

Homotopy transfer for conserved currents and rigid symmetries in gauge theories

(work in progress)

Igor Khavkine

Institute of Mathematics
Czech Academy of Sciences, Prague

Geometry for Higher Spin Gravity
Erwin Schrödinger Institute, Vienna

07 Sep 2021

PDEs

- ▶ A PDE is $\mathcal{E} \hookrightarrow J^\infty F$, where $F \rightarrow M$ is a (ghost) graded super-bundle of fields over spacetime M , with $n = \dim M$.
- ▶ The Cartan distribution on $J^\infty F$ gives rise to d_V , d_H and evolutionary vector fields (evfs), $\mathfrak{X}_{\text{ev}}(J^\infty F)$, which commute with d_H and d_V .
- ▶ The Cartan distribution should be involutive on \mathcal{E} and local symmetries are evfs tangent to \mathcal{E} , $\mathfrak{X}_{\text{ev}}(\mathcal{E})$.
- ▶ We mostly consider horizontal forms $\Omega^{\bullet,\bullet} = \Omega^{\text{ghost}\#, \text{horiz.deg}}(J^\infty F \text{ or } \mathcal{E})$.
- ▶ **Sign Conventions:** Any formula can be written for purely even (odd) objects; general signs recovered by introducing formal parity changing parameters $(\epsilon_1, \epsilon_2, \dots)$.

Adopt $|\mathcal{L}_{(-)}| = |[-, -]_{\mathcal{L}}| = \text{even}$, $|d_V| = |d_H| = |\iota_{(-)}| = \text{odd}$,
 $\#\mathcal{L}_{(-)} = \#[-, -]_{\mathcal{L}} = \#d_V = \#d_H = \#\iota_{(-)} = 0$.

PDEs

- ▶ A PDE is $\mathcal{E} \hookrightarrow J^\infty F$, where $F \rightarrow M$ is a (ghost) graded super-bundle of fields over spacetime M , with $n = \dim M$.
- ▶ The Cartan distribution on $J^\infty F$ gives rise to d_V , d_H and evolutionary vector fields (evfs), $\mathfrak{X}_{\text{ev}}(J^\infty F)$, which commute with d_H and d_V .
- ▶ The Cartan distribution should be involutive on \mathcal{E} and local symmetries are evfs tangent to \mathcal{E} , $\mathfrak{X}_{\text{ev}}(\mathcal{E})$.
- ▶ We mostly consider horizontal forms $\Omega^{\bullet,\bullet} = \Omega^{\text{ghost}\#, \text{horiz.deg}}(J^\infty F \text{ or } \mathcal{E})$.
- ▶ **Sign Conventions:** Any formula can be written for purely even (odd) objects; general signs recovered by introducing formal parity changing parameters $(\epsilon_1, \epsilon_2, \dots)$.

Adopt $|\mathcal{L}_{(-)}| = |[-, -]_{\mathcal{L}}| = \text{even}$, $|d_V| = |d_H| = |\iota_{(-)}| = \text{odd}$,
 $\#\mathcal{L}_{(-)} = \#[-, -]_{\mathcal{L}} = \#d_V = \#d_H = \#\iota_{(-)} = 0$.

PDEs

- ▶ A PDE is $\mathcal{E} \hookrightarrow J^\infty F$, where $F \rightarrow M$ is a (ghost) graded super-bundle of fields over spacetime M , with $n = \dim M$.
- ▶ The Cartan distribution on $J^\infty F$ gives rise to d_V , d_H and evolutionary vector fields (evfs), $\mathfrak{X}_{\text{ev}}(J^\infty F)$, which commute with d_H and d_V .
- ▶ The Cartan distribution should be involutive on \mathcal{E} and local symmetries are evfs tangent to \mathcal{E} , $\mathfrak{X}_{\text{ev}}(\mathcal{E})$.
- ▶ We mostly consider horizontal forms $\Omega^{\bullet,\bullet} = \Omega^{\text{ghost}\#, \text{horiz.deg}}(J^\infty F \text{ or } \mathcal{E})$.

▶ **Sign Conventions:** Any formula can be written for purely even (odd) objects; general signs recovered by introducing formal parity changing parameters $(\epsilon_1, \epsilon_2, \dots)$.

Adopt $|\mathcal{L}_{(-)}| = |[-, -]_{\mathcal{L}}| = \text{even}$, $|d_V| = |d_H| = |\iota_{(-)}| = \text{odd}$,
 $\#\mathcal{L}_{(-)} = \#[-, -]_{\mathcal{L}} = \#d_V = \#d_H = \#\iota_{(-)} = 0$.

PDEs

- ▶ A PDE is $\mathcal{E} \hookrightarrow J^\infty F$, where $F \rightarrow M$ is a (ghost) graded super-bundle of fields over spacetime M , with $n = \dim M$.
- ▶ The Cartan distribution on $J^\infty F$ gives rise to d_V , d_H and evolutionary vector fields (evfs), $\mathfrak{X}_{\text{ev}}(J^\infty F)$, which commute with d_H and d_V .
- ▶ The Cartan distribution should be involutive on \mathcal{E} and local symmetries are evfs tangent to \mathcal{E} , $\mathfrak{X}_{\text{ev}}(\mathcal{E})$.
- ▶ We mostly consider horizontal forms $\Omega^{\bullet,\bullet} = \Omega^{\text{ghost}\#, \text{horiz.deg}}(J^\infty F \text{ or } \mathcal{E})$.
- ▶ **Sign Conventions:** Any formula can be written for purely even (odd) objects; general signs recovered by introducing formal parity changing parameters $(\epsilon_1, \epsilon_2, \dots)$.

Adopt $|\mathcal{L}_{(-)}| = |[-, -]_{\mathcal{L}}| = \text{even}$, $|d_V| = |d_H| = |\iota_{(-)}| = \text{odd}$,
 $\#\mathcal{L}_{(-)} = \#[-, -]_{\mathcal{L}} = \#d_V = \#d_H = \#\iota_{(-)} = 0$.

BV-BRST

- A (sufficiently regular) variational PDE \mathcal{E} has a BV-BRST description.
 - In the BV-BRST extension, fiber coordinates over M come in field Φ^I (F_{BRST}), antifield Φ_I^* pairs ($F_{BV} \rightarrow F_{BRST}$):

$$\begin{array}{ccccccc} F_{BV} = T_V^*[-1]F_{BRST} & \longrightarrow & F_{BRST} & \longrightarrow & F & \longrightarrow & M \\ \uparrow & & \uparrow & & \uparrow & & \\ J^\infty F_{BV} & \longrightarrow & J^\infty F_{BRST} & \longrightarrow & J^\infty F & & \\ \uparrow & & \uparrow & & \uparrow & & \\ \mathcal{E}_{BRST} & \longrightarrow & \mathcal{E} & & & & \end{array}$$
- Antifields give rise to a local shifted symplectic density $\Omega = d_V \Phi_I^* d_V \Phi^I$, together with a homological evfs $Q \in \mathfrak{X}_{ev}(J^\infty F_{BV})$, $Q_{BRST} \in \mathfrak{X}_{ev}(\mathcal{E}_{BRST})$, with ghost $\#Q = -\#\Omega = 1$, such that $H^{\# < 0}(C^\infty(J^\infty F_{BV}), \mathbf{s} = \mathcal{L}_Q) = 0$ and $H^0(C^\infty(J^\infty F_{BV}), \mathbf{s}) = H^0(C^\infty(\mathcal{E}_{BRST}), \mathcal{L}_{Q_{BRST}}) = C^\infty(\mathcal{E})^G$ with G generated by infinitesimal gauge symmetries.
- The antibracket $(\int b, \int c) = \iota_{\Omega-1} d_V b d_V c$ is defined on local functionals $H^{\bullet, n}(d_H)$.
- The BV differential $\mathbf{s}(-) = (S_{BV}, -)$, where the BV extended action $S_{BV} = \int L(\Phi, \Phi^*) d^n x \in H^{0,n}(d_H)$ satisfies the Master Equation $(S_{BV}, S_{BV}) = 0$. This makes $(H^{\bullet, n}(d_H), \mathbf{s}, (-, -))$ into a dg-Lie algebra.
- Symmetries $\beta \in \mathfrak{X}_{ev}(J^\infty F_{BV})$ must now leave $L d^n x$ and Ω invariant mod d_H .
- Conventions: $|(-, -)| = |\mathbf{s}| = \text{odd}$, $\#(-, -) = \#\mathbf{s} = 1$, with $(-\epsilon, -) = (-)^{|\epsilon|} (-, \epsilon -)$.

BV-BRST

- ▶ A (sufficiently regular) variational PDE \mathcal{E} has a BV-BRST description.
 - ▶ In the BV-BRST extension, fiber coordinates over M come in field Φ^I (F_{BRST}), antifield Φ_I^* pairs ($F_{BV} \rightarrow F_{BRST}$):
- $$\begin{array}{ccccccc}
 F_{BV} = T_V^*[-1]F_{BRST} & \longrightarrow & F_{BRST} & \longrightarrow & F & \longrightarrow & M \\
 \uparrow & & \uparrow & & \uparrow & & \\
 J^\infty F_{BV} & \longrightarrow & J^\infty F_{BRST} & \longrightarrow & J^\infty F & & \\
 \uparrow & & \downarrow & & \uparrow & & \\
 \mathcal{E}_{BRST} & \longrightarrow & \mathcal{E} & & & &
 \end{array}$$
- ▶ Antifields give rise to a local shifted symplectic density $\Omega = d_V \Phi_I^* d_V \Phi^I$, together with a homological evfs $Q \in \mathfrak{X}_{ev}(J^\infty F_{BV})$, $Q_{BRST} \in \mathfrak{X}_{ev}(\mathcal{E}_{BRST})$, with ghost $\#Q = -\#\Omega = 1$, such that $H^{\# < 0}(C^\infty(J^\infty F_{BV}), \mathbf{s} = \mathcal{L}_Q) = 0$ and $H^0(C^\infty(J^\infty F_{BV}), \mathbf{s}) = H^0(C^\infty(\mathcal{E}_{BRST}), \mathcal{L}_{Q_{BRST}}) = C^\infty(\mathcal{E})^G$ with G generated by infinitesimal gauge symmetries.
 - ▶ The antibracket $(\int b, \int c) = \iota_{\Omega-1^n} d_V b d_V c$ is defined on local functionals $H^{*,n}(d_H)$.
 - ▶ The BV differential $\mathbf{s}(-) = (S_{BV}, -)$, where the BV extended action $S_{BV} = \int L(\Phi, \Phi^*) d^n x \in H^{0,n}(d_H)$ satisfies the Master Equation $(S_{BV}, S_{BV}) = 0$. This makes $(H^{*,n}(d_H), \mathbf{s}, (-, -))$ into a dg-Lie algebra.
 - ▶ Symmetries $\beta \in \mathfrak{X}_{ev}(J^\infty F_{BV})$ must now leave $L d^n x$ and Ω invariant mod d_H .
 - ▶ Conventions: $|(-, -)| = |\mathbf{s}| = \text{odd}$, $\#(-, -) = \#\mathbf{s} = 1$, with $(-\epsilon, -) = (-)^{|\epsilon|} (-, \epsilon -)$.

