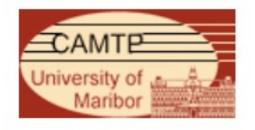
Erwin Schrödinger International Institute for Mathematics and Physics The Landscape versus the Swampland July 1- August 9,2024

# F-theory: Landscape of Particle Physics Models

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Univerza *v Ljubljani* Fakulteta za *matematiko in fiziko* 





#### I. The program:

Globally consistent F-theory compactifications with the gauge symmetry and matter spectrum of the Standard Model Key building blocks: gauge symmetry; matter spectrum; global conditions Upenn-centric

- II. Status of constructions: via toric geometry techniques
  - a) First globally consistent three-family Standard Model [M.C., Klevers, Peña, Oehlmann, Reuter 1503.02068; M.C., Lin, Liu, Oehlmann 1807.01320]
  - b) Landscape of three-family Standard Models

[M.C., Jim Halverson, Ling Lin, Muyang Liu, Jiahua Tian, "A Quadrillion Standard Models from F-theory" 1903.00009, PRL]

#### III. Further analysis:

Time permitting

- a) Constraints on moduli stabilization scenarios [M.C., Cody Long, J. Halverson, L. Lin 2004.00630]
- b) Toward charged vector matter pairs

[Martin Bies, M.C., Ron. Donagi, L. Lin, M. Liu, Fabian Rühle 2007.0009] [Bies, M.C. Donagi, Liu 2102.10115, 2104.08297] [Bies, M.C. Donagi, Marielle Ong 2205.00008, 2307.02535]

#### **Outline:**

I. F-theory: Geometric Approach

Appearance of non-Abelian gauge symmetry, matter and Yukawa couplings Appearance of Abelian continuous and discrete symmetries, global constraints

II. Construction of Particle physics Models

Building blocks for consistent models via toric techniques

Landscape of three-family Standard Models

#### III. Further Analysis:

Time permitting or Standard

Constraints on moduli stabilization scenarios for Standard Model constructions; toward vector-pair matter calculation

IV. Outlook: work in progress & open issues

#### Type II String

- back-reacted D-branes
- regions with large g<sub>s</sub> on non-CY space

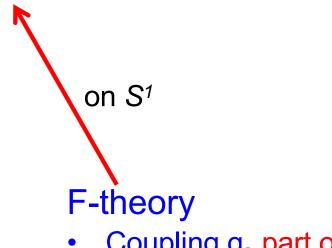
#### F-theory

- Coupling g<sub>s</sub> part of geometry (12dim)
- Torus fibered
   Calabi-Yau manifold

#### Type II String

- back-reacted D-branes
- regions with large g<sub>s</sub> on non-CY space

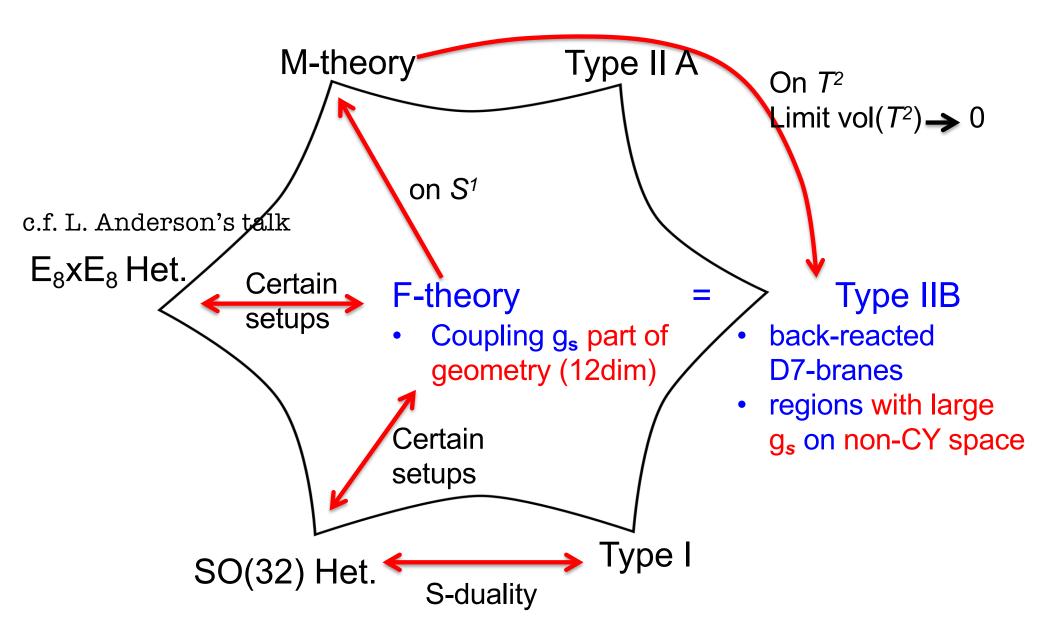
M-theory (11dim SG)



