

Periods, arithmetic & the Hodge structure for Calabi-Yau manifolds II

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The unreasonable effectiveness of toric geometry:
bridging mathematics, computation and string theory

ESI, Universität Wien, 24 June 2026

geometry of families
of Calabi-Yau varieties
and that of their
moduli spaces

arithmetic of families
of Calabi-Yau varieties
and modularity
properties



PERIODS

physics

- black hole solutions in string theory
- scattering amplitudes
- etc, etc, --

Collaborators

- P. Candelas & F Rodriguez-Villegas 2000, 2004
- P. Candelas, D van Straten & M Elmi 2019
- P. Candelas, D van Straten 2021
- P. Candelas, J McGovern & P Kuusela 2021-2024
- P. Candelas & P Kuusela 2024
- P. Candelas, D van Straten, N. Gedda in progress 2026
- P. Candelas & E Svanberg in progress 2026 -

① Calabi-Yau varieties (very brief)

② Calabi-Yau varieties OVER FINITE FIELDS • recap
• very brief

periods
encode
arithmetic
information

{ Exercise: counting points mod p
Counting points exactly

formula in terms of
periods (very briefly!)

③ The attractor mechanism &
arithmetic of attractor varieties

{ illustrate arithmetic
strategy to study the
Hodge structure at
special points

④ Conclusions

①

CALABI-YAU VARIETIES

Mathematical objects of interest:

algebraic varieties with certain special properties

set of solutions of

$$|P(\underline{c}, \underline{x}) = 0, \underline{x} \in \mathbb{A}|$$

polynomials with complex coefficients $\underline{c}$

Calabi-Yau

- Kähler
- $c_1 = 0$
(admit a Ricci-flat metric)

This talk: concerned with $d=3$

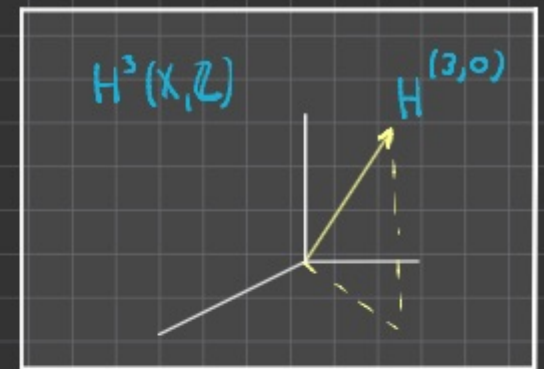
► It is a theorem that $\exists!$ (up to a constant) a nowhere vanishing $(3,0)$ -form Ω which is holomorphic ($d\Omega=0$)

so $h^{(3,0)} = \dim H^{(3,0)} = \dim H^{(0,3)} = 1$

• $H^3 = H^{(3,0)} \oplus H^{(2,1)} \oplus H^{(1,2)} \oplus H^{(0,3)}$

Hodge structure

$\dim H^3 = b_3 = 1 + h^{(2,1)} + h^{(1,2)} + 1 = 2(1 + h^{(2,1)})$



• Vary the complex structure

→ Ω moves inside $H^3(X, \mathbb{C})$ (variations of the HS)
 ($h^{1,1} = \dim \text{moduli space}$)

Examples: very many!

- $\mathbb{P}^4[5]$ eg $\sum_{i=1}^5 x_i^5 - 5\psi x_1 x_2 x_3 x_4 x_5 = 0 \quad (x_1, \dots, x_5) \in \mathbb{P}^4$

$\uparrow$
 $c_1 = 0$

- Verrill 1996, Hulek & Verrill 2005

(quotient of the)
Mirror of a CICY

$$\begin{matrix} \mathbb{P}^1 \\ \mathbb{P}^1 \\ \mathbb{P}^1 \\ \mathbb{P}^1 \\ \mathbb{P}^1 \end{matrix} \left[\begin{array}{cc} 1 & 1 \\ 1 & 1 \\ 1 & 1 \\ 1 & 1 \\ 1 & 1 \end{array} \right]$$

$$h'' = 5$$



quotient

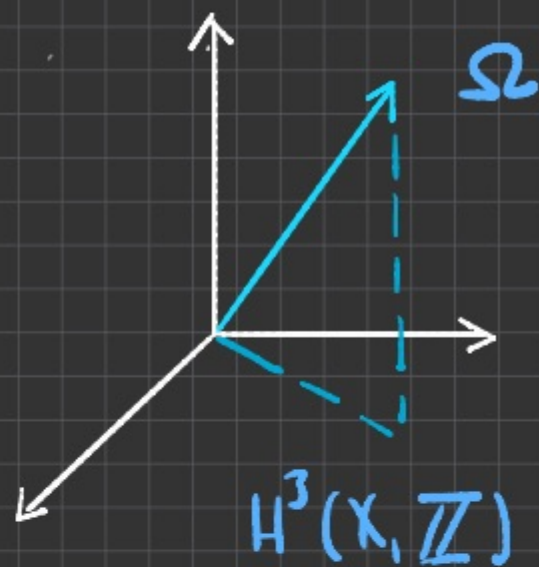
$$h'' = 1$$

why this example? exhibits interesting arithmetic properties
which have an interpretation in Btt solutions
of string theory (also in amplitudes)

Periods

↳ there is a canonical way to give coordinates on the space of complex structures

- study the variations of the complex structure by studying how Ω varies in H^3
- the coordinates of this line are the **periods** (which then vary as we vary the complex structure)



That is,

$$\Omega(\varphi) = z^a(\varphi) \alpha_a - \bar{F}_a(\varphi) \beta^a$$

where

$\{\alpha^a, \beta_b\} \rightarrow$ symplectic basis of $H^3(X, \mathbb{Z})$ $a, b = 0, 1, \dots, h^{2,1}$

$$\int_X \alpha_a \wedge \beta^b = - \int_X \beta^b \wedge \alpha_a = \delta_a^b, \quad \int_X \alpha_a \wedge \alpha_b = 0, \quad \int_X \beta_a \wedge \beta_b = 0$$

let $\{A^a, B_b\}$ the (dual) symplectic basis of $H_3(X, \mathbb{Z})$

$$\text{Then: } z^a(\varphi) = \int_{A^a} \Omega(\varphi)$$

$$\bar{F}_a(\varphi) = \int_{B_a} \Omega(\varphi)$$

$\{z^a, \bar{F}_a\}$ integral basis

$$\Omega \rightarrow \lambda \Omega$$

Bryant & Griffiths: the complex structure is completely determined by the periods $\{z^a\}$ (so $\bar{F}_a = \bar{F}_a(z)$). The set $\{z^a(\varphi)\}$ are projective coordinates on the moduli space and we have that $\dim_{\mathbb{C}} \mathcal{M}_{\text{cs}} = h^{2,1}$

Periods determine the geometry of the moduli space

special geometry

↙ physics!

- Strominger 91
- Candelas & XD 91

periods also have arithmetic content!!