BV-BRST

- ▶ A (sufficiently regular) variational PDE \mathcal{E} has a BV-BRST description.
 - ▶ In the BV-BRST extension, fiber coordinates over M come in field Φ^I (F_{BRST}), antifield Φ_I^* pairs ($F_{BV} \rightarrow F_{BRST}$):
- $$\begin{array}{ccccccc}
F_{BV} = T_V^*[-1]F_{BRST} & \longrightarrow & F_{BRST} & \longrightarrow & F & \longrightarrow & M \\
\uparrow & & \uparrow & & \uparrow & & \\
J^\infty F_{BV} & \longrightarrow & J^\infty F_{BRST} & \longrightarrow & J^\infty F & & \\
\uparrow & & \downarrow & & \uparrow & & \\
\mathcal{E}_{BRST} & \longrightarrow & \mathcal{E} & & & &
\end{array}$$
- ▶ Antifields give rise to a local shifted symplectic density $\Omega = d_V \Phi_I^* d_V \Phi^I$, together with a homological evfs $Q \in \mathfrak{X}_{ev}(J^\infty F_{BV})$, $Q_{BRST} \in \mathfrak{X}_{ev}(\mathcal{E}_{BRST})$, with ghost $\#Q = -\#\Omega = 1$, such that $H^{\# < 0}(C^\infty(J^\infty F_{BV}), \mathbf{s} = \mathcal{L}_Q) = 0$ and $H^0(C^\infty(J^\infty F_{BV}), \mathbf{s}) = H^0(C^\infty(\mathcal{E}_{BRST}), \mathcal{L}_{Q_{BRST}}) = C^\infty(\mathcal{E})^G$ with G generated by infinitesimal gauge symmetries.
 - ▶ The antibracket $(\int b, \int c) = \iota_{\Omega-1} d_V b d_V c$ is defined on local functionals $H^{\bullet, n}(d_H)$.
 - ▶ The BV differential $\mathbf{s}(-) = (S_{BV}, -)$, where the BV extended action $S_{BV} = \int L(\Phi, \Phi^*) d^n x \in H^{0,n}(d_H)$ satisfies the Master Equation $(S_{BV}, S_{BV}) = 0$. This makes $(H^{\bullet, n}(d_H), \mathbf{s}, (-, -))$ into a dg-Lie algebra.
 - ▶ Symmetries $\beta \in \mathfrak{X}_{ev}(J^\infty F_{BV})$ must now leave $L d^n x$ and Ω invariant mod d_H .
 - ▶ Conventions: $|(-, -)| = |\mathbf{s}| = \text{odd}$, $\#(-, -) = \#\mathbf{s} = 1$, with $(-\epsilon, -) = (-)^{|\epsilon|} (-, \epsilon -)$.

BV-BRST

- A (sufficiently regular) variational PDE \mathcal{E} has a BV-BRST description.
 - In the BV-BRST extension, fiber coordinates over M come in field Φ^I (F_{BRST}), antifield Φ_I^* pairs ($F_{BV} \rightarrow F_{BRST}$):
$$\begin{array}{ccccccc} F_{BV} = T_V^*[-1]F_{BRST} & \longrightarrow & F_{BRST} & \longrightarrow & F & \longrightarrow & M \\ \uparrow & & \uparrow & & \uparrow & & \\ J^\infty F_{BV} & \longrightarrow & J^\infty F_{BRST} & \longrightarrow & J^\infty F & & \\ \uparrow & & \uparrow & & \uparrow & & \\ \mathcal{E}_{BRST} & \longrightarrow & \mathcal{E} & & & & \end{array}$$
- Antifields give rise to a local shifted symplectic density $\Omega = d_V \Phi_I^* d_V \Phi^I$, together with a homological evfs $Q \in \mathfrak{X}_{ev}(J^\infty F_{BV})$, $Q_{BRST} \in \mathfrak{X}_{ev}(\mathcal{E}_{BRST})$, with ghost $\#Q = -\#\Omega = 1$, such that $H^{\# < 0}(C^\infty(J^\infty F_{BV}), \mathbf{s} = \mathcal{L}_Q) = 0$ and $H^0(C^\infty(J^\infty F_{BV}), \mathbf{s}) = H^0(C^\infty(\mathcal{E}_{BRST}), \mathcal{L}_{Q_{BRST}}) = C^\infty(\mathcal{E})^G$ with G generated by infinitesimal gauge symmetries.
- The antibracket $(\int b, \int c) = \iota_{\Omega-1} d_V b d_V c$ is defined on local functionals $H^{\bullet,n}(d_H)$.
- The BV differential $\mathbf{s}(-) = (S_{BV}, -)$, where the BV extended action $S_{BV} = \int L(\Phi, \Phi^*) d^n x \in H^{0,n}(d_H)$ satisfies the Master Equation $(S_{BV}, S_{BV}) = 0$. This makes $(H^{\bullet,n}(d_H), \mathbf{s}, (-, -))$ into a dg-Lie algebra.
- Symmetries $\beta \in \mathfrak{X}_{ev}(J^\infty F_{BV})$ must now leave $L d^n x$ and Ω invariant mod d_H .
- Conventions: $|(-, -)| = |\mathbf{s}| = \text{odd}$, $\#(-, -) = \#\mathbf{s} = 1$, with $(-\epsilon, -) = (-)^{|\epsilon|} (-, \epsilon -)$.

BV-BRST

- ▶ A (sufficiently regular) variational PDE \mathcal{E} has a BV-BRST description.
 - ▶ In the BV-BRST extension, fiber coordinates over M come in field Φ^I (F_{BRST}), antifield Φ_I^* pairs ($F_{BV} \rightarrow F_{BRST}$):
- $$\begin{array}{ccccccc}
F_{BV} = T_V^*[-1]F_{BRST} & \longrightarrow & F_{BRST} & \longrightarrow & F & \longrightarrow & M \\
\uparrow & & \uparrow & & \uparrow & & \\
J^\infty F_{BV} & \longrightarrow & J^\infty F_{BRST} & \longrightarrow & J^\infty F & & \\
\uparrow & & \downarrow & & \uparrow & & \\
\mathcal{E}_{BRST} & \longrightarrow & \mathcal{E} & & & &
\end{array}$$
- ▶ Antifields give rise to a local shifted symplectic density $\Omega = d_V \Phi_I^* d_V \Phi^I$, together with a homological evfs $Q \in \mathfrak{X}_{ev}(J^\infty F_{BV})$, $Q_{BRST} \in \mathfrak{X}_{ev}(\mathcal{E}_{BRST})$, with ghost $\#Q = -\#\Omega = 1$, such that $H^{\# < 0}(C^\infty(J^\infty F_{BV}), \mathbf{s} = \mathcal{L}_Q) = 0$ and $H^0(C^\infty(J^\infty F_{BV}), \mathbf{s}) = H^0(C^\infty(\mathcal{E}_{BRST}), \mathcal{L}_{Q_{BRST}}) = C^\infty(\mathcal{E})^G$ with G generated by infinitesimal gauge symmetries.
 - ▶ The antibracket $(\int b, \int c) = \iota_{\Omega-1} d_V b d_V c$ is defined on local functionals $H^{\bullet, n}(d_H)$.
 - ▶ The BV differential $\mathbf{s}(-) = (S_{BV}, -)$, where the BV extended action $S_{BV} = \int L(\Phi, \Phi^*) d^n x \in H^{0,n}(d_H)$ satisfies the Master Equation $(S_{BV}, S_{BV}) = 0$. This makes $(H^{\bullet, n}(d_H), \mathbf{s}, (-, -))$ into a dg-Lie algebra.
 - ▶ Symmetries $\beta \in \mathfrak{X}_{ev}(J^\infty F_{BV})$ must now leave $L d^n x$ and Ω invariant mod d_H .
 - ▶ Conventions: $|(-, -)| = |\mathbf{s}| = \text{odd}$, $\#(-, -) = \#\mathbf{s} = 1$, with $(-\epsilon, -) = (-)^{|\epsilon|} (-, \epsilon -)$.

BV-BRST

- ▶ A (sufficiently regular) **variational PDE** \mathcal{E} has a BV-BRST description.
 - ▶ In the BV-BRST extension, fiber coordinates over M come in **field** Φ^I (F_{BRST}), **antifield** Φ_I^* pairs ($F_{BV} \rightarrow F_{BRST}$):
- $$\begin{array}{ccccccc}
F_{BV} = T_V^*[-1]F_{BRST} & \longrightarrow & F_{BRST} & \longrightarrow & F & \longrightarrow & M \\
\uparrow & & \uparrow & & \uparrow & & \\
J^\infty F_{BV} & \longrightarrow & J^\infty F_{BRST} & \longrightarrow & J^\infty F & & \\
\uparrow & & \uparrow & & \uparrow & & \\
\mathcal{E}_{BRST} & \longrightarrow & \mathcal{E} & & & &
\end{array}$$
- ▶ Antifields give rise to a local shifted symplectic density $\Omega = d_V \Phi_I^* d_V \Phi^I$, together with a homological evfs $Q \in \mathfrak{X}_{ev}(J^\infty F_{BV})$, $Q_{BRST} \in \mathfrak{X}_{ev}(\mathcal{E}_{BRST})$, with ghost $\#Q = -\#\Omega = 1$, such that $H^{\# < 0}(C^\infty(J^\infty F_{BV}), \mathbf{s} = \mathcal{L}_Q) = 0$ and $H^0(C^\infty(J^\infty F_{BV}), \mathbf{s}) = H^0(C^\infty(\mathcal{E}_{BRST}), \mathcal{L}_{Q_{BRST}}) = C^\infty(\mathcal{E})^G$ with G generated by infinitesimal **gauge symmetries**.
 - ▶ The **antibracket** $(\int b, \int c) = \iota_{\Omega-1} d_V b d_V c$ is defined on **local functionals** $H^{\bullet, n}(d_H)$.
 - ▶ The BV **differential** $\mathbf{s}(-) = (S_{BV}, -)$, where the BV **extended action** $S_{BV} = \int L(\Phi, \Phi^*) d^n x \in H^{0,n}(d_H)$ satisfies the **Master Equation** $(S_{BV}, S_{BV}) = 0$. This makes $(H^{\bullet, n}(d_H), \mathbf{s}, (-, -))$ into a **dg-Lie algebra**.
 - ▶ **Symmetries** $\beta \in \mathfrak{X}_{ev}(J^\infty F_{BV})$ must now leave $L d^n x$ and Ω invariant mod d_H .
 - ▶ Conventions: $|(-, -)| = |\mathbf{s}| = \text{odd}$, $\#(-, -) = \#\mathbf{s} = 1$, with $(-\epsilon, -) = (-)^{|\epsilon|}(-, \epsilon -)$.

