

The BV formalism for finite spectral triple: towards the quantum case



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Non-commutative geometry meets topological recursion

ESI, 24 April 2023



$$\begin{aligned} & \text{Left side: } \frac{\partial}{\partial x} \left(\frac{\partial}{\partial x} \right) = \frac{\partial^2}{\partial x^2} \\ & \text{Right side: } \frac{\partial}{\partial x} \left(1 - \delta(x) \right) \cdot \frac{\partial}{\partial x} \left(1 - \delta(x) \right) = \frac{\partial}{\partial x} \left(1 - 2\delta(x) + \delta^2(x) \right) = \frac{\partial}{\partial x} \left(1 - 2\delta(x) \right) = 2\delta'(x) \\ & f_{\delta} : \mathcal{S}(F) \rightarrow \mathcal{S}(F), \quad f_{\delta}(\phi)(x) = \delta(x) \cdot \phi(x) \\ & \text{Left side: } \delta_1(\delta_1 \circ \phi) = \delta_1(\phi) \\ & \text{Right side: } \delta_1(f_{\delta}(\phi)) = \delta_1(\delta_1(\phi)) = \delta_1(\phi) \\ & \delta_1 \circ \delta_2 = \delta_2 \circ \delta_1 = \text{id} \\ & \text{Left side: } \delta_1 \circ \delta_2 = \delta_1(f_{\delta_2}(\phi)) = \delta_1(\delta_2(\phi)) = \delta_2(\phi) \\ & \text{Right side: } \delta_1 \circ \delta_2 = \delta_1(\delta_2(\phi)) = \delta_2(\phi) \\ & \delta_1 \circ \delta_2 = \delta_2 \circ \delta_1 = \text{id} \end{aligned}$$



Standard Model of Elementary Particles									
CHARGES	Strong interactions (hadronic decoupling)					Weak interactions (lepton decoupling)			
	1	2	3	4	5	1	2	3	4
LEPTONS	e ⁻	μ ⁻	τ ⁻	neutrino	neutrino	W ⁺	W ⁻	Z boson	scalar bosons
QUARKS	d	s	u	c	t	u	d	u	u
Gauge bosons	gluon	gluon	gluon	gluon	gluon	W boson	W boson	Z boson	Z boson

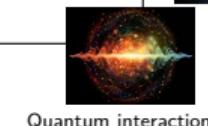
The BV formalism: a cohomological approach to gauge theory

Def. A **gauge theory** (X_0, S_0, \mathcal{G}) consists of a configuration space X_0 , an action functional $S_0 : X_0 \rightarrow \mathbb{R}$, and a group \mathcal{G} acting on X_0 such that $S_0(g \cdot \varphi) = S_0(\varphi)$, $\forall \varphi \in X_0, \forall g \in \mathcal{G}$.

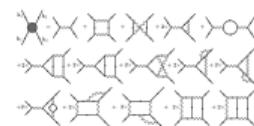


Context: quantization of a theory (X_0, S_0) via a **perturbative approach to path integral**

$$\begin{array}{c} \text{path integral} \\ \downarrow \\ Z := \int_{X_0} e^{\frac{i}{\hbar} S_0} [d\mu] \\ \xrightarrow{\hbar \rightarrow 0} \sum_{x_0 \in \{\text{crit. pts } S_0\}} e^{\frac{i}{\hbar} S_0(x_0)} |\det S_0''(x_0)|^{-\frac{1}{2}} e^{\frac{\pi i}{4} \text{sign}(S_0''(x_0))} (2\pi\hbar)^{\frac{\dim X_0}{2}} \sum_{\Gamma} \frac{\hbar^{-\chi(\Gamma)}}{|Aut(\Gamma)|} \Phi_{\Gamma}. \\ \text{partition function} \end{array}$$