Periods are calculable: they satisfy a differential eq of degree b_3 (Picard-Fuchs equation)

$$\oint \tilde{\omega} = 0, \quad \oint \bar{\omega} = 0$$

Focus today
1-parameter examples (Fuchsian)
($b_3 = 2(1+h_{11}) = 4$)

Intuitively: $\Omega, \Omega', \Omega'', \dots$ are all closed 3-forms but at most b_3 of these are linearly independent
∴ there must be a linear relation between the first b_3+1
ie $\exists$ an operator d st $d\Omega = \text{exact 3-form}$

► There is a prescription to obtain the PF eq
Dwork and Griffiths

Gelfand, Kapranov and Zelevinski $\rightarrow$ GKZ systems

► One parameter families of CY : $b_3 = 4$

$$\deg d = 4$$

Solutions around MUM point $\varphi = 0$

$$\hookrightarrow (T-1)^4 = 0, (T-1)^3 \neq 0$$

$$\mathcal{D}_0(\varphi) = \sum_{m=0}^{\infty} a_m \varphi^m = f_0(\varphi) \leftarrow \text{fundamental period} \quad (f_0 \rightarrow 1 \text{ as } \varphi \rightarrow 0)$$

$$\left\{ \begin{aligned} \mathcal{D}_1(\varphi) &= f_0(\varphi) \log \varphi + f_1(\varphi) \\ \mathcal{D}_2(\varphi) &= f_0(\varphi) \log^2 \varphi + 2f_1(\varphi) \log(\varphi) + f_2(\varphi) \\ \mathcal{D}_3(\varphi) &= f_0(\varphi) \log^3 \varphi + 3f_1(\varphi) \log^2(\varphi) + 3f_2(\varphi) \log \varphi + f_3(\varphi) \end{aligned} \right.$$

f_i are regular maps

(Frobenius basis)

Example: (mirror) quintic

$$P(x, \psi) = \sum_{i=1}^5 x_i^5 - 5\psi x_1 x_2 x_3 x_4 x_5.$$

one parameter
family

$$\mathcal{L} = \theta^4 - 5\psi \prod_{i=1}^4 (\theta + i/r) \quad \theta = \psi \frac{d}{d\psi} \quad \psi = (5\psi)^{-5}$$

3 regular singularities $\left\{ \begin{array}{ll} \psi = 0 & \text{MUM} \\ \psi = 1 & \text{Conifold} \\ \psi = \infty & \text{Fermat} \end{array} \right.$

hypergeometric

Verrill : X_φ is smooth for $\varphi \neq 1, 1/9, 1/25, 0, \infty$ 1-parameter
case

5 singular cases (not hypergeometric !)

$$\varphi = 1, 1/9, 1/25$$

$$\varphi = 0 \quad \text{MUM}$$

$$\varphi = \infty$$

conifold type

2

CALABI-YAU VARIETIES OVER FINITE FIELDS

2.1

Zeta functions

Let X_q be a family of algebraic varieties
 st X_q is a CY hypersurface with
 defining polynomial $P(\underline{x}, q)$

Let $q \in \mathbb{F}_q$, $q = p^k$

The fundamental quantities of interest are

$N_k(q) =$ number of solutions of $P(\underline{x}, q) = 0$ over $\mathbb{F}_{p^k}$

Generating function $\rightarrow$ zeta function

$$Z_X(T, p; \varphi) = \exp \left\{ \sum_{k=1}^{\infty} \frac{1}{k} N_k(\varphi) T^k \right\}$$

$\left\{ \begin{array}{l} \text{depends on both } p \text{ \& } \varphi \end{array} \right.$

$\rightarrow$ properties $\rightarrow$ Weil conjectures

(proven by Dwork, Deligne, Grothendieck)

$$\deg R(T) = b_3$$

$$Z_X(T) = \frac{\cancel{P_1(T)} \quad R(T) \quad \cancel{R_1(T)}}{(1-T) \underbrace{(1-pT)^{h''} (1-p^2T)^{h''}} (1-p^3T)}$$

smooth
CY

(Picard group generated by divisors
defined over the ground field $\overline{\mathbb{F}_p}$)

$$b_2 = b_4 = h''$$

1-parameter families: $h^{2,1} = 1$, $b_3 = 1+1+1+1 = 4$ so $\deg R = 4$

$$R(T) = 1 + a_p T + b_p p T^2 + a_p p^3 T^3 + p^6 T^3$$

$$= \det(1 - T u(\varphi)) \quad (\text{eg, } a_1 = -4\eta, \dots)$$

$$u = \text{Frob}_3^{-1}: H^3(X) \longrightarrow H^3(X)$$

map induced by the Frob automorphism on H^3

4x4 matrix constructed from the periods of Ω $\left\{ \begin{array}{l} u(\varphi) = E^{-1}(\varphi^p) u(0) E(\varphi) ; E_{ij}(\varphi) = \vartheta'(\frac{1}{j}, \vartheta_i) \\ u(0) = \begin{bmatrix} 1 & p & p^2 & p^3 \\ p^3 & p^2 & p^1 & p^0 \end{bmatrix} \end{array} \right. \quad \gamma = \frac{x}{y} S_p(3) \quad \vartheta = \varphi \frac{d}{d\varphi}$

singular CY: $\deg R(T) < 4$

Note: $N_p = p^3 + h^1 p^2 + h^2 p + 1 + a_p$

$$N_{p^2} = p^6 + h^1 p^4 + h^2 p^2 + 1 + 2p b_p - a_p^2$$

Want to study how $R(T)$ varies with φ

Need to compute a_p & b_p

→ counting points exactly

* $\chi(T)$ can be "quickly" computed for $\psi \in \mathbb{F}_p$
($\psi = 0, 1, \dots, p-1$) for many primes p

P Candelas, XD & D van Straten (2021) 1 parameter families

P Candelas, XD, P Kuuskela (2024) } multi-parameter families

P Kuuskela, M Lathwood, M Moser Rojas, M Stepniczka, (2026)

→ based on the deformation method
developed first by Dwork & Landau

* formulas for point-counts ← next

Many questions arise: certainly ready for substantial experimentation.

► One can construct L-functions

1 parameter families

$$L(s, \psi) = \prod_p \underbrace{R(p^{-s}, \psi)^{-1}}_{\text{spinor zeta function } w=3}$$

What are the properties of this L-function?

- functional equation
- formula for the conductor

Langlands $\longleftrightarrow$ modularity $\longleftrightarrow$ paramodular forms of $\Gamma(N) \leq \mathrm{Sp}(4)$ for 1-parameter cases

[G. Toharia, A. Pacetti, J. Voight, V. Golyshev

P. Candelas, XD, N. Gegeia, D. van Straten, in progress]

2.2

AN EXERCISE: COUNTING POINTS MOD P

► Fermat's
little theorem

$$a^p \equiv a$$

$$\Rightarrow a(a^{p-1} - 1) \equiv 0$$

$$\Rightarrow a^{p-1} \equiv \begin{cases} 1 & a \neq 0 \\ 0 & a = 0 \end{cases}$$

Then: $1 - P(\underline{x}, \varphi)^{p-1} \equiv \begin{cases} 1, & \forall \underline{x} \text{ st } P(\underline{x}, \varphi) = 0 \\ 0, & \forall \underline{x} \text{ st } P(\underline{x}, \varphi) \neq 0 \end{cases}$

...

$$v_p(\varphi) \equiv \sum_{\underline{x} \in \mathbb{F}_p^n} (1 - P(\underline{x}, \varphi)^{p-1})$$

↖ n-variables

Chevalley, Waring 1935

⋮

P. Candelas, XD, F. Rodríguez-Villegas 2000, 2004

⋮

E. Parnitz, F. Brown

(mirror) quintic:

$$\nu_p(\varphi) \equiv \sum_{\underline{x} \in \mathbb{F}_p^5} \left\{ 1 - \left(\sum_{i=1}^5 x_i^5 - 54x_1x_2x_3x_4x_5 \right)^{p-1} \right\}$$

expand (multinomial...)

$$\equiv \sum_{m=0}^{[p/5]} a_m \varphi^m \equiv \text{[p/5]} \varrho_0(\varphi)$$

truncated
fundamental
period.

$$\sum_{a \in \mathbb{F}_p} a^n \equiv \begin{cases} 0, & p-1 \nmid n \\ -1, & p-1 \mid n \end{cases}$$

- result extends to all hypersurfaces in toric varieties and $h^{2,1} \geq 1$

- Want to count points exactly!