L_∞ -algebras

- ▶ **Def:** On a $(\mathbb{Z}, \mathbb{Z}_2)$ -graded vector space V , the $(1, \text{odd})$ -degree brackets $[\mathcal{S}(V)] \rightarrow V$ are an L_∞ -algebra when $[e^B[e^B]] = 0$ for any even $B \in V$, while $[1] = 0$ and $\epsilon[(\cdots)] = (-)^{|\epsilon|}[\epsilon(\cdots)]$, for odd ϵ .
- ▶ Writing $\mathbf{s}B := [B]$ and decoding the higher Jacobi identities, $\mathbf{s}^2B = 0$, $2[B\mathbf{s}B] + \mathbf{s}[B^2] = 0$, $3[B^2\mathbf{s}B] + 3[B[B^2]] + \mathbf{s}[B^3] = 0$, ... ,

$$\sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0.$$

- ▶ **Ex:** For a dg-Lie algebra $(\mathfrak{g}, \mathbf{s}, [-, -])$, the L_∞ -algebra is
$$(\mathfrak{g}[-1, \text{odd}], 0 \oplus \mathbf{s} \oplus [-, -] \oplus 0 \oplus \cdots).$$
- ▶ **Dually:** L_∞ -algebra $(V, [-]) \iff (\mathcal{S}(V^*), D = [-]^*)$ dgca.
- ▶ **(non-)Ex:** $(C^\infty(J^\infty F_{BV}), \mathbf{s}_{BV})$ is an L_∞ -algebra(-oid) in the dual picture. But it is **not** the L_∞ -structure that I will talk about!

L_∞ -algebras

- ▶ **Def:** On a $(\mathbb{Z}, \mathbb{Z}_2)$ -graded vector space V , the $(1, \text{odd})$ -degree brackets $[\mathcal{S}(V)] \rightarrow V$ are an L_∞ -algebra when $[e^B[e^B]] = 0$ for any even $B \in V$, while $[1] = 0$ and $\epsilon[(\cdots)] = (-)^{|\epsilon|}[\epsilon(\cdots)]$, for odd ϵ .
- ▶ Writing $\mathbf{s}B := [B]$ and decoding the **higher Jacobi** identities, $\mathbf{s}^2B = 0$, $2[B\mathbf{s}B] + \mathbf{s}[B^2] = 0$, $3[B^2\mathbf{s}B] + 3[B[B^2]] + \mathbf{s}[B^3] = 0$, ... ,

$$\sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0.$$

- ▶ **Ex:** For a dg-Lie algebra $(\mathfrak{g}, \mathbf{s}, [-, -])$, the L_∞ -algebra is
$$(\mathfrak{g}[-1, \text{odd}], 0 \oplus \mathbf{s} \oplus [-, -] \oplus 0 \oplus \cdots).$$
- ▶ **Dually:** L_∞ -algebra $(V, [-]) \iff (\mathcal{S}(V^*), D = [-]^*)$ dgca.
- ▶ **(non-)Ex:** $(C^\infty(J^\infty F_{BV}), \mathbf{s}_{BV})$ is an L_∞ -algebra(-oid) in the dual picture. But it is **not** the L_∞ -structure that I will talk about!

L_∞ -algebras

- ▶ **Def:** On a $(\mathbb{Z}, \mathbb{Z}_2)$ -graded vector space V , the $(1, \text{odd})$ -degree brackets $[\mathcal{S}(V)] \rightarrow V$ are an L_∞ -algebra when $[e^B[e^B]] = 0$ for any even $B \in V$, while $[1] = 0$ and $\epsilon[(\cdots)] = (-)^{|\epsilon|}[\epsilon(\cdots)]$, for odd ϵ .
- ▶ Writing $\mathbf{s}B := [B]$ and decoding the **higher Jacobi** identities, $\mathbf{s}^2B = 0$, $2[B\mathbf{s}B] + \mathbf{s}[B^2] = 0$, $3[B^2\mathbf{s}B] + 3[B[B^2]] + \mathbf{s}[B^3] = 0$, ... ,

$$\sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0.$$

- ▶ **Ex:** For a dg-Lie algebra $(\mathfrak{g}, \mathbf{s}, [-, -])$, the L_∞ -algebra is

$$(\mathfrak{g}[-1, \text{odd}], 0 \oplus \mathbf{s} \oplus [-, -] \oplus 0 \oplus \cdots).$$

- ▶ **Dually:** L_∞ -algebra $(V, [-]) \iff (\mathcal{S}(V^*), D = [-]^*)$ dgca.
- ▶ **(non-)Ex:** $(C^\infty(J^\infty F_{BV}), \mathbf{s}_{BV})$ is an L_∞ -algebra(-oid) in the dual picture. But it is **not** the L_∞ -structure that I will talk about!

L_∞ -algebras

- ▶ **Def:** On a $(\mathbb{Z}, \mathbb{Z}_2)$ -graded vector space V , the $(1, \text{odd})$ -degree brackets $[\mathcal{S}(V)] \rightarrow V$ are an L_∞ -algebra when $[e^B[e^B]] = 0$ for any even $B \in V$, while $[1] = 0$ and $\epsilon[(\cdots)] = (-)^{|\epsilon|}[\epsilon(\cdots)]$, for odd ϵ .
- ▶ Writing $\mathbf{s}B := [B]$ and decoding the **higher Jacobi** identities, $\mathbf{s}^2B = 0$, $2[B\mathbf{s}B] + \mathbf{s}[B^2] = 0$, $3[B^2\mathbf{s}B] + 3[B[B^2]] + \mathbf{s}[B^3] = 0$, ... ,

$$\sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0.$$

- ▶ **Ex:** For a dg-Lie algebra $(\mathfrak{g}, \mathbf{s}, [-, -])$, the L_∞ -algebra is

$$(\mathfrak{g}[-1, \text{odd}], 0 \oplus \mathbf{s} \oplus [-, -] \oplus 0 \oplus \cdots).$$

- ▶ **Dually:** L_∞ -algebra $(V, [-]) \iff (\mathcal{S}(V^*), D = [-]^*)$ dgca.
- ▶ **(non-)Ex:** $(C^\infty(J^\infty F_{BV}), \mathbf{s}_{BV})$ is an L_∞ -algebra(-oid) in the dual picture. But it is **not** the L_∞ -structure that I will talk about!

L_∞ -algebras

- ▶ **Def:** On a $(\mathbb{Z}, \mathbb{Z}_2)$ -graded vector space V , the $(1, \text{odd})$ -degree brackets $[\mathcal{S}(V)] \rightarrow V$ are an L_∞ -algebra when $[e^B[e^B]] = 0$ for any even $B \in V$, while $[1] = 0$ and $\epsilon[(\cdots)] = (-)^{|\epsilon|}[\epsilon(\cdots)]$, for odd ϵ .
- ▶ Writing $\mathbf{s}B := [B]$ and decoding the **higher Jacobi** identities, $\mathbf{s}^2B = 0$, $2[B\mathbf{s}B] + \mathbf{s}[B^2] = 0$, $3[B^2\mathbf{s}B] + 3[B[B^2]] + \mathbf{s}[B^3] = 0$, ... ,

$$\sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0.$$

- ▶ **Ex:** For a dg-Lie algebra $(\mathfrak{g}, \mathbf{s}, [-, -])$, the L_∞ -algebra is

$$(\mathfrak{g}[-1, \text{odd}], 0 \oplus \mathbf{s} \oplus [-, -] \oplus 0 \oplus \cdots).$$

- ▶ **Dually:** L_∞ -algebra $(V, [-]) \iff (\mathcal{S}(V^*), D = [-]^*)$ dgca.
- ▶ **(non-)Ex:** $(C^\infty(J^\infty F_{BV}), \mathbf{s}_{BV})$ is an L_∞ -algebra(-oid) in the dual picture. But it is **not** the L_∞ -structure that I will talk about!

Higher Structure Constants in Extended BV

- ▶ Brandt, Henneaux, Wilch (1998)
Extended antifield formalism Nucl Phys B**510** 640–656
 - ▶ Start with a BV description of a **gauge theory** and consider a basis S_A of **representatives** for $H^{•<0,n}(\mathbf{s}|d_H)$ in local functionals, $\int b(\Phi, \Phi^*) d^n x$.
 - ▶ Then take some **antibrackets** $(-, -)$, then some more antibrackets, ...

At each stage, you get new structure constants $f_{AB}^P, f_{ABC}^P, \dots$, and new local functionals S_{AB}, S_{ABC}, \dots .

- ▶ Putting the **dg-Lie** properties of \mathbf{s} and $(-, -)$ to full use, after **some magic**, they end up with ...

$$\iff \sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0 \quad L_\infty\text{-algebra identities!}$$

Higher Structure Constants in Extended BV

- ▶ Brandt, Henneaux, Wilch (1998)
Extended antifield formalism Nucl Phys B**510** 640–656
 - ▶ Start with a BV description of a **gauge theory** and consider a basis S_A of **representatives** for $H^{\bullet < 0,n}(\mathbf{s}|d_H)$ in local functionals, $\int b(\Phi, \Phi^*) d^n x$.
 - ▶ Then take some **antibrackets** $(-, -)$, then some more antibrackets, ...

At each stage, you get new structure constants $t_{AB}^P, t_{ABC}^P, \dots$, and new local functionals S_{AB}, S_{ABC}, \dots .

- ▶ Putting the **dg-Lie** properties of \mathbf{s} and $(-, -)$ to full use, after **some magic**, they end up with ...

$$\iff \sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0 \quad L_\infty\text{-algebra identities!}$$

Higher Structure Constants in Extended BV

- ▶ Brandt, Henneaux, Wilch (1998)
Extended antifield formalism Nucl Phys B**510** 640–656
 - ▶ Start with a BV description of a **gauge theory** and consider a basis S_A of **representatives** for $H^{\bullet < 0,n}(\mathbf{s}|d_H)$ in local functionals, $\int b(\Phi, \Phi^*) d^n x$.
 - ▶ Then take some **antibrackets** $(-, -)$, then some more **antibrackets**, ...

$$(-1)^{e_A}(S_A, S_B) = f_{AB}^D S_D + (S, S_{AB}) \quad (3.3)$$

At each stage, you get new **structure constants** $f_{AB}^D, f_{ABC}^D, \dots$, and new local functionals S_{AB}, S_{ABC}, \dots .