Quantum interactions



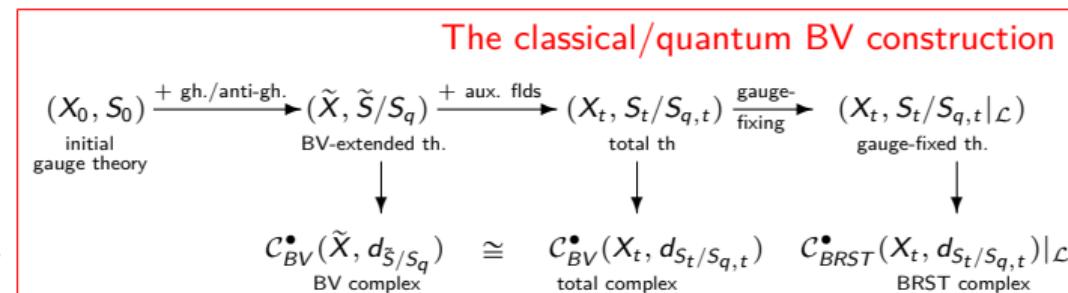
Feynman diagrams

Problem: for gauge theories, the critical points appear in *orbits*

Solution: to add extra variables, called **ghost fields**

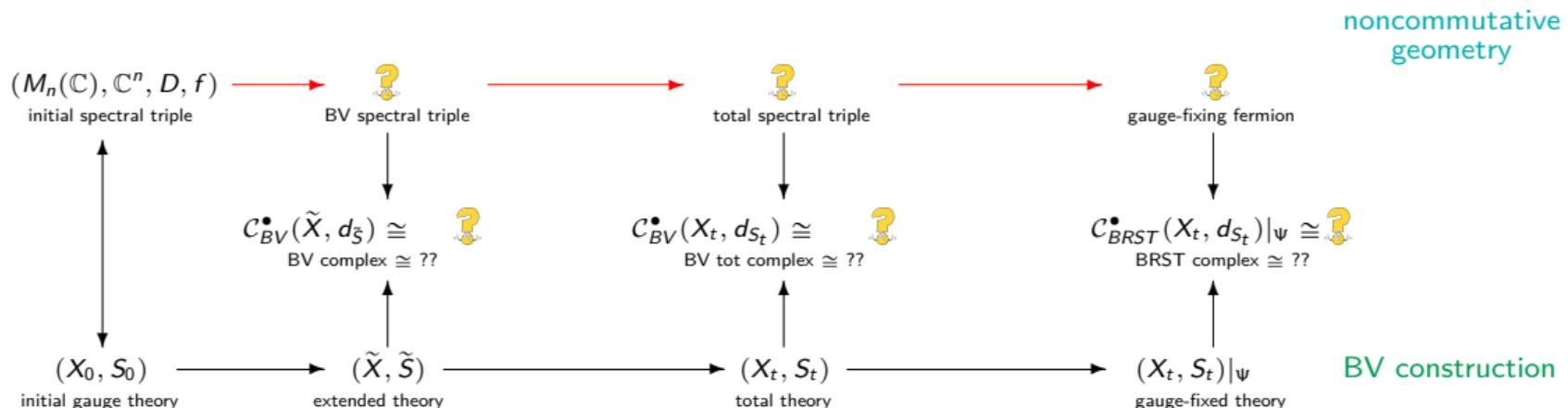
$$\begin{array}{l} \langle g \rangle := \int_{X_0} g e^{\frac{i}{\hbar} S_0} [d\mu] \\ \quad \downarrow \\ \langle \mathcal{O} \rangle := \int_{\mathcal{L} \subset \mathcal{F}_{BV}} \mathcal{O} e^{\frac{i}{\hbar} S_q|_{\mathcal{L}}} \sqrt{\mu|_{\mathcal{L}}} \end{array}$$

- ▶ \mathcal{L} Lagrangian $\subset \mathcal{F}_{BV}$ ghost sector
- ▶ $S_q \in \mathcal{C}^\infty(\mathcal{F}_{BV})[[-i\hbar]]$, sol. quant. master eq.
- ▶ $\mathcal{O} \in \mathcal{C}^\infty(\mathcal{F}_{BV})[[\hbar]]$, quantum BV cocycle



The BV formalism for spectral triples: towards the quantum case

spectral triple	gauge theory
$(\mathcal{A}, \mathcal{H}, D)$	(X_0, S_0, \mathcal{G})
<ul style="list-style-type: none"> ► \mathcal{A} = unital *-alg., $\mathcal{A} \subset \mathcal{B}(\mathcal{H})$ ► \mathcal{H} = Hilbert space ► $D : \mathcal{H} \rightarrow \mathcal{H}$ = self-adj. op. 	<ul style="list-style-type: none"> ► $X_0 = \{\varphi = \sum_j a_j [D, b_j] : \varphi^* = \varphi\} \rightsquigarrow$ conf. sp = inner fluctuations ► $S_0[D + \varphi] = \text{Tr}(f(D + \varphi)) \rightsquigarrow$ action func. = spectral action ► $\mathcal{G} = \mathcal{U}(\mathcal{A}) \rightsquigarrow$ gauge group = unitary elements in \mathcal{A}



The BV formalism for spectral triples: towards the quantum case

