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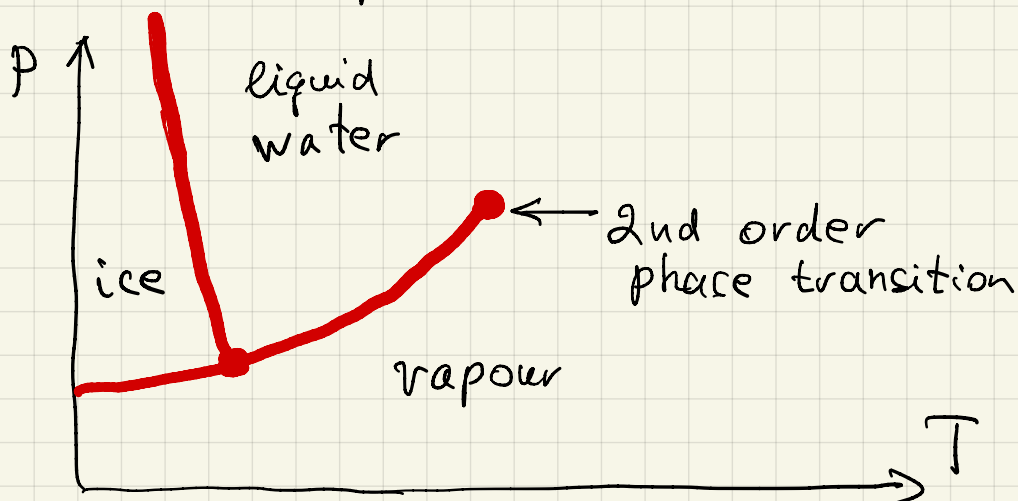
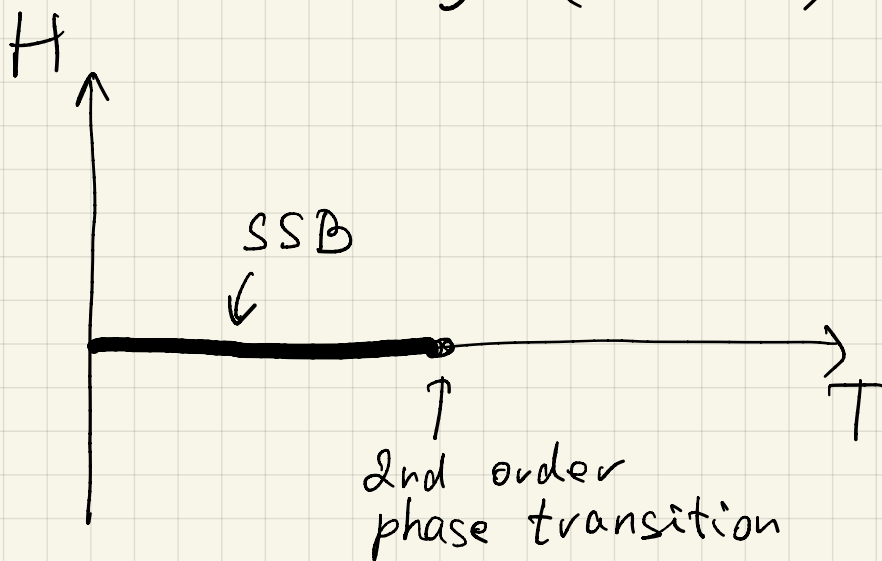


# 0. Gapped many-body systems and TQFT.

Many-body physics = physics of  $\infty$ -volume systems.