- ▶ Putting the **dg-Lie** properties of \mathbf{s} and $(-, -)$ to full use, after **some magic**, they end up with ...

$$\iff \sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0 \quad L_\infty\text{-algebra identities!}$$

Higher Structure Constants in Extended BV

- ▶ Brandt, Henneaux, Wilch (1998)
Extended antifield formalism Nucl Phys B**510** 640–656
 - ▶ Start with a BV description of a **gauge theory** and consider a basis S_A of **representatives** for $H^{\bullet < 0,n}(\mathbf{s}|d_H)$ in local functionals, $\int b(\Phi, \Phi^*) d^n x$.
 - ▶ Then take some **antibrackets** $(-, -)$, then some more **antibrackets**, ...

$$(-1)^{e_A}(S_A, S_B) = f_{AB}^D S_D + (S, S_{AB}) \quad (3.3)$$

$$(-)^{e_A}(S_{[A}, S_{BC]}) = S_{D[C} f_{AB]}^D + \frac{1}{3} f_{ABC}^D S_D + \frac{1}{3} (S, S_{ABC}) \quad (3.5)$$

At each stage, you get new **structure constants** $f_{AB}^D, f_{ABC}^D, \dots$, and new local functionals S_{AB}, S_{ABC}, \dots .

- ▶ Putting the **dg-Lie** properties of \mathbf{s} and $(-, -)$ to full use, after **some magic**, they end up with ...

$$\iff \sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0 \quad L_\infty\text{-algebra identities!}$$

Higher Structure Constants in Extended BV

- ▶ Brandt, Henneaux, Wilch (1998)
Extended antifield formalism Nucl Phys B**510** 640–656
 - ▶ Start with a BV description of a **gauge theory** and consider a basis S_A of **representatives** for $H^{\bullet < 0,n}(\mathbf{s}|d_H)$ in local functionals, $\int b(\Phi, \Phi^*) d^n x$.
 - ▶ Then take some **antibrackets** $(-, -)$, then some more **antibrackets**, ...

$$(-1)^{e_A}(S_A, S_B) = f_{AB}^D S_D + (S, S_{AB}) \quad (3.3)$$

$$(-)^{e_A}(S_{[A}, S_{BC]}) = S_{D[C} f_{AB]}^D + \frac{1}{3} f_{ABC}^D S_D + \frac{1}{3} (S, S_{ABC}) \quad (3.5)$$

At each stage, you get new **structure constants** $f_{AB}^D, f_{ABC}^D, \dots$, and new local functionals S_{AB}, S_{ABC}, \dots .

- ▶ Putting the **dg-Lie** properties of \mathbf{s} and $(-, -)$ to full use, after **some magic**, they end up with ...

$$\iff \sum_{k=1}^n \frac{n!}{(n-k)!k!} [B^{n-k}[B^k]] = 0 \quad L_\infty\text{-algebra identities!}$$

Higher Structure Constants in Extended BV

- ▶ Brandt, Henneaux, Wilch (1998)
Extended antifield formalism Nucl Phys B**510** 640–656
 - ▶ Start with a BV description of a **gauge theory** and consider a basis S_A of **representatives** for $H^{\bullet < 0,n}(\mathbf{s}|d_H)$ in local functionals, $\int b(\Phi, \Phi^*) d^n x$.
 - ▶ Then take some **antibrackets** $(-, -)$, then some more **antibrackets**, ...

$$(-1)^{e_A} (S_A, S_B) = f_{AB}^D S_D + (S, S_{AB}) \quad (3.3)$$

$$(-)^{e_A} (S_{[A}, S_{BC]}) = S_{D[C} f_{AB]}^D + \frac{1}{3} f_{ABC}^D S_D + \frac{1}{3} (S, S_{ABC}) \quad (3.5)$$

At each stage, you get new **structure constants** $f_{AB}^D, f_{ABC}^D, \dots$, and new local functionals S_{AB}, S_{ABC}, \dots .

- ▶ Putting the **dg-Lie** properties of **s** and $(-, -)$ to full use, after **some magic**, they end up with ...

$$\sum_{r=2}^{p-1} \frac{1}{r! (p-r)!} f_{C[A_{r+1} \dots A_p}^D f_{A_1 \dots A_r]}^C = 0, \quad (3.8)$$

$$\iff \sum_{k=1}^n \frac{n!}{(n-k)! k!} [B^{n-k} [B^k]] = 0 \quad L_\infty\text{-algebra identities!}$$

Homotopy Transfer

- ▶ Dually, an L_∞ -morphism $\lambda^*: (\mathcal{S}(V_2^*), D_2) \rightarrow (\mathcal{S}(V_1^*), D_1)$ is a dgca-morphism, with $\lambda(1) = 0$. Equivalently, for even $B \in V_1$:

$$[e^{\lambda(e^B)}]_2 = \lambda(e^B[e^B]_1).$$

- ▶ Expansion: $s_2\lambda(B) = \lambda(s_1B)$ then ...

- ▶ **Homotopy transfer** (Thm 10.3.{1,7}, Loday-Valette Algebraic Operads): Given an L_∞ -algebra $(V_2, [-]_2)$ and a quasi-isomorphism $(V_1, s_1) \rightarrow (V_2, s_2)$ of dg-vector spaces, it can be extended to an L_∞ -morphism $(V_1, [-]_1 = s_1 + ?) \rightarrow (V_2, [-]_2)$. The $[-]_1$ can be built "explicitly" provided the q-iso is presented by a zigzag of (dg-vector) homotopy retracts, using bar and co-bar constructions.
- ▶ Homotopy transfer interpretation of Brandt-Henneaux-Wilch'98:

$$(H^{<0,n}(\mathbf{s}|d_H), 0 \oplus f_{AB}^D \oplus f_{ABC}^D \oplus \dots) \xrightarrow{S_A, S_{AB}, S_{ABC}, \dots} (H^{<0,n}(d_H), \mathbf{s} \oplus (-, -))$$

Homotopy Transfer

- ▶ Dually, an L_∞ -morphism $\lambda^*: (\mathcal{S}(V_2^*), D_2) \rightarrow (\mathcal{S}(V_1^*), D_1)$ is a dgca-morphism, with $\lambda(1) = 0$. Equivalently, for even $B \in V_1$:

$$[e^{\lambda(e^B)}]_2 = \lambda(e^B[e^B]_1).$$

- ▶ Expansion: $\mathbf{s}_2\lambda(B) = \lambda(\mathbf{s}_1B)$, then ...

- ▶ **Homotopy transfer** (Thm 10.3.{1,7}, Loday-Valette Algebraic Operads): Given an L_∞ -algebra $(V_2, [-]_2)$ and a quasi-isomorphism $(V_1, \mathbf{s}_1) \rightarrow (V_2, \mathbf{s}_2)$ of dg-vector spaces, it can be extended to an L_∞ -morphism $(V_1, [-]_1 = \mathbf{s}_1 + ?) \rightarrow (V_2, [-]_2)$. The $[-]_1$ can be built "explicitly" provided the q-iso is presented by a zigzag of (dg-vector) homotopy retracts, using bar and co-bar constructions.
- ▶ Homotopy transfer interpretation of Brandt-Henneaux-Wilch'98:

$$(H^{<0,n}(\mathbf{s}|d_H), 0 \oplus f_{AB}^D \oplus f_{ABC}^D \oplus \dots) \xrightarrow{S_A, S_{AB}, S_{ABC}, \dots} (H^{<0,n}(d_H), \mathbf{s} \oplus (-, -))$$

Homotopy Transfer

- ▶ Dually, an L_∞ -morphism $\lambda^*: (\mathcal{S}(V_2^*), D_2) \rightarrow (\mathcal{S}(V_1^*), D_1)$ is a dgca-morphism, with $\lambda(1) = 0$. Equivalently, for even $B \in V_1$:

$$[e^{\lambda(e^B)}]_2 = \lambda(e^B[e^B]_1).$$

- ▶ Expansion: $\mathbf{s}_2\lambda(B) = \lambda(\mathbf{s}_1B)$, then ...

$$\begin{aligned} \frac{1}{2}(\mathbf{s}_2\lambda(B^2) + [\lambda(B)^2]_2) &= \frac{1}{2}\lambda(2Bs_1B + [B^2]_1) \\ (-1)^{e_A}(S_A, S_B) &= f_{AB}^D S_D + (S, S_{AB}) \quad (3.3) \end{aligned}$$

- ▶ **Homotopy transfer** (Thm 10.3.{1,7}, Loday-Valette Algebraic Operads): Given an L_∞ -algebra $(V_2, [-]_2)$ and a quasi-isomorphism $(V_1, \mathbf{s}_1) \rightarrow (V_2, \mathbf{s}_2)$ of dg-vector spaces, it can be extended to an L_∞ -morphism $(V_1, [-]_1 = \mathbf{s}_1 + ?) \rightarrow (V_2, [-]_2)$. The $[-]_1$ can be built "explicitly" provided the q-iso is presented by a zigzag of (dg-vector) homotopy retracts, using bar and co-bar constructions.
- ▶ Homotopy transfer interpretation of Brandt-Henneaux-Wilch'98:

$$(H^{<0,n}(\mathbf{s}|d_H), 0 \oplus f_{AB}^D \oplus f_{ABC}^D \oplus \dots) \xrightarrow{S_A, S_{AB}, S_{ABC}, \dots} (H^{<0,n}(d_H), \mathbf{s} \oplus (-, -))$$

Homotopy Transfer

- ▶ Dually, an L_∞ -morphism $\lambda^*: (\mathcal{S}(V_2^*), D_2) \rightarrow (\mathcal{S}(V_1^*), D_1)$ is a dgca-morphism, with $\lambda(1) = 0$. Equivalently, for even $B \in V_1$:

$$[e^{\lambda(e^B)}]_2 = \lambda(e^B[e^B]_1).$$

- ▶ Expansion: $\mathbf{s}_2\lambda(B) = \lambda(\mathbf{s}_1B)$, then . . .

$$\frac{1}{3!}(\mathbf{s}_2\lambda(B^3) + [3\lambda(B^2)\lambda(B)]_2 + [\lambda(B)^3]_2) = \frac{1}{3!}\lambda(3B^2\mathbf{s}_1B + 3B[B^2]_1 + [B^3]_1)$$

$$(-)^{\varepsilon_A}(S_{[A}, S_{BC]}) = S_{D[C} f_{AB]}^D + \frac{1}{3}f_{ABC}^D S_D + \frac{1}{3}(S, S_{ABC}) \quad (3.5)$$