Coupling g<sub>s</sub> part of geometry (12dim)

#### Type IIB

- back-reacted D7-branes
- regions with large
   g<sub>s</sub> on non-CY space



## I. F-theory basic ingredients

Type IIB string perspective

## F-theory compactification

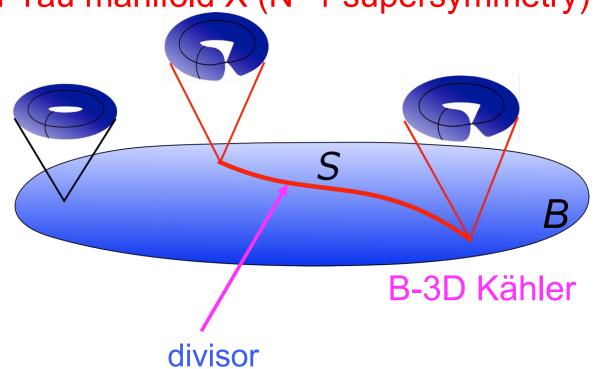
[Vafa'96], [Morrison, Vafa'96]

review [Weigand 1806.01854] Singular torus fibered Calabi-Yau manifold X (N=1 supersymmetry)

To B add torus: Modular parameter of torus (elliptic curve)

$$\tau \equiv C_0 + ig_s^{-1}$$

(SL(2,Z) of Type IIB string)



#### Weierstrass normal form for torus (elliptic) fibration of X

$$y^2 = x^3 + fxz^4 + gz^6$$

[z:x:y] - homogeneous coordinates on  $P^2(1,2,3)$   $(x,y,z) \simeq (\lambda^2 x, \lambda^3 y, \lambda z)$ weighted projective space

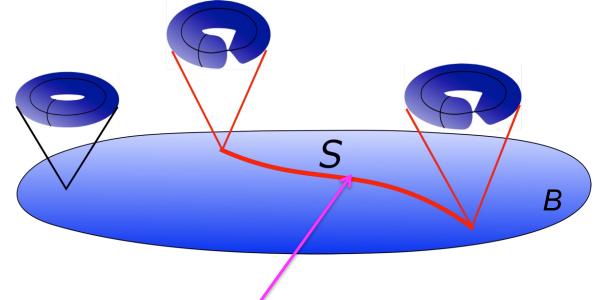
f, g – sections of  $\overline{\mathcal{K}}_B^6$  and  $\overline{\mathcal{K}}_B^4$  on B  $\overline{\mathcal{K}}_B$  -anti-canonical bundle on B

## F-theory compactification

Singular elliptically fibered Calabi-Yau manifold X

Modular parameter of two-torus (elliptic curve)

$$\tau \equiv C_0 + ig_s^{-1}$$



divisor- singular elliptic-fibration  $g_s \rightarrow \infty$  location of (p,q) 7-branes

Non-Abelian gauge symmetry (co-dim 1) – ADE singularities

## I.a Non-Abelian Gauge Symmetry

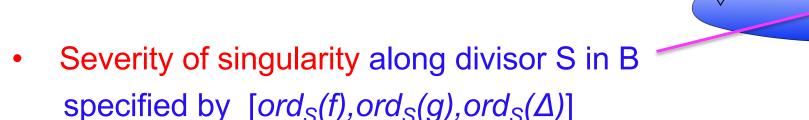
Standard Model has SU(2)<sub>L</sub> x SU(3)<sub>C</sub>

## Non-Abelian Gauge Symmetry

[Kodaira],[Tate], [Vafa], [Morrison, Vafa],...[Esole, Yau], [Hayashi, Lawrie, Schäfer-Nameki], [Morrison], ...