## Key facts

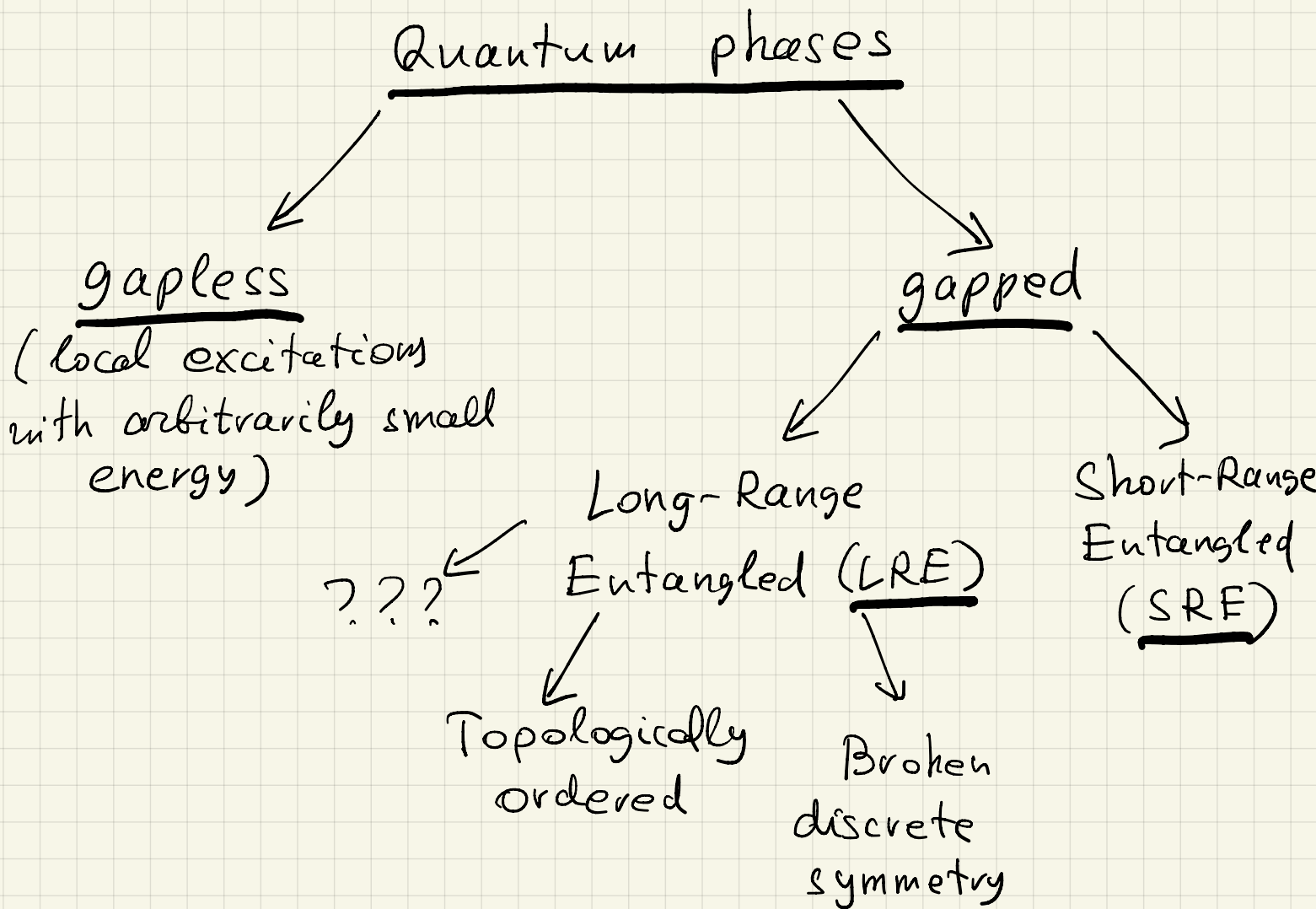
- Sharp phase transitions as one varies parameters
- Possibility of spontaneous symmetry breaking (SSB)



Quantum phase transitions = transitions  
at  $T=0$ .

Phase transitions separate phases.

What are possible phases at  $T=0$ ?



Can one use QFT to describe  
all these phases?

Topologically ordered phases  $\leftrightarrow$  non-trivial TQFT

SRE phases  $\overset{?}{\leftrightarrow}$  "almost trivial TQFT".

"almost trivial" = invertible.

### Invertible TQFT

- unique ground state on a closed spatial slice of any topology
- no non-trivial "anyons".
- partition function  $Z_X \neq 0$  for any closed Euclidean space-time  $X$ .

Unitary invertible TQFTs have been completely classified (Freed, Hopkins).

Completely determined by  $Z_X$  for all  $X$ .

E.g. let  $Z_X = \exp(i\pi \int_X W_2(TX)^2)$   
for any 4d manifold  $X$ .



which SRE phase corresponds to  
which invertible TQFT?

### Problems

- What is a "quantum phase" anyway?
- What is SRE anyway?
- How does one compute  $Z_X$  for any given SRE phase?

### Rough answers

- A quantum phase is an equivalence class of gapped lattice Hamiltonians.
- SRE phase = invertible phase  
(A. Kitaev)
- ???

a test case: Hall conductance at  $T=0$

## Expectations

- Defined for any gapped <sup>2d</sup> system with a  $U(1)$  symmetry
- a numerical invariant of a gapped phase.
- $\sigma_H \neq 0 \Rightarrow$  no gapped edge is possible
- For SRE systems  $\sigma_H = n \cdot \frac{e^2}{h}$ ,  $n \in \mathbb{Z}$
- For bosonic SRE systems  $n \in 2 \cdot \mathbb{Z}$ .