- ▶ **Homotopy transfer** (Thm 10.3.{1,7}, Loday-Valette Algebraic Operads): Given an L_∞ -algebra $(V_2, [-]_2)$ and a quasi-isomorphism $(V_1, \mathbf{s}_1) \rightarrow (V_2, \mathbf{s}_2)$ of dg-vector spaces, it can be extended to an L_∞ -morphism $(V_1, [-]_1 = \mathbf{s}_1 + ?) \rightarrow (V_2, [-]_2)$. The $[-]_1$ can be built “explicitly” provided the q-iso is presented by a zigzag of (dg-vector) homotopy retracts, using bar and co-bar constructions.
- ▶ Homotopy transfer interpretation of Brandt-Henneaux-Wilch'98:

$$(H^{<0,n}(\mathbf{s}|d_H), 0 \oplus f_{AB}^D \oplus f_{ABC}^D \oplus \dots) \xrightarrow{S_A, S_{AB}, S_{ABC}, \dots} (H^{<0,n}(d_H), \mathbf{s} \oplus (-, -))$$

Homotopy Transfer

- ▶ Dually, an L_∞ -morphism $\lambda^*: (\mathcal{S}(V_2^*), D_2) \rightarrow (\mathcal{S}(V_1^*), D_1)$ is a dgca-morphism, with $\lambda(1) = 0$. Equivalently, for even $B \in V_1$:

$$[e^{\lambda(e^B)}]_2 = \lambda(e^B[e^B]_1).$$

- ▶ Expansion: $\mathbf{s}_2\lambda(B) = \lambda(\mathbf{s}_1B)$, then . . .

$$\frac{1}{3!}(\mathbf{s}_2\lambda(B^3) + [3\lambda(B^2)\lambda(B)]_2 + [\lambda(B)^3]_2) = \frac{1}{3!}\lambda(3B^2\mathbf{s}_1B + 3B[B^2]_1 + [B^3]_1)$$

$$(-)^{\varepsilon_A}(S_{[A}, S_{BC]}) = S_{D[C} f_{AB]}^D + \frac{1}{3}f_{ABC}^D S_D + \frac{1}{3}(S, S_{ABC}) \quad (3.5)$$

- ▶ **Homotopy transfer** (Thm 10.3.{1,7}, Loday-Valette **Algebraic Operads**): Given an L_∞ -algebra $(V_2, [-]_2)$ and a quasi-isomorphism $(V_1, \mathbf{s}_1) \rightarrow (V_2, \mathbf{s}_2)$ of dg-vector spaces, it can be extended to an L_∞ -morphism $(V_1, [-]_1 = \mathbf{s}_1 + ?) \rightarrow (V_2, [-]_2)$. The $[-]_1$ can be built “explicitly” provided the **q-iso** is presented by a **zigzag** of (dg-vector) **homotopy retracts**, using **bar** and **co-bar** constructions.
- ▶ Homotopy transfer interpretation of Brandt-Henneaux-Wilch'98:

$$(H^{<0,n}(\mathbf{s}|d_H), 0 \oplus f_{AB}^D \oplus f_{ABC}^D \oplus \dots) \xrightarrow{S_A, S_{AB}, S_{ABC}, \dots} (H^{<0,n}(d_H), \mathbf{s} \oplus (-, -))$$

Homotopy Transfer

- ▶ Dually, an L_∞ -morphism $\lambda^*: (\mathcal{S}(V_2^*), D_2) \rightarrow (\mathcal{S}(V_1^*), D_1)$ is a dgca-morphism, with $\lambda(1) = 0$. Equivalently, for even $B \in V_1$:

$$[e^{\lambda(e^B)}]_2 = \lambda(e^B[e^B]_1).$$

- ▶ Expansion: $\mathbf{s}_2\lambda(B) = \lambda(\mathbf{s}_1B)$, then . . .

$$\frac{1}{3!}(\mathbf{s}_2\lambda(B^3) + [3\lambda(B^2)\lambda(B)]_2 + [\lambda(B)^3]_2) = \frac{1}{3!}\lambda(3B^2\mathbf{s}_1B + 3B[B^2]_1 + [B^3]_1)$$

$$(-)^{\varepsilon_A}(S_{[A}, S_{BC]}) = S_{D[C} f_{AB]}^D + \frac{1}{3}f_{ABC}^D S_D + \frac{1}{3}(S, S_{ABC}) \quad (3.5)$$

- ▶ **Homotopy transfer** (Thm 10.3.{1,7}, Loday-Valette **Algebraic Operads**): Given an L_∞ -algebra $(V_2, [-]_2)$ and a quasi-isomorphism $(V_1, \mathbf{s}_1) \rightarrow (V_2, \mathbf{s}_2)$ of dg-vector spaces, it can be extended to an L_∞ -morphism $(V_1, [-]_1 = \mathbf{s}_1 + ?) \rightarrow (V_2, [-]_2)$. The $[-]_1$ can be built “explicitly” provided the q-iso is presented by a zigzag of (dg-vector) homotopy retracts, using bar and co-bar constructions.
- ▶ Homotopy transfer interpretation of Brandt-Henneaux-Wilch'98:

$$(H^{<0,n}(\mathbf{s}|d_H), 0 \oplus f_{AB}^D \oplus f_{ABC}^D \oplus \dots) \xrightarrow{S_A, S_{AB}, S_{ABC}, \dots} (H^{<0,n}(d_H), \mathbf{s} \oplus (-, -))$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- Recall the hierarchy $F_{BV} \rightarrow F_{BRST} \rightarrow F \rightarrow M$ (topologically trivial M, F).
$$\begin{array}{ccccc} \text{antifields} & \text{ghosts} & \text{fields} \\ \downarrow & \downarrow & \downarrow \\ \text{---on-shell} & \text{---gauge} \\ \text{Koszul-Tate} & \text{Chevalley-Eilenberg} \\ \mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE} & & = \mathcal{Z}_{KT} + \mathcal{Z}_{CE}, \mathbf{s}_{KT}^2 = 0. \end{array}$$
- evl $\mathfrak{X}_\alpha(J^\infty F_{BRST}) \ni \beta = \beta^i(\partial/\partial\Phi^i) \leadsto \Phi_i^*\beta^i = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ no form
- Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{Z}_\beta \Omega = 0 \pmod{d_H}$:
$$H_{F_{BRST}}^*(\{[Q_{CE}, -]\}_\alpha) \cong H_{J^\infty F_{BV}, \Omega}^*(\{[Q, -]\}_\alpha) \cong H^{*,<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{*,<0,n}(\mathbf{s}|d_H)$$

- Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- Conserved currents: $d_H j = 0, j \in \Omega^{0,n-\sigma}(\mathcal{S}) = H_{\text{anti}\#-0}^{0,n-\sigma}(\mathbf{s}_{KT})$.
- Dickey bracket: $[j_1, j_2]_\sigma = \mathcal{Z}_{KT} j_2 \pmod{d_H}$ on $H_E^{0,n-\sigma}(d_H) \cong H_{\text{anti}\#-0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- Noether's first theorem (Lie algebra isomorphisms):

$$(H_{F_{BRST}}^0(\{[Q, -]\}), \{ -, - \}_\alpha) \cong (H^{-1,0}(\mathbf{s}|d_H), \{ -, - \}) \cong (H_E^{0,n-1}(d_H), \{ -, - \}_D)$$

- Tantalizing hint of homotopy transfer?

$$H_{F_{BRST}}^{0,0}(\{[Q_{CE}, -]\}_\alpha) \cong H^{-p,0}(\mathbf{s}|d_H) \cong H_E^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I*
CMP **174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
 $\xrightarrow{\text{on-shell}}$
 $\xrightarrow{\sim/\text{gauge}}$
- ▶ BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate Chevalley-Eilenberg $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0$.
- ▶ evf $\mathfrak{X}_{ev}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:
 $H_{\mathcal{E}_{BRST}}^*(([Q_{CE}, -]_{ev}) \cong H_{J^\infty F_{BV}, \Omega}^*([Q, -]_{ev}) \cong H^{*,<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{*,<0,n}(\mathbf{s}|d_H)$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* JMP **37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-\sigma}(\mathcal{E}) = H_{\text{anti}\#-0}^{0,n-\sigma}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{Q_{KT}} j_2 \pmod{d_H}$ on $H_E^{0,n-\sigma}(d_H) \cong H_{\text{anti}\#-0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{ev}) \cong (H^{-1,0}(\mathbf{s}|d_H), [-, -]) \cong (H_E^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\text{canon}}^{0,0}([Q_{CE}, -]_{ev}) \cong H^{-p,0}(\mathbf{s}|d_H) \cong H_E^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
 - ▶ BV differential decomposition $s = \underset{\substack{\rightarrow \text{on-shell} \\ \text{Koszul-Tate}}}{s_{KT}} + \underset{\substack{\sim/\text{gauge} \\ \text{Chevalley-Eilenberg}}}{s_{CE}} = \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, s^2_{KT} = 0$.
 - ▶ evf $\mathfrak{X}_{ev}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ loc.form
 - ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:
 $H_{\mathcal{E}_{BRST}}^*(([Q_{CE}, -]_{ev}) \cong H_{J^\infty F_{BV}, \Omega}^*([Q, -]_{ev}) \cong H^{*,<0,n}(s_{KT}|d_H) \cong H^{*,<0,n}(s|d_H)$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-\sigma}(\mathcal{S}) = H_{\text{anti}\#-0}^{0,n-\sigma}(s_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{Q_D} j_2 \pmod{d_H}$ on $H_E^{0,n-\sigma}(d_H) \cong H_{\text{anti}\#-0}^{0,n-1}(d_H|s_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{ev}) \cong (H^{-1,0}(s|d_H), [-, -]) \cong (H_E^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\text{canon}}^{0,0}([Q_{CE}, -]_{ev}) \cong H^{-p,0}(s|d_H) \cong H_E^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate $\xrightarrow{\text{on-shell}}$ Chevalley-Eilenberg $\xrightarrow{\sim/\text{gauge}}$ $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0$.
- ▶ evf $\mathfrak{X}_{ev}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:
 $H_{\mathcal{E}_{BRST}}^*(([Q_{CE}, -]_{ev}) \cong H_{J^\infty F_{BV}, \Omega}^*([Q, -]_{ev}) \cong H^{*,<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{*,<0,n}(\mathbf{s}|d_H)$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-\sigma}(\mathcal{E}) = H_{\text{anti}\#-0}^{0,n-\sigma}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{Q_D} j_2 \pmod{d_H}$ on $H_E^{0,n-\sigma}(d_H) \cong H_{\text{anti}\#-0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{ev}) \cong (H^{-1,0}(\mathbf{s}|d_H), [-, -]) \cong (H_E^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\text{canon}}^{0,0}([Q_{CE}, -]_{ev}) \cong H^{-p,0}(\mathbf{s}|d_H) \cong H_E^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT}^{\text{on-shell}} + \mathbf{s}_{CE}^{\sim/\text{gauge}}$ $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}$, $\mathbf{s}_{KT}^2 = 0$.
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:
 $H_{\mathcal{E}_{BRST}}^*(([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^*([Q, -]_{\text{ev}}) \cong H^{*,<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{*,<0,n}(\mathbf{s}|d_H)$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0$, $j \in \Omega^{0,n-\sigma}(\mathcal{S}) = H_{\text{anti}\#-0}^{0,n-\sigma}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_\sigma = \mathcal{L}_{Q_{CE}} j_2 \pmod{d_H}$ on $H_E^{0,n-\sigma}(d_H) \cong H_{\text{anti}\#-0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{\text{ev}}) \cong (H^{-1,0}(\mathbf{s}|d_H), [-, -]) \cong (H_E^{0,n-1}(d_H), [-, -]_\sigma)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\text{Cohom}}^{0,0}([Q_{CE}, -]_{\text{ev}}) \cong H^{-p,0}(\mathbf{s}|d_H) \cong H_E^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ →on-shell
Koszul-Tate ~/gauge
Chevalley-Eilenberg $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0.$
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:

$$H_{\mathcal{E}_{BRST}}^\bullet([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^\bullet([Q, -]_{\text{ev}}) \cong H^{*,<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{*,<0,n}(\mathbf{s}|d_H)$$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-\sigma}(\mathcal{E}) = H_{\text{anti}\#-0}^{0,n-\sigma}(\mathbf{s}_{KT}).$
- ▶ Dickey bracket: $[j_1, j_2]_\sigma = \mathcal{L}_{Q_{CE}} j_2 \pmod{d_H}$ on $H_E^{0,n-\sigma}(d_H) \cong H_{\text{anti}\#-0}^{0,n-1}(d_H|\mathbf{s}_{KT}).$
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^\bullet([Q, -]), [-, -]_\sigma) \cong (H^{-1,\sigma}(\mathbf{s}|d_H), [-, -]) \cong (H_E^{0,n-1}(d_H), [-, -]_\sigma)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\mathcal{E}_{BRST}}^{0,\sigma}([Q_{CE}, -]_{\text{ev}}) \cong H^{-\sigma,\sigma}(\mathbf{s}|d_H) \cong H_E^{0,n-\sigma}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate on-shell + Chevalley-Eilenberg ~gauge $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0$.
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:

$$H_{\mathcal{E}_{BRST}}^{\bullet}([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^{\bullet}([Q, -]_{\text{ev}}) \cong H^{<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{<0,n}(\mathbf{s}|d_H)$$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-p}(\mathcal{E}) = H_{\text{anti}\#_0}^{0,n-p}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{\xi_1} j_2 \pmod{d_H}$ on $H_{\mathcal{E}}^{0,n-p}(d_H) \cong H_{\text{anti}\#_0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{\text{ev}}) \cong (H^{-1,n}(\mathbf{s}|d_H), (-, -)) \cong (H_{\mathcal{E}}^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\mathcal{E}_{BRST}}^{-p}([Q_{CE}, -]_{\text{ev}}) \cong H^{-p,n}(\mathbf{s}|d_H) \cong H_{\mathcal{E}}^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate → on-shell + Chevalley-Eilenberg ~ gauge $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0$.
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{*,<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:

$$H_{\mathcal{E}_{BRST}}^{\bullet}([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^{\bullet}([Q, -]_{\text{ev}}) \cong H^{*,<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{*,<0,n}(\mathbf{s}|d_H)$$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-p}(\mathcal{E}) = H_{\text{anti}\#_0}^{0,n-p}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{\xi_1} j_2 \pmod{d_H}$ on $H_{\mathcal{E}}^{0,n-p}(d_H) \cong H_{\text{anti}\#_0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{\text{ev}}) \cong (H^{-1,n}(\mathbf{s}|d_H), (-, -)) \cong (H_{\mathcal{E}}^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\mathcal{E}_{BRST}}^{-p}([Q_{CE}, -]_{\text{ev}}) \cong H^{-p,n}(\mathbf{s}|d_H) \cong H_{\mathcal{E}}^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate on-shell + Chevalley-Eilenberg ~gauge $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0$.
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:

$$H_{\mathcal{E}_{BRST}}^{\bullet}([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^{\bullet}([Q, -]_{\text{ev}}) \cong H^{<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{<0,n}(\mathbf{s}|d_H)$$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-p}(\mathcal{E}) = H_{\text{anti}\#_0}^{0,n-p}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{\xi_1} j_2 \pmod{d_H}$ on $H_{\mathcal{E}}^{0,n-p}(d_H) \cong H_{\text{anti}\#_0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{\text{ev}}) \cong (H^{-1,n}(\mathbf{s}|d_H), (-, -)) \cong (H_{\mathcal{E}}^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\mathcal{E}_{BRST}}^{-p}([Q_{CE}, -]_{\text{ev}}) \cong H^{-p,n}(\mathbf{s}|d_H) \cong H_{\mathcal{E}}^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate on-shell + Chevalley-Eilenberg ~gauge $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}, \mathbf{s}_{KT}^2 = 0.$
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:

$$H_{\mathcal{E}_{BRST}}^{\bullet}([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^{\bullet}([Q, -]_{\text{ev}}) \cong H^{<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{<0,n}(\mathbf{s}|d_H)$$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0, j \in \Omega^{0,n-p}(\mathcal{E}) = H_{\text{anti}\#_0}^{0,n-p}(\mathbf{s}_{KT}).$
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{\xi_1} j_2 \pmod{d_H}$ on $H_{\mathcal{E}}^{0,n-p}(d_H) \cong H_{\text{anti}\#_0}^{0,n-1}(d_H|\mathbf{s}_{KT}).$
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{\text{ev}}) \cong (H^{-1,n}(\mathbf{s}|d_H), (-, -)) \cong (H_{\mathcal{E}}^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\mathcal{E}_{BRST}}^{-p}([Q_{CE}, -]_{\text{ev}}) \cong H^{-p,n}(\mathbf{s}|d_H) \cong H_{\mathcal{E}}^{0,n-p}(d_H)$$

Noether's Theorem

“symmetries” \simeq “conserved currents”

- ▶ Barnich, Brandt, Henneaux (1995) *Local BRST cohomology in the antifield formalism: I* **CMP 174** 57–92

- ▶ Recall the hierarchy $F_{BV} \xrightarrow{\text{antifields}} F_{BRST} \xrightarrow{\text{ghosts}} F \xrightarrow{\text{fields}} M$ (topologically trivial M, F).
BV differential decomposition $\mathbf{s} = \mathbf{s}_{KT} + \mathbf{s}_{CE}$ Koszul-Tate Chevalley-Eilenberg $= \mathcal{L}_{Q_{KT}} + \mathcal{L}_{Q_{CE}}$, $\mathbf{s}_{KT}^2 = 0$.
- ▶ **evf** $\mathfrak{X}_{\text{ev}}(J^\infty F_{BRST}) \ni \beta = \beta^I(\partial/\partial\Phi^I) \rightsquigarrow \Phi_I^* \beta^I = b \in \Omega^{<0,n}(J^\infty F_{BV})$ loc.form
- ▶ Symmetry $[Q_{KT}, \beta] = 0$, gauge symmetry $\beta = [Q_{KT}, \gamma]$, among $\mathcal{L}_\beta \Omega = 0 \pmod{d_H}$:

$$H_{\mathcal{E}_{BRST}}^{\bullet}([Q_{CE}, -]_{\text{ev}}) \cong H_{J^\infty F_{BV}, \Omega}^{\bullet}([Q, -]_{\text{ev}}) \cong H^{<0,n}(\mathbf{s}_{KT}|d_H) \cong H^{<0,n}(\mathbf{s}|d_H)$$

- ▶ Barnich, Henneaux (1996) *Isomorphisms between the Batalin-Vilkovisky antibracket and the Poisson bracket* **JMP 37** 5273–5296

- ▶ Conserved currents: $d_H j = 0$, $j \in \Omega^{0,n-p}(\mathcal{E}) = H_{\text{anti}\# = 0}^{0,n-p}(\mathbf{s}_{KT})$.
- ▶ Dickey bracket: $[j_1, j_2]_D = \mathcal{L}_{\xi_1} j_2 \pmod{d_H}$ on $H_{\mathcal{E}}^{0,n-p}(d_H) \cong H_{\text{anti}\# = 0}^{0,n-1}(d_H|\mathbf{s}_{KT})$.
- ▶ Noether's first theorem (Lie algebra isomorphisms):

$$(H_{\mathcal{E}_{BRST}}^0([Q, -]), [-, -]_{\text{ev}}) \cong (H^{-1,n}(\mathbf{s}|d_H), (-, -)) \cong (H_{\mathcal{E}}^{0,n-1}(d_H), [-, -]_D)$$

- ▶ Tantalizing hint of homotopy transfer?

$$H_{\mathcal{E}_{BRST}}^{-p}([Q_{CE}, -]_{\text{ev}}) \cong H^{-p,n}(\mathbf{s}|d_H) \cong H_{\mathcal{E}}^{0,n-p}(d_H)$$

(New?) Local Antibracket

- The **antibracket** $\left(\int b(\Phi, \Phi^*), \int c(\Phi, \Phi^*) \right)$ is traditionally defined on **local functionals** or $H^{\bullet,n}(d_H)$ classes $[b], [c]$.
- **Local antibracket:** lift to local forms $\Omega^{\bullet,\bullet}$. *Barnich-Henneaux'96* tried

$$(b, c)_{\text{loc}} = \mathcal{L}_\beta c \quad \text{or} \quad -\mathcal{L}_\gamma b \quad \text{or} \quad \iota_\beta \iota_\gamma \Omega$$

where $d_V b = \iota_\beta \Omega - d_H \theta_b$, $d_V c = \iota_\gamma \Omega - d_H \theta_c$ and $\Omega = d_V \Phi_i^* d_V \Phi^I$ is the local antibracket shifted symplectic form (density).

- All these choices satisfy at least one of anti-symmetry, Jacobi, (Leibniz) identities **only up to d_H** . By the transfer theorem, there **might only be** an L_∞ -transfer to $\Omega^{\bullet,\bullet}$ (cf. *Barnich-Fulp-Lada-Stasheff'98*).
- **Theorem:** (via Prop 17.2.3 *Delgado (PhD, Bonn 2017)*; via Eq (2.100) *Deligne-Freed'99*)
 $(\Omega^{\bullet,\bullet}[1, \text{odd}], d_H + \tilde{s}, (-, -)_{\text{loc}})$ is a **dg-Lie algebra** on the nose, with (for odd $b, c \in \Omega^{\bullet,\bullet}$):

$$(b, c)_{\text{loc}} = \begin{cases} \mathcal{L}_\beta c - \mathcal{L}_\gamma b - \iota_\beta \iota_\gamma \Omega & \text{if } b, c \in \Omega^{\bullet,n}, \\ \mathcal{L}_\beta c & \text{if } b \in \Omega^{\bullet,n}, c \in \Omega^{\bullet,<n}, \\ 0 & \text{if } b, c \in \Omega^{\bullet,<n}. \end{cases}$$

Proof. For Jacobi, use $\Omega = d_V(\Phi_i^* d_V \Phi^I)$ and $(b, c)_{\text{loc}} \rightsquigarrow [\beta, \gamma]$ (*Barnich-Henneaux'96*), while $\tilde{s}(-) = (L + J + \dots, -)_{\text{loc}}$, with $d_H J = (L, L)_{\text{loc}}$.