- Igusa 1958 class of elliptic curves

⋮

- mirror quintic (P Candelas, XD, F Rodriguez Villegas)

- many other examples eg S. Kadir 2-parameter CY ~ 2004
F Rodriguez V + D Roberts HGM in 2024

K3 surfaces • C Doran, T Kelly, A Salerno, S Sperber, J Voight, L Whitcher 2018, 2020

• Z Davis, J Dukes, T Gomes-Libeiro, E Orvis, A Salerno,
L Sturman & L Whitcher 2025

• F Beukens & M Vlasenko, M Kerr 2020-2025

- In progress with P. Candelas & E Swanberg

2.3

COMPUTING THE NUMBER OF POINTS EXACTLY

$\nu_p(\varphi)$ is a definite number: we seek to compute it exactly

We can expand

$$\nu_p(\varphi) = \nu_p^{(0)}(\varphi) + p \nu_p^{(1)}(\varphi) + p^2 \nu_p^{(2)}(\varphi) + p^3 \nu_p^{(3)}(\varphi) + \dots$$

with $0 \leq \nu_p^{(i)} \leq p-1$

and evaluate $\bmod p^2, \bmod p^3, \dots$

This leads to p -adic analysis!

Counting pants exactly:

In progress with P Candelas & E **Swanberg**: schematically

primitive part of N_P

$$= \sum_{\underline{e}, \underline{1} = \underline{0}}^{\infty} \frac{h_{\underline{e}, \underline{1}}}{\underline{e}!} \left(\frac{P}{1-P} \underline{\Theta} \right)^{\underline{e} + \underline{1}} [1-z] f^{\underline{e}}(\varphi)$$

truncated
f-series

$$\underline{e}, \underline{1} = (e_1, \dots, e_n)$$

$$\underline{\Theta} = \Theta_1 \Theta_2 \dots \Theta_n$$

$$\Theta_i = \varphi_i \frac{\partial}{\partial \varphi_i}$$

mirror

$$\text{quantic} = [1-z] f_0(\varphi) + \left(\frac{P}{1-P} \right) [1-z] f_1(\varphi) + \frac{1}{2} \left(\frac{P}{1-P} \right)^2 [1-z] f_2(\varphi) + \dots$$

[also in F Beukens & M Vlasenko, M Kerr]

3

THE ATTRACTOR MECHANISM

(Furawa, Kallosh 85, Greg Moore 98 ...)

Physics : supersymmetric black hole solutions
of type IIB superstrings

→ 10 dimensional

include { Einstein's equations for gravity
+
Maxwell equations for electromagnetism

10 dim space-time: a CY, $X_{\varphi(r)}$, at each point of 4 dim space-time
 4 dim space-time = spherically symmetric asymptotically flat charged BH



CY with
complex
structure φ
varying
with r

horizon ($r=0$)
(coordinate singularity)
 $u(r) \rightarrow -\infty$

last surface visible from infinity
(surface separating interior
& exterior of the BH)

r

flat 4-dim space-time
 $r \rightarrow \infty$ ($u(r) \rightarrow 0$)

$$ds^2 = e^{2u(r)} dt^2 + e^{-2u(r)} d\vec{x}^2, \quad \vec{x} = (x, y, z)$$

metric dependent only on the radial
coordinate r

Type IIB supergravity contains gravity but also extra
 $h(1)$ gauge fields (h of them)

So: the BH has electric & magnetic charges

$$Q = \begin{pmatrix} q_a \\ p^b \end{pmatrix} \quad a, b = 0, 1, \dots, h^{2,1}$$

↪ these must all be integers

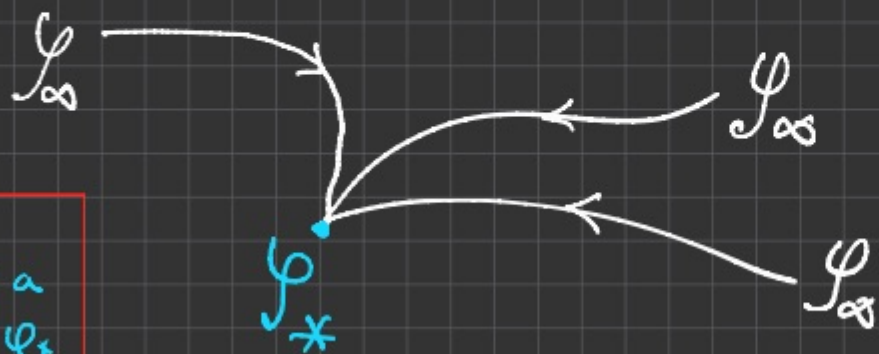
Let $\Gamma = p^a \alpha_a - q_a \beta^a \in H^3(X, \mathbb{Z})$
↪ charge vector

(α_a, β^b) symplectic basis of $H^3(X, \mathbb{Z})$

Black hole solutions which preserve **supersymmetry** need to satisfy 1st order differential eqs for $\psi(r)$ & $\varphi(r)$. These are the attractor equations which represent a non-linear dynamical system on the $\mathbb{G}$ -structure moduli space with flow parameter $\rho = 1/r$.

Given a choice of $\Gamma \in H^3(X, \mathbb{Z})$ one can **prove** using the attractor eq that

- ▶ the $\mathbb{G}$ -structure parameters flow to a value $\varphi_* = \varphi(r=0)$ where $|Z_\Gamma| = |e^{K/2} \int_\Gamma \Omega|$ reaches a minimum, and it is independent of the starting value $\varphi_\infty = \varphi(r=\infty)$



φ evolves smoothly to a fixed point φ_* at $r=0$

95: Ferrara+Kallosh

...

98: G. Moore

conjectures on the arithmetic nature of attractor varieties

$$X_{\varphi_*} = X_*$$

► Moreover,

the $\mathbb{Q}$ -structure at an attractor point $\varphi = \varphi_*$ is st

$$\Gamma = p^a \alpha_a - q_a \beta^a \in H^{(3,0)} \oplus H^{(0,3)}$$

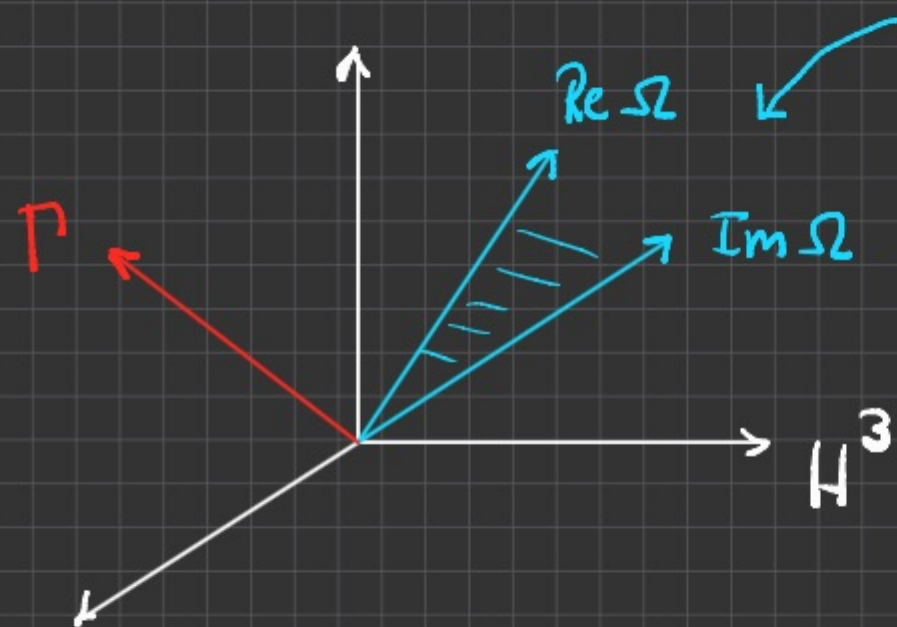
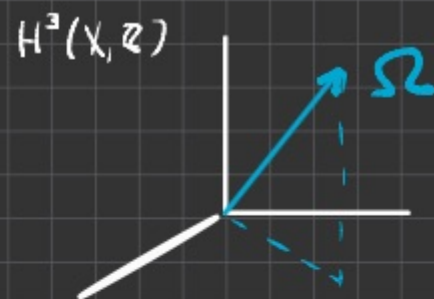
ie $\Gamma^{(1,2)} = \Gamma^{(2,1)} = 0$