Weierstrass normal form for elliptic fibration of X

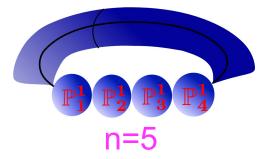
$$y^2 = x^3 + fxz^4 + gz^6$$











Cartan gauge bosons: supported by (1,1) form  $\omega_i \leftrightarrow \mathbb{P}^1_i$  on resolved X

(via M-theory Kaluza-Klein reduction of  $C_3$  potential  $C_3 \supset A^i \omega_i$ )

Non-Abelian gauge bosons: light M2-brane excitations on P<sup>1</sup>s

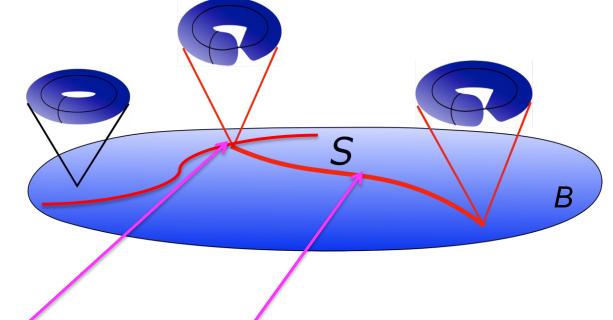
[Witten]

## F-theory compactification

Singular elliptically fibered Calabi-Yau manifold X

Modular parameter of two-torus (elliptic curve)

$$\tau \equiv C_0 + ig_s^{-1}$$



Matter (co-dim 2)

divisor- singular elliptic-fibration  $g_s \rightarrow \infty$  location of (p,q) 7-branes

Non-Abelian gauge symmetry (co-dim 1) – ADE singularities Abelian symmetries different

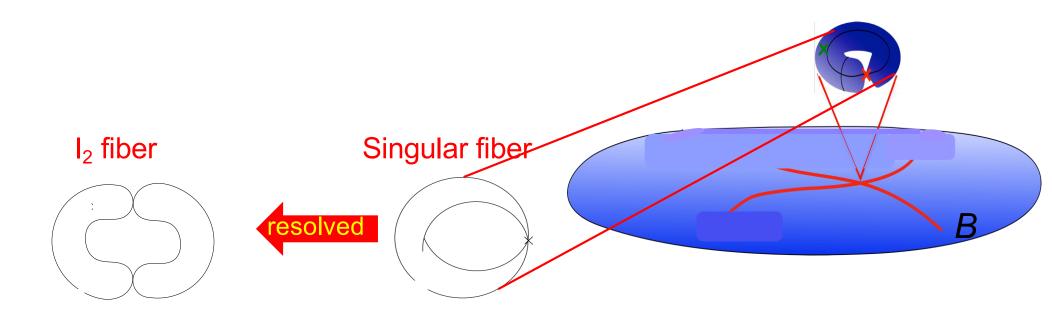
 $\rightarrow$  More later

#### Matter

Standard Model  $SU(2)_L \times SU(3)_C$  has quarks Q ~ ( 2 , 3 )