(New?) Local Antibracket

- The **antibracket** $\left(\int b(\Phi, \Phi^*), \int c(\Phi, \Phi^*) \right)$ is traditionally defined on **local functionals** or $H^{\bullet,n}(d_H)$ classes $[b], [c]$.
- **Local antibracket:** lift to local forms $\Omega^{\bullet,\bullet}$. *Barnich-Henneaux'96* tried

$$(b, c)_{\text{loc}} = \mathcal{L}_\beta c \quad \text{or} \quad -\mathcal{L}_\gamma b \quad \text{or} \quad \iota_\beta \iota_\gamma \Omega$$

where $d_V b = \iota_\beta \Omega - d_H \theta_b$, $d_V c = \iota_\gamma \Omega - d_H \theta_c$ and $\Omega = d_V \Phi_i^* d_V \Phi^I$ is the **local antibracket shifted symplectic form (density)**.

- All these choices satisfy at least one of anti-symmetry, Jacobi, (Leibniz) identities **only up to d_H** . By the transfer theorem, there **might only be** an L_∞ -transfer to $\Omega^{\bullet,\bullet}$ (cf. *Barnich-Fulp-Lada-Stasheff'98*).
- **Theorem:** (via Prop 17.2.3 *Delgado (PhD, Bonn 2017)*; via Eq (2.100) *Deligne-Freed'99*)
 $(\Omega^{\bullet,\bullet}[1, \text{odd}], d_H + \tilde{s}, (-, -)_{\text{loc}})$ is a **dg-Lie algebra** on the nose, with (for odd $b, c \in \Omega^{\bullet,\bullet}$):

$$(b, c)_{\text{loc}} = \begin{cases} \mathcal{L}_\beta c - \mathcal{L}_\gamma b - \iota_\beta \iota_\gamma \Omega & \text{if } b, c \in \Omega^{\bullet,n}, \\ \mathcal{L}_\beta c & \text{if } b \in \Omega^{\bullet,n}, c \in \Omega^{\bullet,<n}, \\ 0 & \text{if } b, c \in \Omega^{\bullet,<n}. \end{cases}$$

Proof. For Jacobi, use $\Omega = d_V(\Phi_i^* d_V \Phi^I)$ and $(b, c)_{\text{loc}} \rightsquigarrow [\beta, \gamma]$ (*Barnich-Henneaux'96*), while $\tilde{s}(-) = (L + J + \dots, -)_{\text{loc}}$, with $d_H J = (L, L)_{\text{loc}}$.

(New?) Local Antibracket

- The **antibracket** $\left(\int b(\Phi, \Phi^*), \int c(\Phi, \Phi^*) \right)$ is traditionally defined on **local functionals** or $H^{\bullet,n}(d_H)$ classes $[b], [c]$.
- **Local antibracket:** lift to local forms $\Omega^{\bullet,\bullet}$. *Barnich-Henneaux'96* tried

$$(b, c)_{\text{loc}} = \mathcal{L}_\beta c \quad \text{or} \quad -\mathcal{L}_\gamma b \quad \text{or} \quad \iota_\beta \iota_\gamma \Omega$$

where $d_V b = \iota_\beta \Omega - d_H \theta_b$, $d_V c = \iota_\gamma \Omega - d_H \theta_c$ and $\Omega = d_V \Phi_i^* d_V \Phi^I$ is the **local antibracket shifted symplectic form** (density).

- All these choices satisfy at least one of anti-symmetry, Jacobi, (Leibniz) identities **only up to** d_H . By the transfer theorem, there **might only be** an L_∞ -**transfer** to $\Omega^{\bullet,\bullet}$ (cf. *Barnich-Fulp-Lada-Stasheff'98*).
- **Theorem:** (via Prop 17.2.3 *Delgado (PhD, Bonn 2017)*; via Eq (2.100) *Deligne-Freed'99*) $(\Omega^{\bullet,\bullet}[1, \text{odd}], d_H + \tilde{s}, (-, -)_{\text{loc}})$ is a **dg-Lie algebra** on the nose, with (for odd $b, c \in \Omega^{\bullet,\bullet}$):

$$(b, c)_{\text{loc}} = \begin{cases} \mathcal{L}_\beta c - \mathcal{L}_\gamma b - \iota_\beta \iota_\gamma \Omega & \text{if } b, c \in \Omega^{\bullet,n}, \\ \mathcal{L}_\beta c & \text{if } b \in \Omega^{\bullet,n}, c \in \Omega^{\bullet,<n}, \\ 0 & \text{if } b, c \in \Omega^{\bullet,<n}. \end{cases}$$

Proof. For Jacobi, use $\Omega = d_V(\Phi_i^* d_V \Phi^I)$ and $(b, c)_{\text{loc}} \rightsquigarrow [\beta, \gamma]$ (*Barnich-Henneaux'96*), while $\tilde{s}(-) = (L + J + \dots, -)_{\text{loc}}$, with $d_H J = (L, L)_{\text{loc}}$.

(New?) Local Antibracket

- The **antibracket** $\left(\int b(\Phi, \Phi^*), \int c(\Phi, \Phi^*) \right)$ is traditionally defined on **local functionals** or $H^{\bullet,n}(d_H)$ classes $[b], [c]$.
- **Local antibracket:** lift to local forms $\Omega^{\bullet,\bullet}$. *Barnich-Henneaux'96* tried

$$(b, c)_{\text{loc}} = \mathcal{L}_\beta c \quad \text{or} \quad -\mathcal{L}_\gamma b \quad \text{or} \quad \iota_\beta \iota_\gamma \Omega$$

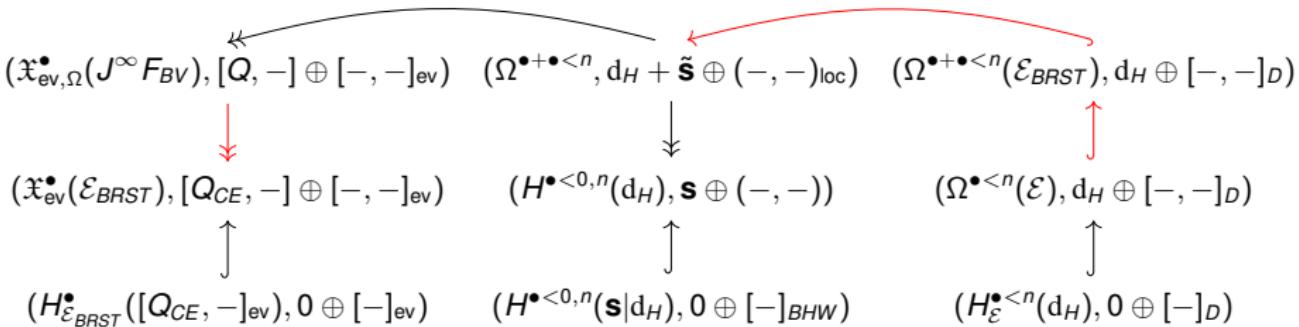
where $d_V b = \iota_\beta \Omega - d_H \theta_b$, $d_V c = \iota_\gamma \Omega - d_H \theta_c$ and $\Omega = d_V \Phi_i^* d_V \Phi^I$ is the **local antibracket shifted symplectic form** (density).

- All these choices satisfy at least one of anti-symmetry, Jacobi, (Leibniz) identities **only up to** d_H . By the transfer theorem, there **might only be** an L_∞ -transfer to $\Omega^{\bullet,\bullet}$ (cf. *Barnich-Fulp-Lada-Stasheff'98*).
- **Theorem:** (via Prop 17.2.3 *Delgado (PhD, Bonn 2017)*; via Eq (2.100) *Deligne-Freed'99*) $(\Omega^{\bullet,\bullet}[1, \text{odd}], d_H + \tilde{s}, (-, -)_{\text{loc}})$ is a **dg-Lie algebra** on the nose, with (for odd $b, c \in \Omega^{\bullet,\bullet}$):

$$(b, c)_{\text{loc}} = \begin{cases} \mathcal{L}_\beta c - \mathcal{L}_\gamma b - \iota_\beta \iota_\gamma \Omega & \text{if } b, c \in \Omega^{\bullet,n}, \\ \mathcal{L}_\beta c & \text{if } b \in \Omega^{\bullet,n}, c \in \Omega^{\bullet,<n}, \\ 0 & \text{if } b, c \in \Omega^{\bullet,<n}. \end{cases}$$

Proof. For Jacobi, use $\Omega = d_V(\Phi_i^* d_V \Phi^I)$ and $(b, c)_{\text{loc}} \rightsquigarrow [\beta, \gamma]$ (*Barnich-Henneaux'96*), while $\tilde{s}(-) = (L + J + \dots, -)_{\text{loc}}$, with $d_H J = (L, L)_{\text{loc}}$.

L_∞ -zigzags (WIP)

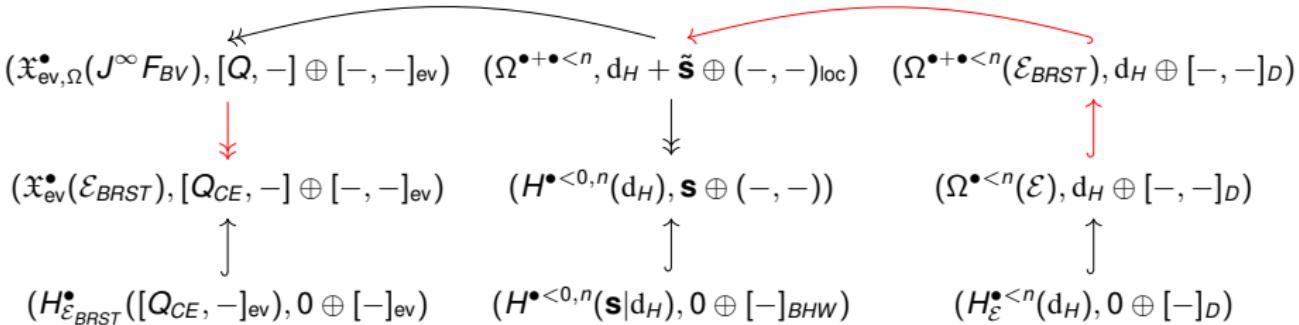


- ▶ Arrows indicate **dg-Lie morphisms**, or inclusion of **dg-vector cocycles**.
All arrows should be **dg-vector quasi-isomorphisms**. (**Work In Progress**)
- ▶ (conj.) **Extended Noether's theorem**:

$$(H_{E_{BRST}}^bullet([Q_{CE}, -]_{ev}), 0 \oplus [-]_{ev}) \xrightleftharpoons{L_\infty \text{ equiv.}} (H_E^{bullet < n}(d_H), 0 \oplus [-]_D)$$

- ▶ Topologically non-trivial $E \rightarrow M$: $H_E^*(d_H) \rightsquigarrow H_E^*(d_H)/H_E^*(d)$.
~~~ Central  $L_\infty$ -extension?
- ▶ Bonus:  $(\Omega^*(E), d_H, \wedge)$  is a **dgca**.  
~~~ Homotopy Transfer of  $C_\infty$ -algebra structure?