proof: exercise in special geometry

rank 1 attractors

Recall: Ω defines a line in $H^3(X, \mathbb{Z})$

Ω moves inside $H^3(X, \mathbb{Z})$ as we change the CS



Consider

$V_{\mathbb{R}}(\varphi)$ = plane spanned
over $\mathbb{R}$ by $\text{Re } \Omega$ & $\text{Im } \Omega$

$V_{\mathbb{R}}(\varphi)$ moves with φ

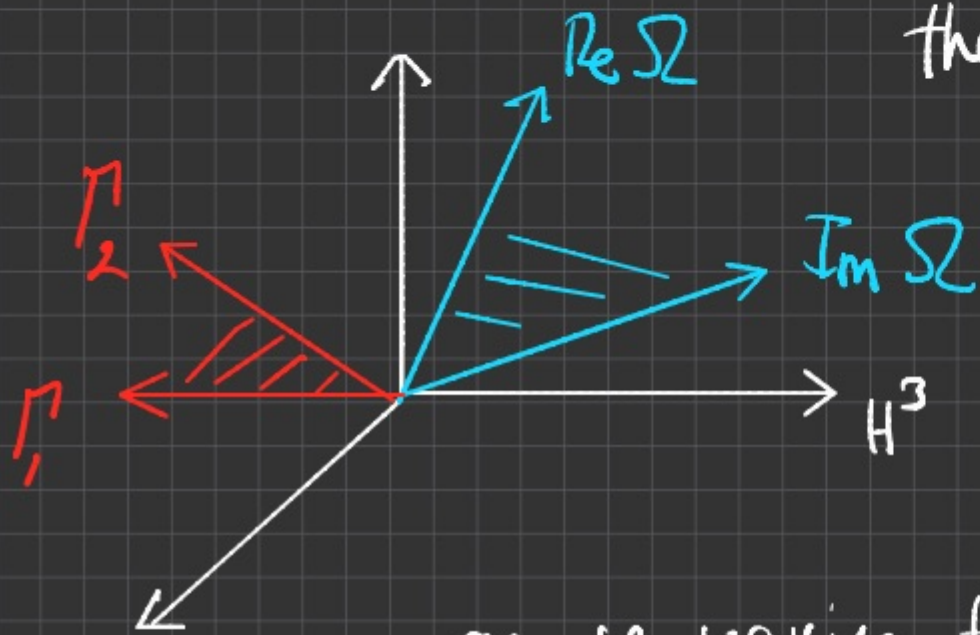
OTOH: Inside $H^3(X, \mathbb{R})$ we have a lattice of vectors
 $\Gamma \subset H^3(X, \mathbb{Z})$ which are fixed

rank 1 A.P: φ st $V_{\mathbb{R}}(\varphi)$ contains the line Γ

rank 2 attractors

at an attractor point of rank 2 there are two vectors in $H^3(X, \mathbb{Z})$

P_1, P_2 st $P_{1,2} \in H^{(3,0)} \oplus H^{(0,3)}$



as u varies the plane $V_{\mathbb{R}}(u)$ moves and at a rank 2 attractor point $u = u_*$ the plane $V_{\mathbb{R}}(u)$ coincides with the plane generated by P_1 & P_2 .

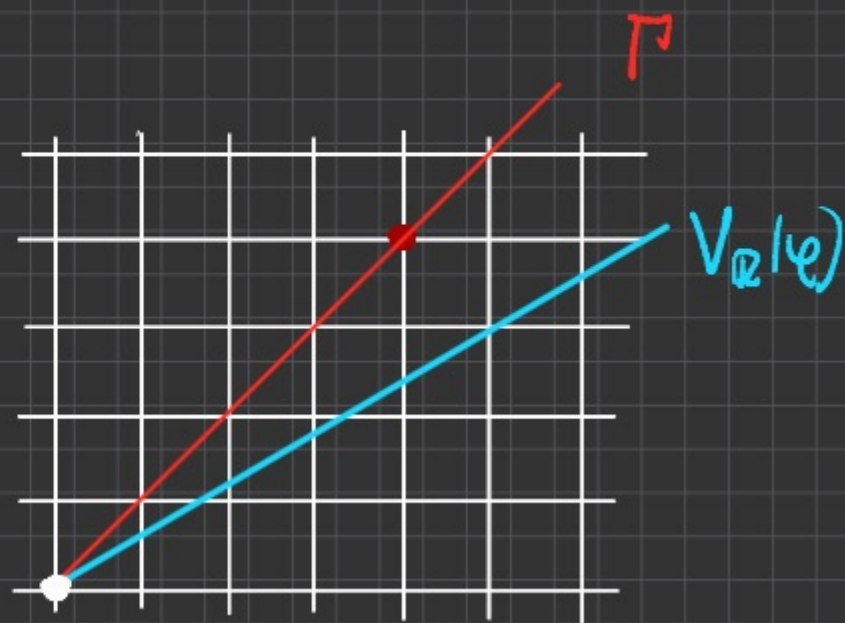
RARE

very difficult to find a CY which has rank 2 attractor points

Why?

A line which passes through the origin in general will not pass through another lattice point unless the slope is rational.

Not too hard to find φ st $V_{\mathbb{Q}}(\varphi)$ coincides with T



For rank 2 attractors we have a **plane** and it is then much harder to find φ st **it** coincides with $V_{\mathbb{Q}}(\varphi)$

Geometrically: at $\varphi = \varphi^*$

$\Gamma_1, \Gamma_2 \in \Lambda := V_{\mathbb{R}}(\varphi) \cap H^3(X, \mathbb{Z})$ is a lattice plane

and $\Lambda \otimes \mathbb{C} = H^{(3,0)} \oplus H^{(0,3)}$

Also: $\Lambda^{\perp} \subset H^3(X, \mathbb{Z})$ lattice orthogonal to Λ
(under the natural symplectic product on 3-forms)

satisfies $\Lambda^{\perp} \otimes \mathbb{C} = H^{(2,1)} \oplus H^{(1,2)}$

We obtain an isomorphism

$$H^3(X, \mathbb{Q}) = \Lambda_{\mathbb{Q}} \oplus \Lambda_{\mathbb{Q}}^{\perp}$$

Λ has finite index
so extends to $\mathbb{Q}$

splitting of the Hodge structure $H^3(X, \mathbb{Q})$!

Hodge conjecture $\Rightarrow$ splitting has a geometrical origin

VERY HARD

$\rightsquigarrow$ open question

Problem: how do we find varieties with
rank 2 attractor points?

However: the splitting becomes apparent in the
arithmetic structure of X

$\therefore$ arithmetic strategy !