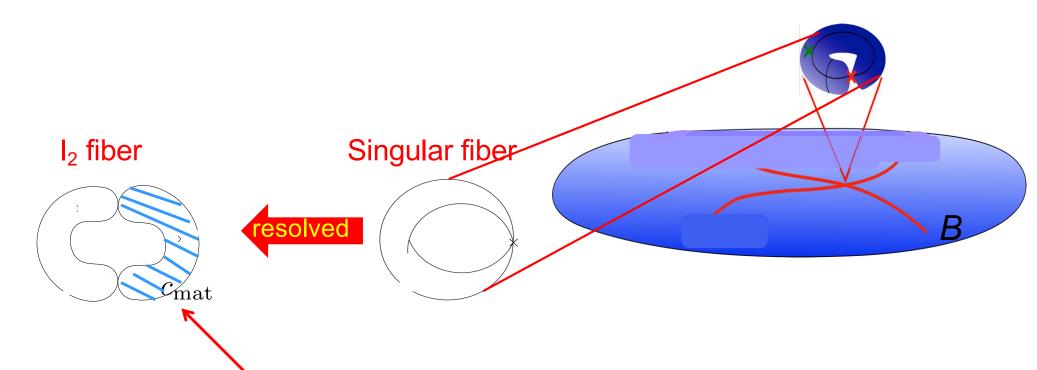
#### Matter

#### Singularity at codimension-two in B:



#### Matter

#### Singularity at codimension-two in B:



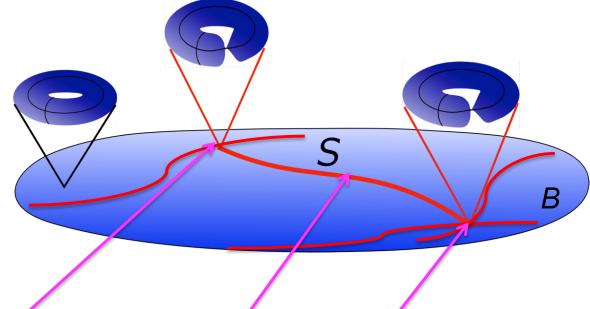
w/isolated (M2-matter) curve wrapping  $\mathbb{P}^1 \to \text{charged matter}$  (determine rep. via intersection theory)

## F-theory compactification

Singular elliptically fibered Calabi-Yau manifold X

Modular parameter of two-torus (elliptic curve)

$$\tau \equiv C_0 + ig_s^{-1}$$



Matter (co-dim 2)

Chirality: Should add  $G_4$ -flux  $\rightarrow$  More later

c.f., [Lin, Mayrhofer, Till, Weigand, 1508.00162] [M.C., Grassi, Klevers, Piragua, 1306.3987]

Yukawa couplings (co-dim 3) → No time

Divisor- singular elliptic-fibration  $g_s \rightarrow \infty$  location of (p,q) 7-branes

Non-Abelian gauge symmetry (co-dim 1) – ADE singularities Abelian symmetries different

 $\rightarrow$  More later

I.b Abelian Gauge Symmetry - U(1)

Standard Model has SU(2)<sub>L</sub> x SU(3)<sub>C</sub> x U(1)<sub>Y</sub>

## I.b Abelian Gauge Symmetry - U(1)

**Different**: (1,1) forms  $\omega_m$ , supporting U(1) gauge bosons, isolated & associated with  $I_1$ -fibers, only

[Morrison, Vafa'96]

(1,1) - form  $\omega_m$  rational section of elliptic fibration

#### Abelian Gauge Symmetry & Mordell-Weil Group

rational sections of elliptic fibr. 
rational points of elliptic curve

#### Abelian Gauge Symmetry & Mordell-Weil Group

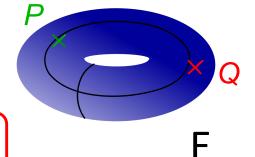
rational sections of elliptic fibr. - rational points of elliptic curve