L_∞ -zigzags (WIP)

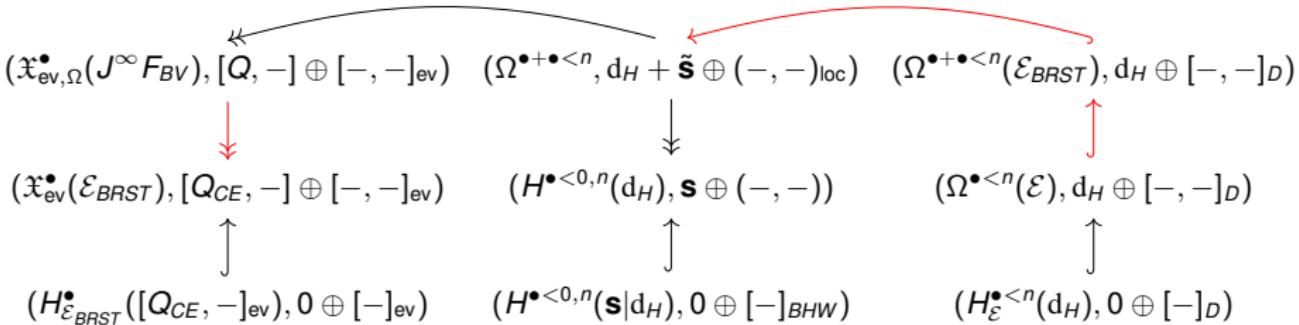


- ▶ Arrows indicate **dg-Lie morphisms**, or inclusion of **dg-vector cocycles**.
All arrows should be **dg-vector quasi-isomorphisms**. (**Work In Progress**)
- ▶ (conj.) **Extended Noether's theorem:**

$$(H_{E_{BRST}}^bullet([Q_{CE}, -]_{ev}), 0 \oplus [-]_{ev}) \xleftarrow{L_\infty \text{ equiv.}} (H_E^{bullet < n}(d_H), 0 \oplus [-]_D)$$

- ▶ Topologically non-trivial $E \rightarrow M$: $H_E^*(d_H) \rightsquigarrow H_E^*(d_H)/H_E^*(d)$.
~~~ Central  $L_\infty$ -extension?
- ▶ Bonus:  $(\Omega^*(E), d_H, \wedge)$  is a **dgca**.  
~~~ Homotopy Transfer of  $C_\infty$ -algebra structure?

L_∞ -zigzags (WIP)

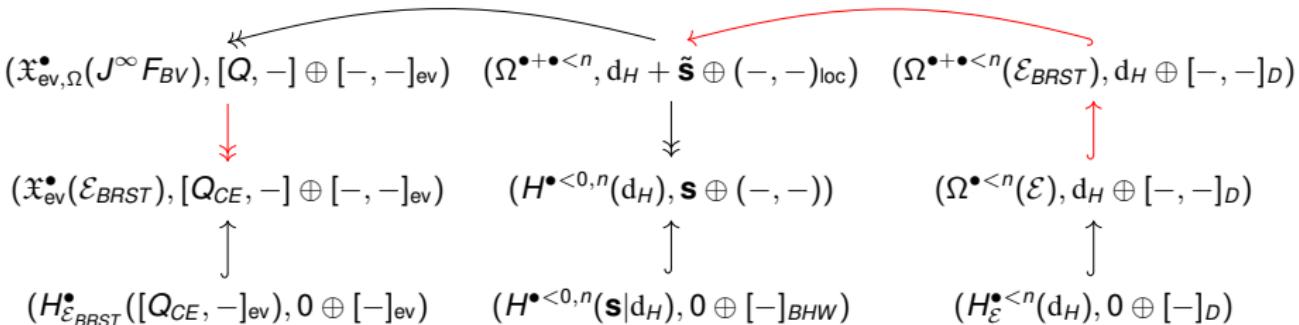


- ▶ Arrows indicate **dg-Lie morphisms**, or inclusion of **dg-vector cocycles**.
All arrows should be **dg-vector quasi-isomorphisms**. (**Work In Progress**)
- ▶ (conj.) **Extended Noether's theorem**:

$$(H_{\mathcal{E}_{BRST}}^\bullet([Q_{CE}, -]_{ev}), 0 \oplus [-]_{ev}) \xleftarrow{L_\infty \text{ equiv.}} (H_{\mathcal{E}}^\bullet(d_H), 0 \oplus [-]_D)$$

- ▶ **Topologically non-trivial** $\mathcal{E} \rightarrow M$: $H_{\mathcal{E}}^\bullet(d_H) \rightsquigarrow H_{\mathcal{E}}^\bullet(d_H)/H_{\mathcal{E}}^\bullet(d)$.
~~~ Central  $L_\infty$ -extension?
- ▶ **Bonus**:  $(\Omega^\bullet(\mathcal{E}), d_H, \wedge)$  is a **dgca**.  
~~~ Homotopy Transfer of  $C_\infty$ -algebra structure?

L_∞ -zigzags (WIP)



- ▶ Arrows indicate **dg-Lie morphisms**, or inclusion of **dg-vector cocycles**.
All arrows should be **dg-vector quasi-isomorphisms**. (**Work In Progress**)
- ▶ (conj.) **Extended Noether's theorem**:

$$(H_{\mathcal{E}_{BRST}}^\bullet([Q_{CE}, -]_{\text{ev}}), 0 \oplus [-]_{\text{ev}}) \xrightleftharpoons{L_\infty \text{ equiv.}} (H_{\mathcal{E}}^\bullet(d_H), 0 \oplus [-]_D)$$

- ▶ **Topologically non-trivial** $\mathcal{E} \rightarrow M$: $H_{\mathcal{E}}^\bullet(d_H) \rightsquigarrow H_{\mathcal{E}}^\bullet(d_H)/H_{\mathcal{E}}^\bullet(d)$.
 ↵ Central L_∞ -extension?
- ▶ **Bonus:** $(\Omega^\bullet(\mathcal{E}), d_H, \wedge)$ is a **dgca**.
 ↵ Homotopy Transfer of C_∞ -algebra structure?

Discussion

- ▶ L_∞ Homotopy Transfer interpretation of constant ghost/antifield extended BV (*Brandt-Henneaux-Wilch'98*).
Cf. talk in Prague Mathematical Physics Seminar by *Hiroaki Matsunaga* (05.2021).
- ▶ Successful dg-Lie local lift $(-, -)_{\text{loc}}$ of antibracket.
- ▶ Geometric interpretation of $H^{\bullet < 0, n}(\mathbf{s}|d_H)$ via higher symmetries and conserved currents.
- ▶ L_∞ -extension of Noether's theorem.

Thank you for your attention!

Discussion

- ▶ L_∞ Homotopy Transfer interpretation of constant ghost/antifield extended BV (*Brandt-Henneaux-Wilch'98*).
Cf. talk in Prague Mathematical Physics Seminar by *Hiroaki Matsunaga* (05.2021).
- ▶ Successful dg-Lie local lift $(-, -)_{\text{loc}}$ of antibracket.
- ▶ Geometric interpretation of $H^{\bullet < 0, n}(\mathbf{s}|d_H)$ via higher symmetries and conserved currents.
- ▶ L_∞ -extension of Noether's theorem.

Thank you for your attention!

Discussion

- ▶ L_∞ Homotopy Transfer interpretation of constant ghost/antifield extended BV (*Brandt-Henneaux-Wilch'98*).
Cf. talk in Prague Mathematical Physics Seminar by *Hiroaki Matsunaga* (05.2021).
- ▶ Successful dg-Lie local lift $(-, -)_{\text{loc}}$ of antibracket.
- ▶ Geometric interpretation of $H^{\bullet < 0, n}(\mathbf{s}|d_H)$ via higher symmetries and conserved currents.
- ▶ L_∞ -extension of Noether's theorem.

Thank you for your attention!

Discussion

- ▶ L_∞ Homotopy Transfer interpretation of constant ghost/antifield extended BV (*Brandt-Henneaux-Wilch'98*).
Cf. talk in Prague Mathematical Physics Seminar by *Hiroaki Matsunaga* (05.2021).
- ▶ Successful dg-Lie local lift $(-, -)_{\text{loc}}$ of antibracket.
- ▶ Geometric interpretation of $H^{\bullet < 0, n}(\mathbf{s}|d_H)$ via higher symmetries and conserved currents.
- ▶ L_∞ -extension of Noether's theorem.

Thank you for your attention!

References

-  G. Barnich, R. Fulp, T. Lada, and J. Stasheff, "The sh Lie structure of Poisson brackets in field theory," *Communications in Mathematical Physics* **191** (1998) 585–601, arXiv:hep-th/9702176.
-  G. Barnich, F. Brandt, and M. Henneaux, "Local BRST cohomology in gauge theories," *Physics Reports* **338** no. 5, (Nov., 2000) 439–569, arXiv:hep-th/0002245.
-  G. Barnich and M. Henneaux, "Isomorphisms between the Batalin–Vilkovisky antibracket and the poisson bracket," *Journal of Mathematical Physics* **37** (1996) 5273–5296, arXiv:hep-th/9601124.
-  F. Brandt, M. Henneaux, and A. Wilch, "Extended antifield formalism," *Nuclear Physics B* **510** (1998) 640–656, arXiv:hep-th/9705007.
-  N. L. Delgado, *Lagrangian field theories: ind/pro-approach and L_∞ -algebra of local observables*.
PhD thesis, MPI Bonn, 2017.
arXiv:1805.10317.
-  J.-L. Loday and B. Vallette, *Algebraic Operads*, vol. 346 of *Grundlehren der mathematischen Wissenschaften*.
Springer, Berlin, 2012.