P. Candelas, M. Elmi
& D van Straten, 2021



P Candelas, XD, D van Straten April 2021

K Bönish + A Klemm + Scheidegger + Hagler 2022

K Bönish + M Elmi + A.K. Kashani-Poor + A Klemm 2022

Candelas, XD, J McGovern, P Kuusela 2021 & Feb 2023

P Candelas, XD, P Kuusela 2024

④

Arithmetic of **attractor** varieties

At $\varphi = \varphi_x$

Recall $H^3(X_x, \mathbb{Z}) = V \oplus V^\perp$
 $(1,0)+(0,3) \quad (2,1)+(1,2)$

$$R(T) = \det(1 - T \mathcal{U}) \quad \mathcal{U} \rightsquigarrow \begin{pmatrix} 1 & \\ & -1 \end{pmatrix}$$

$$= (1 - \alpha_p T + p^3 T^2) (1 - \beta_p T + p^3 T^2)$$

$H^{2,1} \oplus H^{1,2} \quad H^{3,0} \oplus H^{0,3}$

factor
over $\mathbb{Z}$
 $\forall p$

Moreover: expect α_p, β_p to be coeffs of modular forms
 (Tate & Serre conjectures)

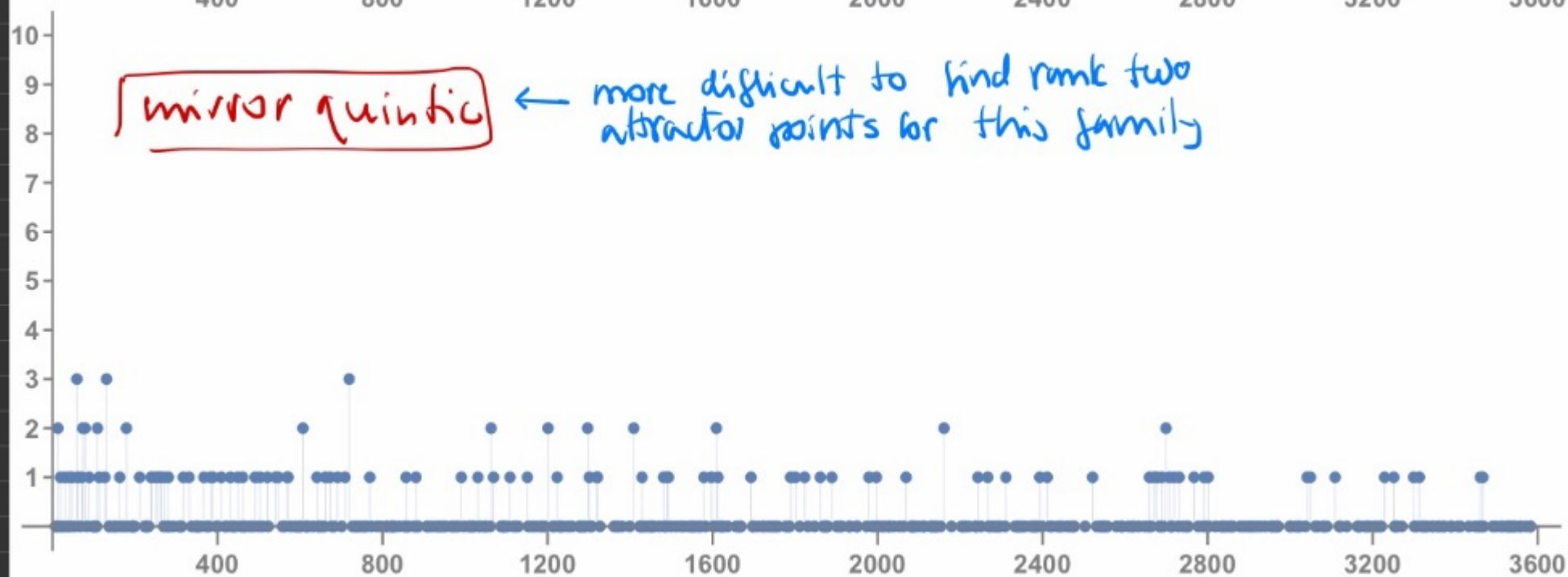
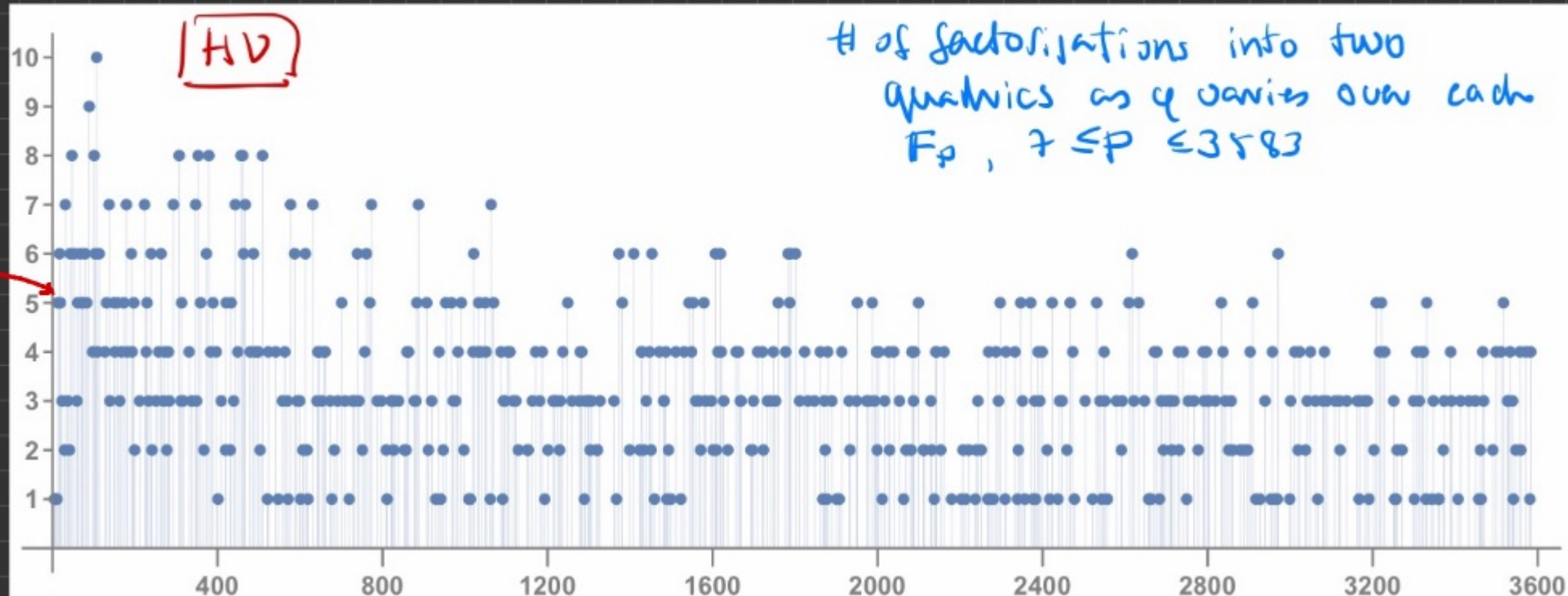
Arithmetic strategy:

make tables of $R(T, \varphi)$ for many p & φ
and look for **persistent factorisations** of R
into two quadrics

↑
factorisations occurring when
 φ_* is a root of a polynomial
with integer coeffs