#### Rational point Q on elliptic curve E with zero point P

• is solution  $(x_Q, y_Q, z_Q)$  in field K of Weierstrass form

$$y^2 = x^3 + fxz^4 + gz^6$$

Rational points form group (addition) on E

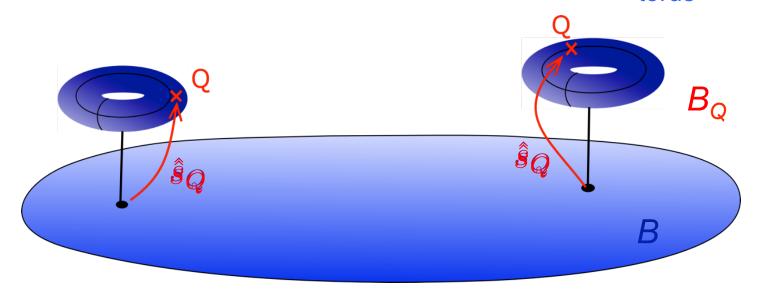




Mordell-Weil group of rational points

## U(1)'s-Abelian Symmetry & Mordell-Weil Group

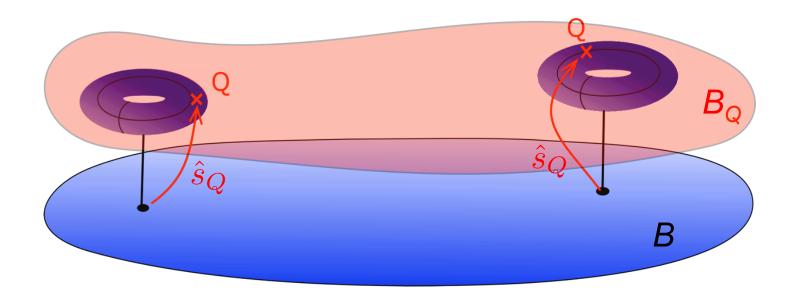
Point Q E induces a rational section  $\hat{s}_Q: B \to X$  of elliptic fibration torus



 $\hat{S}_Q$  gives rise to a second copy of B in X: new divisor  $B_Q$  in X

#### U(1)'s-Abelian Symmetry & Mordell-Weil Group

Point Q E induces a rational section  $\hat{s}_Q: B \to X$  of elliptic fibration



 $\hat{S}_Q$  gives rise to a second copy of B in X:

new divisor  $B_Q$  in X Construct (1,1) form  $\omega_m$  from  $B_Q$ 

Shioda map of  $\hat{s}_Q$ , complementary to  $\mathcal{B}_{\mathcal{P}}$ - zero section &  $\mathcal{E}_{\mathsf{i}}$  - Cartan divisors:  $\sigma\left(\hat{s}_{Q}\right) = B_{Q} - B_{P} - \sum_{i} l_{i} E_{i} + \cdots$ [M.C., Lin, 1706.08521]

Implications for global constraints on gauge symmetries  $\rightarrow$  shortly

Earlier work: [Grimm, Weigand 1006.0226]...[Grassi, Perduca 1201.0930] [M.C., Grimm, Klevers 1210.6034]...

## Explicit Examples: (n+1)-rational sections – U(1)<sup>n</sup> [Deligne]

[via line bundle constr. on elliptic curve E- CY in (blow-up) of  $W\mathbb{P}^m$ ]

```
n=0: with P - generic CY in \mathbb{P}^2(1,2,3) (Tate form)
```

*n*=1: with *P*, *Q* - generic CY in  $Bl_1\mathbb{P}^2(1,1,2)$  [Morrison,Park 1208.2695]...

## Explicit Examples: (n+1)-rational sections – U(1)<sup>n</sup>

[Morrison, Park 1208.2695]...

[Borchmann, Mayerhofer, Palti, Weigand

with P - generic CY in  $\mathbb{P}^2(1,2,3)$  (Tate form)

with P, Q, R - specific example: generic CY in  $dP_2$ 

with P, Q - generic CY in  $Bl_1\mathbb{P}^2(1,1,2)$ 

n=0:

```
[M.C.,Klevers,Piragua 1303.6970,1307.6425]

[M.C.,Grassi,Klevers,Piragua 1306.0236]

generalization to nongeneric cubic in \mathbb{P}^2[u:v:w] \rightarrow \text{collisions of P,Q,R} \rightarrow \text{gauge enhancement \& higher index representations....}

[M.C.,Klevers,Piragua,Taylor 1507.05954]

n=3: \text{ with } P, Q, R, S - \text{CICY in Bl}_3\mathbb{P}^3
n=4: \text{ determinantal variety in } \mathbb{P}^4 [M.C.,Klevers,Piragua,Song 1310.0463] higher n, not clear...

No time \rightarrow \text{c.f.,[M.C. Lin,1809.00012]} TASI review
```

