$$

4 Twist:



5 Fourier split:



## Goal:

Construct simpler graph invariants with those symmetries.

- ① point count

[Schnetz]

$$c_2(p) = \frac{1}{p^2} \left| \left\{ \vec{x} \in (\mathbb{Z}/p\mathbb{Z})^N : \mathcal{U}(\vec{x}) = 0 \right\} \right| \pmod{p}$$

- ② extended graph permanent

[Crump]

$P_{7,11}$

$p$	2	3	5	7	11	13	17	19	23
$c_2(p)$	1	0	1	-1	1	-1	1	-1	1
Perm( $p$ )		0	1	1	1	11	5	0	22

- ③  $\mathcal{H}$  ( $\in \mathbb{Q}$ )
- ④ # {minimal 6-cuts} ( $\in \mathbb{Z}$ )
- ⑤  $O(-2)$  symmetry factor ( $\in \mathbb{Z}$ )

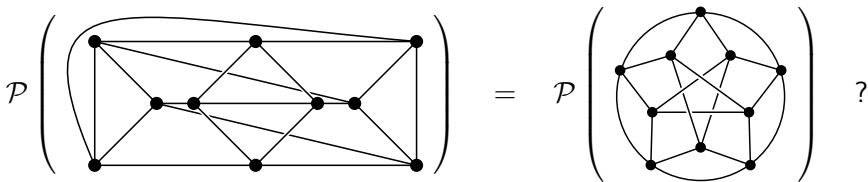
## Conjecture ( $\phi^4$ )

$$\mathcal{H}(G_1) = \mathcal{H}(G_2) \quad \Leftrightarrow \quad \mathcal{P}(G_1) = \mathcal{P}(G_2)$$

### Example:

- $\mathcal{H}(P_{8,30}) = \frac{1724488}{3} = \mathcal{H}(P_{8,36})$
- $\mathcal{P}(P_{8,30}) \approx 505.5 \approx \mathcal{P}(P_{8,36})$

(exact period unknown)



➔ there are further symmetries of Feynman integrals