[P. Candelas, XD, **A Thorne**, DuStraten 2018

P. Candelas, XD, DuStraten 04/2021]



We find that for the HV manifold there is always a factorisation when (P Candelas, XD, M Elmi & D van Straten)

$$\bullet \quad 7\varphi + 1 = 0 \quad : \quad \varphi = -1/7 \quad (\forall p \neq 7)$$

and

$$\bullet \quad \varphi^2 - 66\varphi + 1 = 0 \quad : \quad \varphi_{\pm} = 33 \pm 8\sqrt{17}$$

(exists in $\mathbb{F}_p$ when 17 is a square modulo p)

For $p = 19$ (say)

$$\varphi = -\frac{1}{7} \equiv 8$$

$$\varphi_{\pm} \equiv 4, 5$$

($17 \equiv 6^2$)

Remark: X_{φ_*} is smooth

$p = 19$			
φ	smooth/sing.	singularity	$R(T)$
1	singular	1	$(1 - pT)(1 - 20T + p^3T^2)$
2	smooth		$1 + 4pT + 2pT^2 + 4p^4T^3 + p^6T^4$
3	smooth		$1 - 8T + 242pT^2 - 8p^3T^3 + p^6T^4$
4	smooth		$(1 + 4pT + p^3T^2)(1 - 60T + p^3T^2)$
5	smooth		$(1 + 4pT + p^3T^2)(1 - 60T + p^3T^2)$
6	smooth		$1 + 8T - 318pT^2 + 8p^3T^3 + p^6T^4$
7	smooth		$1 - 44T - 238pT^2 - 44p^3T^3 + p^6T^4$
8	smooth		$(1 - 2pT + p^3T^2)(1 - 80T + p^3T^2)$
9	smooth		$(1 + 4pT + p^3T^2)(1 - 160T + p^3T^2)$
10	smooth		$1 + 12T + 562pT^2 + 12p^3T^3 + p^6T^4$
11	smooth		$(1 + 4pT + p^3T^2)(1 - 140T + p^3T^2)$
12	smooth		$1 + 12T + 82pT^2 + 12p^3T^3 + p^6T^4$
13	smooth		$1 + 178T + 1082pT^2 + 178p^3T^3 + p^6T^4$
14	smooth		$1 + 12T - 158pT^2 + 12p^3T^3 + p^6T^4$
15	smooth		$1 + 42T - 2p^2T^2 + 42p^3T^3 + p^6T^4$
16	singular	$\frac{1}{25}$	$(1 - pT)(1 + 76T + p^3T^2)$
17	singular	$\frac{1}{9}$	$(1 - pT)(1 - 20T + p^3T^2)$
18	smooth		$1 - 54T + 322pT^2 - 54p^3T^3 + p^6T^4$

conifold

φ_+ $\rightarrow$
 φ_- $\rightarrow$
 $\varphi = -1/7$ $\rightarrow$

conifold

conifold

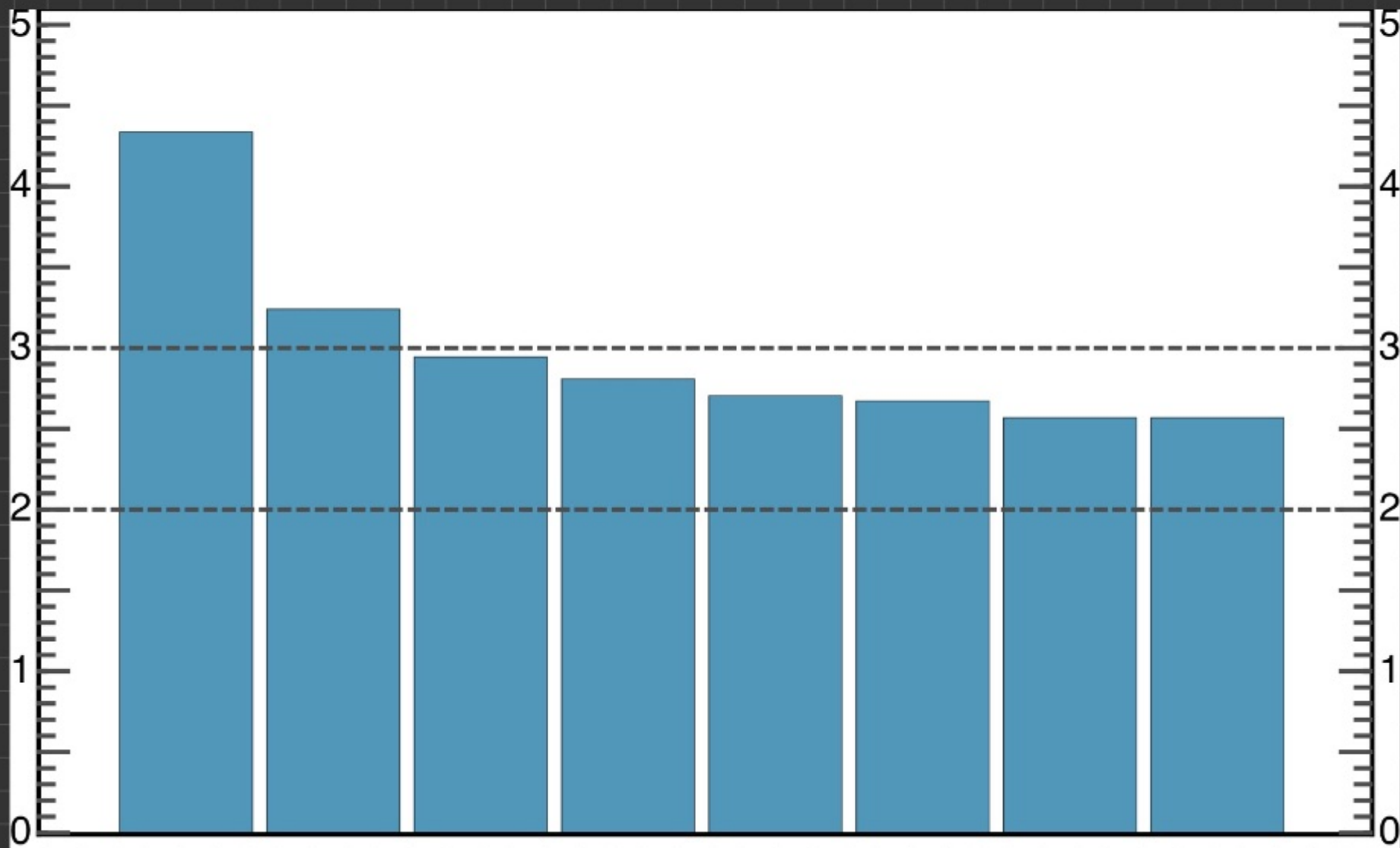
Table 1: The R -factors for $\varphi \in \mathbb{F}_{19}$. Note the factorisations into two quadrics for the five values $\varphi = 4, 5, 8, 9, 11$.