[M.C. and Ling Lin 1706.08521]

c.f., review [M.C. Lin,1809.00012]

## I.b U(1) & Non-Abelian Gauge symmetry

## Shioda map of section $\hat{s}_Q$ more involved than $\mathcal{B}_Q$ :

a map onto divisor complementary to  $\mathcal{B}_{\mathcal{P}}$  divisor of zero section  $\hat{s}_P$ 

&  $\mathcal{L}_i$  – resolution (Cartan) divisors of non-Abelian gauge symmetry

$$\sigma\left(\hat{s}_{Q}\right) = B_{Q} - B_{P} - \sum_{i} l_{i} E_{i} + \cdots$$

Ensures proper F-theory interpretation of U(1) (via M-theory/F-theory duality)

$$l_i = C_{ij}^{-1}(B_Q - B_P) \cdot \mathbb{P}^1_j$$
 - fractional # always an integer  $\kappa$  s.t.  $\forall i$ :  $\kappa$   $l_i \in \mathbb{Z}$ 

Cartan matrix Fiber of divisor E<sub>j</sub>

R'S SPIN SYSTEM ANALOGY mics from the outer and inner horis

Curir\_regards a rota**tpositacke**holemaperatolereysamdwtheoimmericherizahourio, rhe Employing (a)  $q_{\mathfrak{g}(1)} = \frac{n}{n}$ ,  $n \in \mathbb{Z}$  & Leading the outer and inner horizon. The outer horizon is takener more and angular momentum J are common to both superature and the inner horizon to have a negative temperature,  $\mathbf{v} = \mathbf{v} + \mathbf{v} +$ 

ar momentum J are communa Guriystelasims  $^1[1]$ .  $\sup_{\mathbf{u}\in \mathbf{F}}(\mathbf{w})_{\mathbf{m}}=^1[\mathbf{e}]^{2\pi iq(\mathbf{w})}\otimes (\mathbf{e}^{-2\pi i\,l_i\,\mathbf{w}_i}\times \mathbb{1})]$   $\mathbf{v}$  I. CURIR'S SPIN SYSTEM A

Anna Curir cla defines element in centre of  $U(1) \times G$ 

her  $T_{\pm}$ 

 $dM = T_{\pm}dA_{\pm} + \Omega_{\pm}dJ$  and angular momentum (3) are conthis lacks a factor of 1/4 in front of the first term compared with the first term compared with the first gauge grant of the first term compared with the first

Generale  $= that \text{ this } lacks \frac{U(1) \times G}{a \text{ factor}} of 1/4 in front of value . Note that this lacks a factor of <math>1/4 in front of \Omega_{\pm}$ Anna Curir's L. value.

### Global Constraint on Gauge Symmetry:

$$G_{\mathsf{global}} = rac{U(1) imes G}{\langle C 
angle} \cong rac{U(1) imes G}{\mathbb{Z}_{\kappa}}$$

Exemplify for SU(5) GUT's and Standard Model constructions Including for globally consistent three family SM

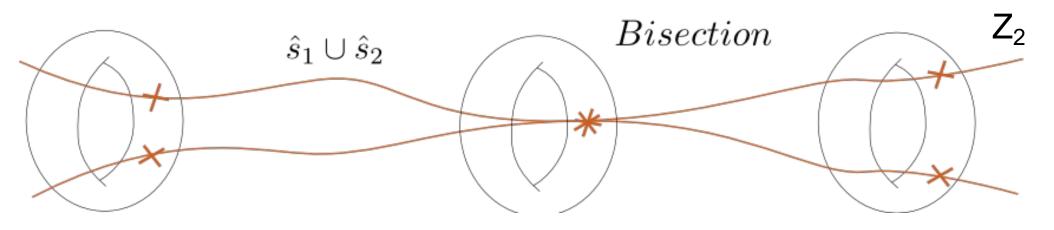
More later

## I.c Discrete Abelian Gauge Symmetries Z<sub>n</sub>

In Standard Model one often introduces  $\mathbb{Z}_2$  (R-parity) to prevent baryon number violating couplings

## I.c Discrete Abelian Gauge Symmetries - Z<sub>n</sub>

Geometric origin: torus fibrations that do not admit a section, but a multi-section



Earlier work: [Witten; deBoer, Dijkgraaf, Hori, Keurentjes, Morgan, Morrison, Sethi;...]