Last two properties follow from TQFT:

$$Z_X(A) = \exp\left(\frac{in}{4\pi} \int_X A dA\right) =$$

classical  
 $U(1)$  gauge field

$$= \exp\left(\frac{in}{4\pi} \int_Y F \wedge F\right)$$

$$\partial Y = X$$

$$Z_X(A) \text{ gauge-invariant} \Rightarrow n \in \mathbb{Z}$$

$$Z_X(A) \text{ independent of spin-structure}$$

$$\Rightarrow n \in 2 \cdot \mathbb{Z}$$

# 1. Quantum lattice systems and quantum phases.

Problem: define and classify possible phases of lattice systems.

$\Lambda \subset \mathbb{R}^d$  : countable uniformly discrete ("lattice").

$A = \bigotimes_{p \in \Lambda} A_p$  : algebra of observables  
( $A_p \cong \text{Mat}(n_p, \mathbb{C})$ ,  $n_p \in \mathbb{N}$ )

$A_X = \bigotimes_{p \in X} A_p \subset A$  : observables localized on  $X \subset \Lambda$ .

$H = \sum_{p \in \Lambda} H_p$  ,  $H_p \in A_{B_p(r)}$  ,  $r > 0$   
("Hamiltonian").  
↑  
ball of radius  $r$  and center at  $p$

Want to study ground states of such Hamiltonians in infinite volume, up to some equivalence.

Remark The Hilbert space where  $H$  acts is not specified from the outset. It is constructed from the ground state via the GNS construction.

## 2. Gapped vs. gapless

Want to require an energy gap  $\Delta$  between the (unique) ground state  $|0\rangle$  and excited states.

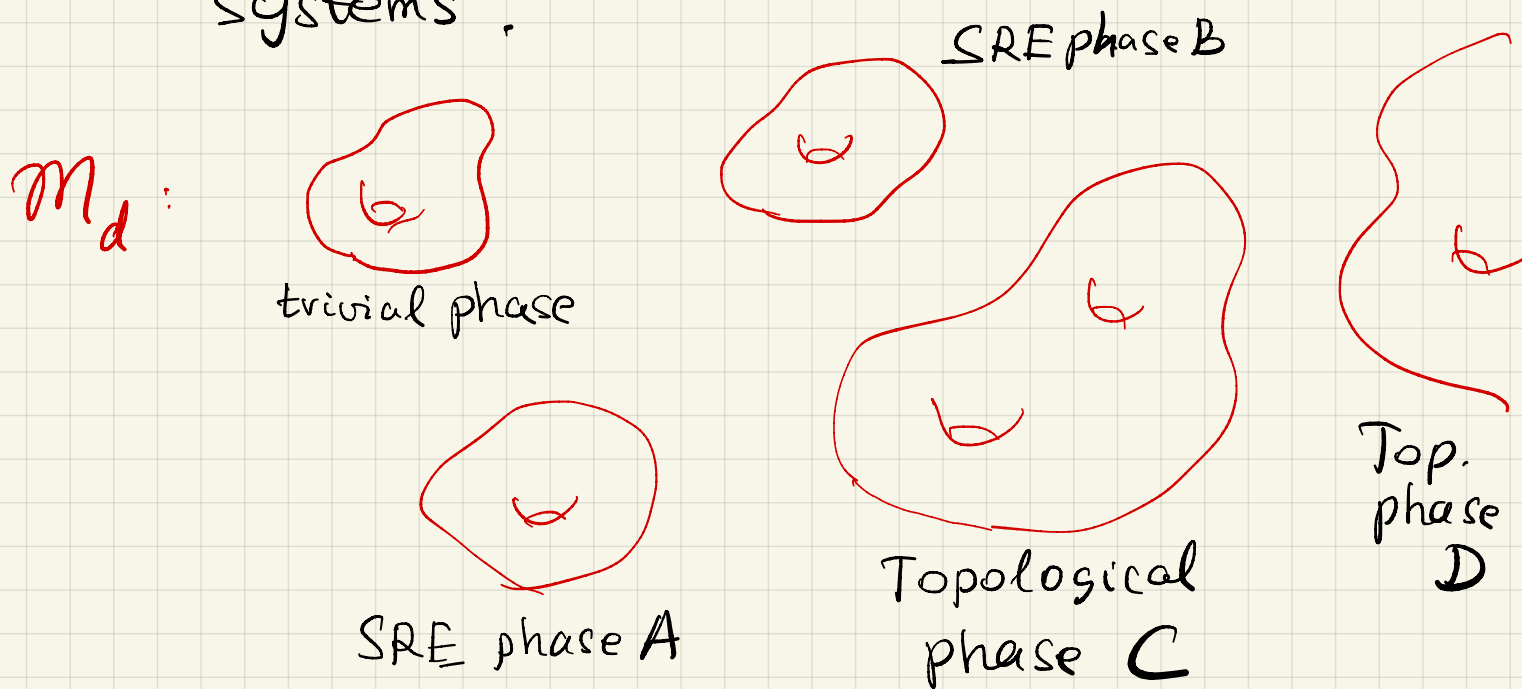
more precisely, there is a gap in the spectrum of the unbounded operator

$$\hat{H} \equiv \pi_{\text{GNS}}(H), \quad \hat{H}|0\rangle = 0.$$

Want to classify such  $H$  up to a suitable equivalence. Very hard.