Corollary to the Chebotariev-theorem

Over $\mathbb{F}_p$, an irreducible polynomial $f(x)$ with integer coefficients has on average one root (over $\mathbb{F}_p$)

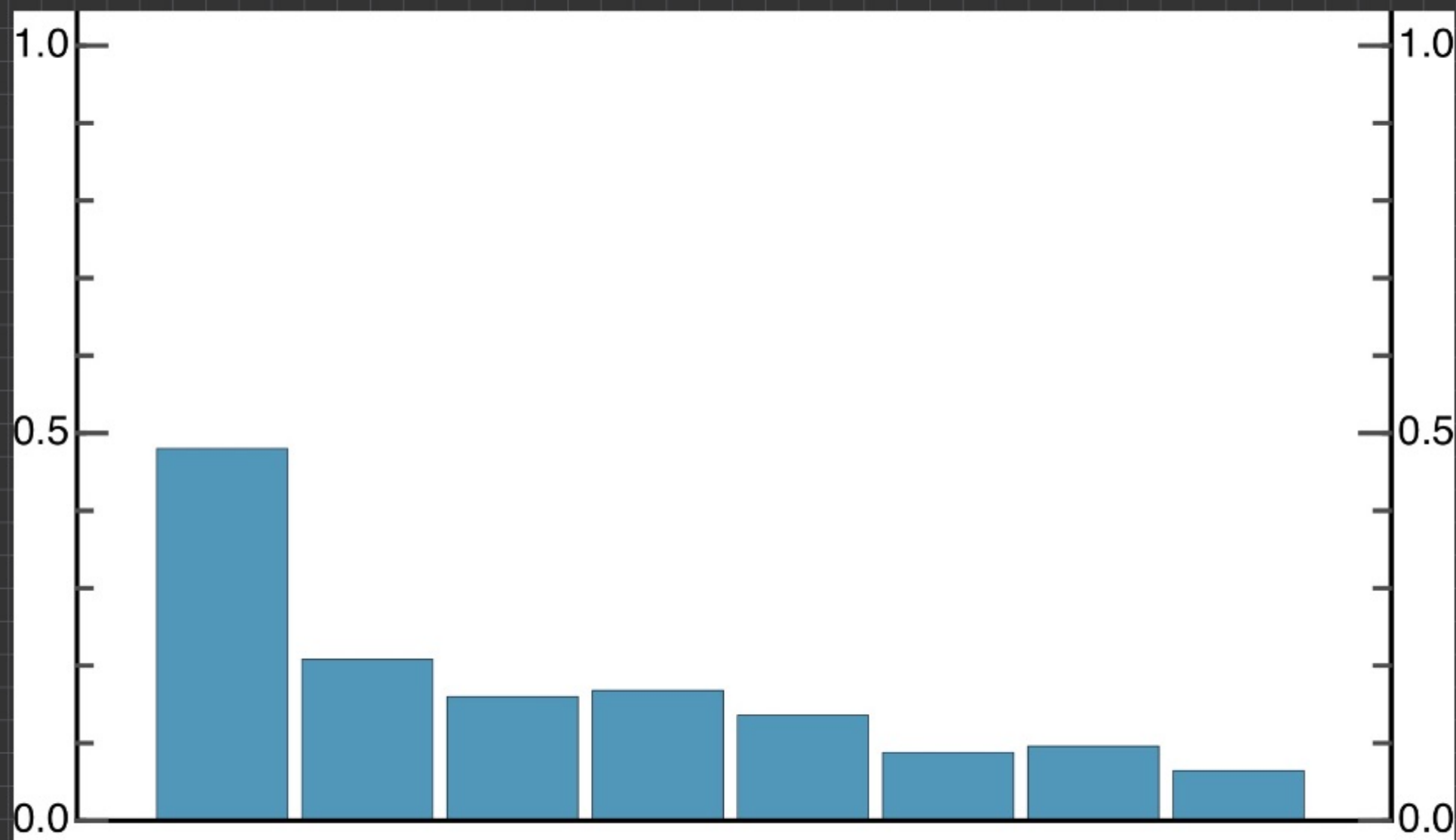
$\Rightarrow$ if $G(x)$ has k irreducible factors it has on average k roots

For Verrill's manifold



→ 1000 pri ms
175

For the mirror quintic



There is more information in the tables:
there are modular forms

$$R(T) = (1 - \rho \alpha_p T + p^3 T^2)(1 - \beta_p T + p^3 T^2)$$

Serre's conjecture (generalizing Taniyama - Weil)

↳ "motives" of length two are modular
 ↳ algebraically defined part of cohomology

[proof: Dieulefait, Khare & Wintenberger, Kisin]

For $\psi = -1/7$

α_p & β_p are Fourier coefficients of a modular form for $\Gamma_0(14)$

LMFDB

$$f_2 = \sum_n \alpha_n q^n$$

weight 2

14.2.a.a

$$f_4 = \sum_n \beta_n q^n$$

weight 4

14.4.a.a

$$q = e^{2\pi i \tau} \quad \tau \in \mathbb{H}$$

Similarly, for $\psi_{\pm}$

$\psi_p \in \mathcal{B}_p \rightarrow$ coeffs of modular forms
for $\Gamma_1(34) \subset \Gamma_0(34)$

$$SL(2, \mathbb{Z}) \ni \begin{pmatrix} a & b \\ c & d \end{pmatrix} \equiv \begin{pmatrix} 1 & b \\ 0 & 1 \end{pmatrix} \pmod{34}$$

$$f_2 \rightarrow 34.2.b.a$$

$$f_4 \rightarrow 34.4.b.a$$

Associated with these modular forms are L -functions

Mellin
transform
of f

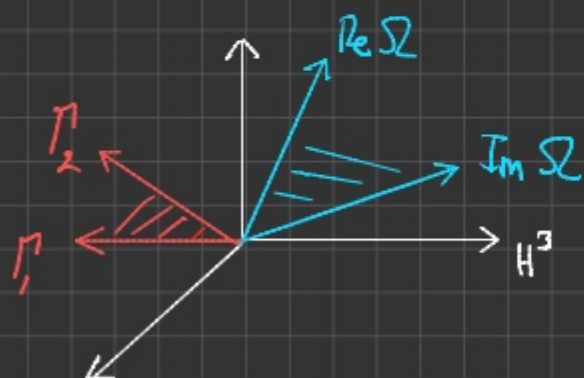
$$L_w(s) = \frac{(2\pi)^s}{\Gamma(s)} \int_0^\infty dy \, y^{s-1} f_w(iy)$$

$$w=2,4$$

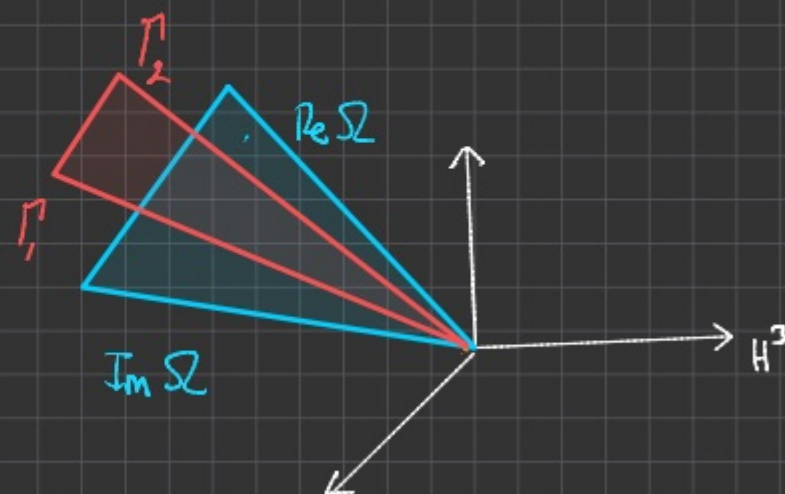
At φ_v the periods of Ω are given by
simple rational multiples of $L_4(1)$ and $L_4(2)$

[$L_w(s, f_w)$ critical values at $s=1, 2, \dots, w-1$]

Deligne: critical values of L related to periods of modular c/s]