Recent extensive efforts'14-'16: [Braun, Morrison; Morrison, Taylor;

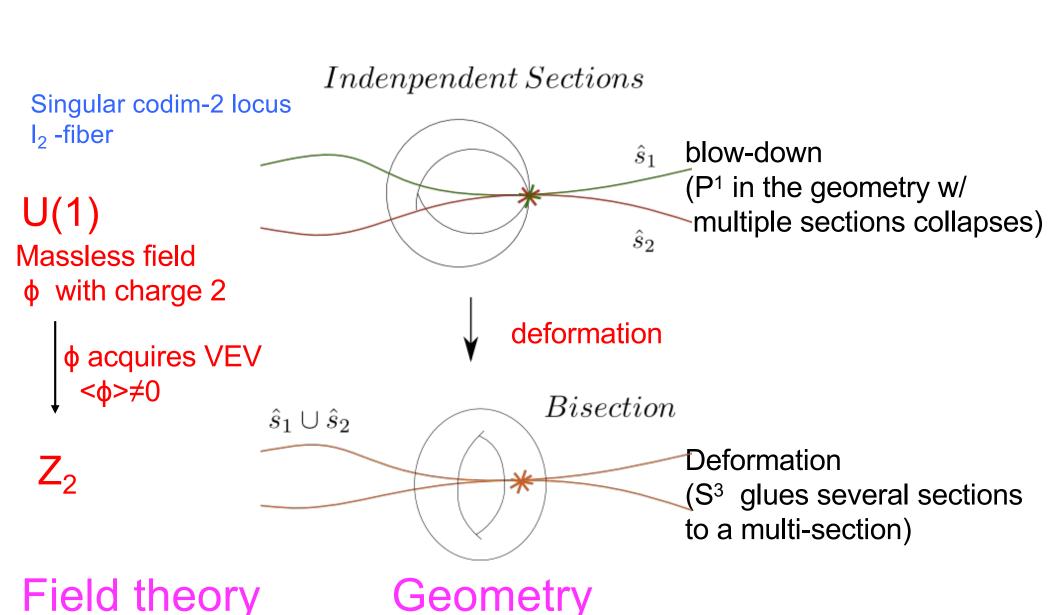
Klevers, Mayorga-Pena, Oehlmann, Piragua, Reuter; Anderson, Garcia-Etxebarria,

Grimm; Braun, Grimm, Keitel; Mayrhofer, Palti, Till, Weigand; M.C., Donagi, Klevers,

Piragua, Poretschkin; Grimm, Pugh, Regalado; M.C., Grassi, Poretschkin;...]

Standard Model with Z<sub>2</sub> matter parity [M.C., Lin, Liu, Oehlmann 1807.01320]