Can we at least construct some "topological invariants"?

### 3. The "space" of gapped lattice systems.



$\pi_0(\mathcal{M}_d)$  = set of gapped phases in dimension  $d$ .

$$\mathcal{M}_d \rightsquigarrow \mathcal{M}_d^B \text{ (bosonic)}$$

$$\mathcal{M}_d \rightsquigarrow \mathcal{M}_d^F \text{ (fermionic)}$$

$$\pi_0(\mathcal{M}_1^F) \stackrel{?}{=} \mathbb{Z}_2 \text{ (Kitaev chain)}$$

$$\pi_0(\mathcal{M}_1^B) \stackrel{?}{=} 0$$

$$\pi_0(\mathcal{M}_2^{B/F}) \stackrel{?}{=} \mathbb{Q} \text{ (chiral central charge)}$$

Kitaev argued that spaces  $\mathcal{M}_d^{\text{SRE}}$  are all related:

$$\Omega \mathcal{M}_d^{\text{SRE}} \sim \mathcal{M}_{d-1}^{\text{SRE}} \quad (\Rightarrow \quad \Omega^k \mathcal{M}_d^{\text{SRE}} \sim \mathcal{M}_{d-k}^{\text{SRE}})$$

That is, they form a "loop spectrum".

"Practical" consequence:

$$\pi_n(\mathcal{M}_d^{\text{SRE}}) = \pi_n(\Omega^k \mathcal{M}_{d+k}^{\text{SRE}}) = \pi_{n+k}(\mathcal{M}_{d+k}^{\text{SRE}}) \\ \forall k > 0.$$

But we are still very far from proving this.

#### 4. Topological invariants of families of gapped lattice systems.

Choose a connected component of  $\mathcal{M}_d$ .

Say, the one in a trivial phase.

What are its homotopy groups?

What are its (co)homology groups?

What is its homotopy type?

How does one detect the nontrivial topology of the space of systems in a particular gapped phase?

Let's look at the case  $d=0$  first.

## 5. Berry connection, Berry curvature, Berry class.

Consider a Hamiltonian  $H(\lambda)$  depending on parameters  $\lambda \in M$ .

The Hilbert space  $V$  is fixed (is a trivial bundle over  $M$ ).

Ground state is unique  $\forall \lambda \in M \Rightarrow$  get a line bundle  $B$  over  $M$ .

- Canonical connection  $\nabla$  on  $B$
- $\frac{1}{i} \nabla^2 = F \in \Omega^2(M)$  is a closed form with periods  $\in 2\pi \mathbb{Z}$
- $[\frac{F}{2\pi}] = c_1(B)$ .

If  $[\frac{F}{2\pi}] \neq 0$ , the family of Hamiltonians is topologically non-trivial.



$\left[\frac{F}{2\pi}\right]$  is a complete invariant of a family.

Space of  $d=0$  Hamiltonians with a unique ground state  $\sim$  Space of rank-1 projectors  $\sim U(V)/U(1)$   $\sim$  odd Hilbert space  $\downarrow$

$U(V)/U(1) = PU(V)$  is a  $K(\mathbb{Z}, 2)$ :

$$\pi_k(PU(V)) = \begin{cases} \mathbb{Z}, & k=2 \\ 0, & k \neq 2 \end{cases}$$

$$\Rightarrow \pi(M, PU(V)) = H^2(M, \mathbb{Z}).$$

(Equivalently, a line bundle of ground states is determined (topologically) by  $\left[\frac{F}{2\pi}\right]$ .)

Note:  $\mathcal{M}_0^{SRE} \sim K(\mathbb{Z}, 2)$  implies (assuming Kitaev)

$$\pi_{d+2}(\mathcal{M}_d^{SRE}) = \mathbb{Z}, \quad \pi_k(\mathcal{M}_d^{SRE}) = 0 \quad \text{if } k = d, d+1, \text{ or } k > d+2$$

In particular,  $\mathcal{M}_1^{triv} \sim K(\mathbb{Z}, 3)$ .

Can we check all this directly?