$\varphi \rightarrow \varphi_v$
~~~~~





For example

$$\pi \left( -\frac{1}{7} \right) = i \frac{L_4(1)}{4\pi} \begin{pmatrix} 8K \\ -30K \\ 0 \\ 5 \end{pmatrix} + \frac{7}{2} \frac{L_4(2)}{\pi^2} \begin{pmatrix} 0 \\ 0 \\ 2 \\ 1 \end{pmatrix}$$

The two integral vectors define a lattice plane

$\Rightarrow$  a two dimensional family of charge vectors.

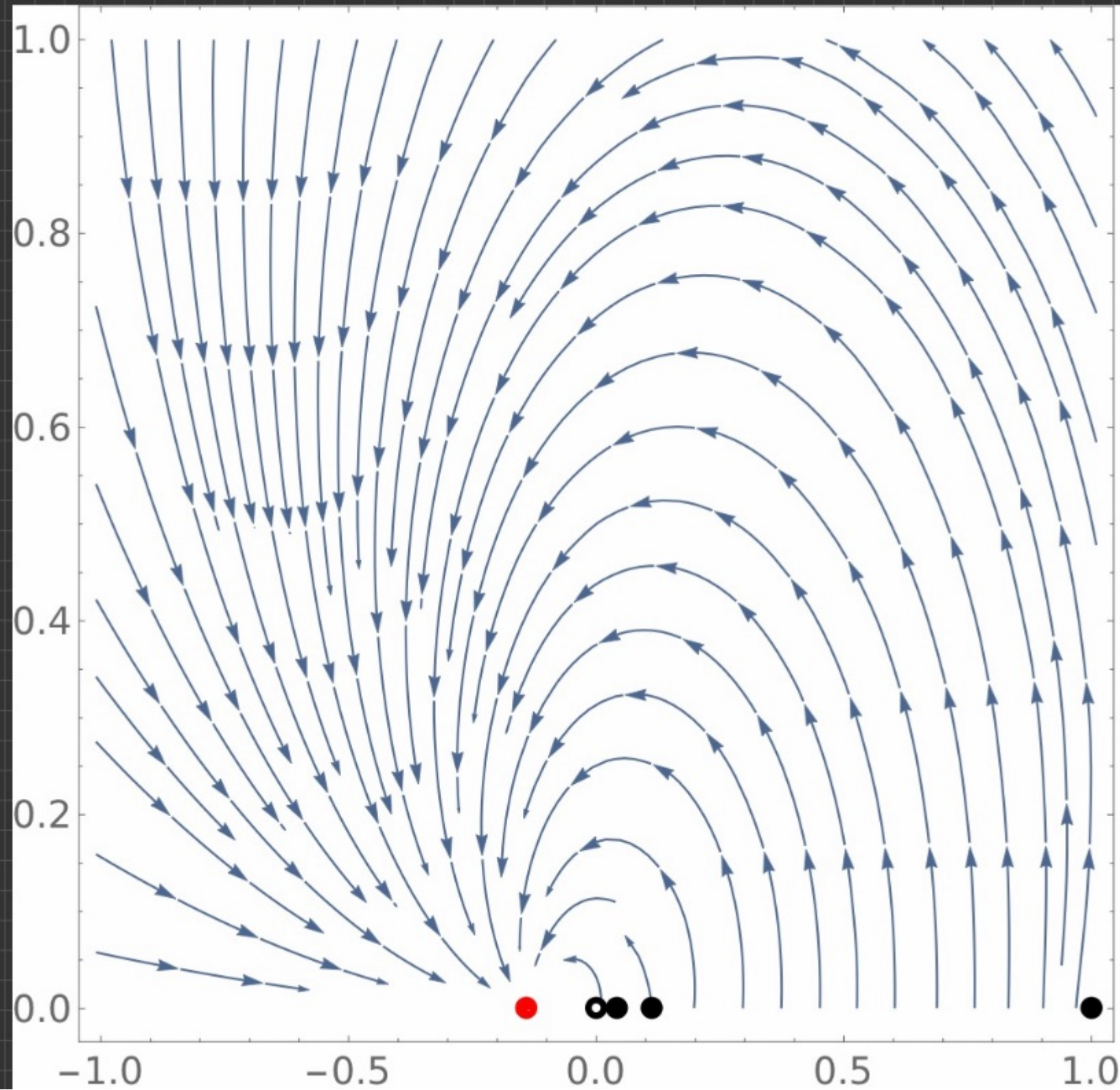
$$K = \begin{cases} 1 & \text{Verrill}/G_1 & h^{21}=1 & h^{11}=9 \\ 2 & \text{Verrill}/G_1 \times G_2 & h^{21}=1 & h^{11}=5 \end{cases}$$

$$P = \left( \sum_{i=1}^r x_i \right) \left( \sum_{i=1}^r \frac{a_i}{x_i} \right) - \psi \quad (x_1, \dots, x_r) \in \mathbb{P}^r, \quad \prod x_i \neq 0$$

$$G_1: x_i \longmapsto x_{i+1} \quad \mathbb{Z}_r$$

$$G_2: x_i \longmapsto \frac{a_i}{x_i} \quad \mathbb{Z}_2$$

$$\varphi = -\frac{1}{7}$$



## Area of the horizon of the BH

$$\varphi = -1/7$$

Chargess  $Q_{kl} = k(4, -15, -5, 0) + l(0, 0, 2, 1)$

(so a two parameter family of BHs)

let  $V_* = \frac{7}{\pi} \frac{L_4(2)}{L_4(1)}$

$L_4 \rightarrow$  L-function associated to  $f_4$

$$L_4(s) = \frac{(2\pi)^s}{\Gamma(s)} \int_0^\infty dy y^{s-1} f_4(iy)$$

Then

$$A_{kl}(\varphi_*) = 4\pi |Z_8(\varphi_*)|^2 = 14\pi \left\{ k^2 V_* + \left( l - \frac{5k}{2} \right)^2 \frac{1}{V_*} \right\}$$

↗ BH entropy

What is the precise meaning of this?



## Area of the horizon of the BH

$$\varphi_{\pm} = 33 \pm 8\sqrt{17}$$

$$\varphi_+ \quad Q_{ke} = k(4, -9, 7, 4) + e(4, -30, -30, -5)$$

$$\varphi_- \quad Q_{ke} = k(-2, 0, 0, 5) + e(0, 3, 1, 0)$$

$$v_* = \frac{17}{2} \left( \frac{9 - \sqrt{17}}{2} \right) \frac{\lambda_4(2)}{\pi \lambda_4(1)} \quad (\lambda_4 = \rho_e L_4)$$

$$A_{ke}(\varphi_{\pm}) = 34\pi \left( \frac{k^2}{v_*} + e^2 v_* \right)$$



# Outlook

- arithmetic strategy to study the Hodge structure
  - **Hope** to understand HS better eg using the formulas for  $N_q(\psi)$  explicitly in terms of the periods
  - **attractor mechanism**  $\rightarrow$  **still open questions**
    - the meaning of the formula for the area of the BH solns.
    - what is the property of a given family of CYs that allows for rank 2 attractor points in its moduli space
    - rigorous proof?
      - $\rightarrow$  find the geometric reason for the splitting of the Hodge structure
      - $\rightarrow$  Donnigan 2026 proof for HV  $\psi = -1/7$
    - other splittings/degenerations?

## • IIB [X] flux vacua

Conjecture: (Kachru, Nally, Yang) A CY variety  $X_\varphi$  with dS  $\varphi$  which gives a supersymmetric flux vacuum is modular in the sense that it is associated to a modular form of  $w=2$

(..... Candelas, XD, M. Goodwin, Kuusela 2023)

## ► Modularity of CY varieties!

1-parameter families  $\leftrightarrow$  paramodular forms of  $\Gamma(N) \subseteq \mathrm{Sp}(4)^v$

$\rightarrow$  at certain values of the parameter this might "degenerate" into **known** modular forms of  $\Gamma_0(N) \subseteq \mathrm{SL}(2, \mathbb{Z})$

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THANKS!