## Transition between continuous and discrete symmetry: $U(1) \rightarrow Z_2$ example



### I.c Discrete Abelian Gauge Symmetries



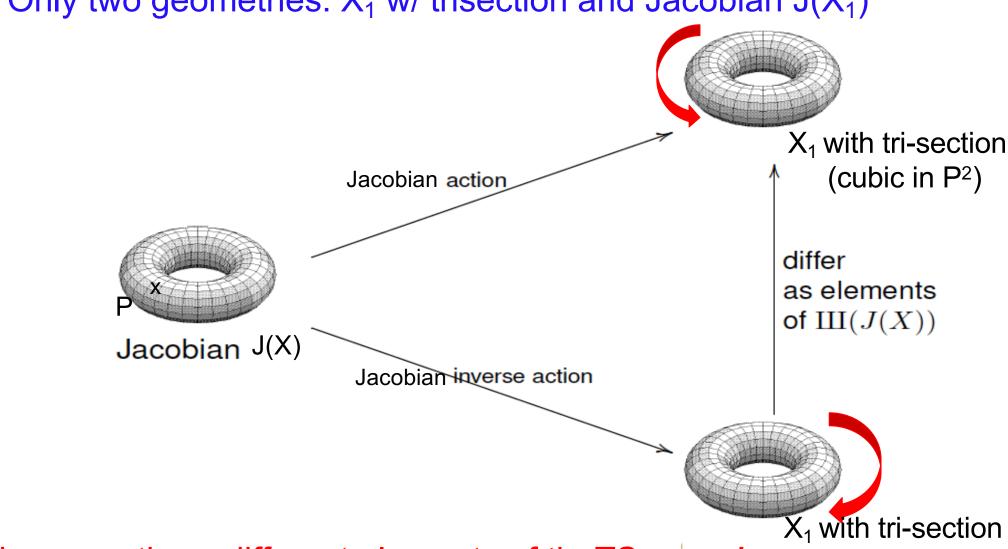
Geometries with n-section Tate-Shafarevich Group Z<sub>n</sub>

- Z<sub>2</sub> [Anderson, Garcia-Etxebarria, Grimm; Braun, Grimm, Keitel; Mayrhofer, Palti, Till, Weigand'14]
- Z<sub>3</sub> [M.C.,Donagi,Klevers,Piragua,Poretschkin 1502.06953] No time, but

## Tate-Shafarevich group and Z<sub>3</sub>

[M.C., Donagi, Klevers, Piragua, Poretschkin 1502.06953]

Only two geometries:  $X_1$  w/ trisection and Jacobian  $J(X_1)$ 



There are three different elements of theTS group!

(cubic in P<sup>2</sup>)

Shown to be in one-to-one correspondence with three M-theory vacua.

## II. Particle physics Constructions

#### Initial focus: F-theory with SU(5) Grand Unification

[Donagi, Wijnholt'08] [Beasley, Heckman, Vafa'08]...

#### **Model Constructions:**

local

[Donagi, Wijnholt'09-10]...[Marsano, Schäfer-Nameki, Saulina'09-11]...

Review: [Heckman]

#### global

```
[Blumehagen, Grimm, Jurke, Weigand'09] [M.C., Garcia-Etxebarria, Halverson'10]... [Marsano, Schäfer-Nameki'11-12]... [Clemens, Marsano, Pantev, Raby, Tseng'12]...
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#### Progress on Standard Model:

Standard Model building blocks (via toric techniques) [Lin, Weigand'14]



## II. Particle physics Constructions

Globally consistent models via torique techniques

he starting polypersurface fiber Grofand base Ene The following discussion ypersurface emitting central mountainentianal parablain Theregises of cients  $a_q$  and the savarhears of  $\pm 2.23$ . have to for proportional of sections of a topic les efficieles de Le Le detertaine plaiset linfet Benches try citie realist per gan us fibre oric variety  $\mathbb{P}_{E}$  hwhich exthefance contactions (2.23) of ethers we base B, orde  $E_{F_i} \not\models_{F_i} = \underbrace{\frac{\text{coefficies Rs}}{\text{out(Che)}^n \text{of}}}_{\text{letib}} \xrightarrow{\text{prod the variables } x} \underbrace{\frac{\text{letiber of the variables } x}{\text{letiber}^n}}_{F_i} \xrightarrow{\text{prod the variables } x} \underbrace{\frac{\text{letiber of the letiber}}{\text{letiber}^n}}_{\text{letiber of the letiber}} \xrightarrow{\text{letiber of the variables } x}_{\text{letiber of the letiber of the letibe$ 2D toric varietyklevers, Pena, behlmant, interpresente of the pov  $\mathbb{P}_{F_i} \longrightarrow \mathbb{P}_{F_i}^B(\mathcal{S}_7, \mathcal{S}_9)$  $E_{F_i} = \{p_{F_i} = 0\}$  in  $\mathbb{P}_{F_i}^{\mathbb{P}_{F_i}}$ denotes the total space of this fibration. The structure of its  $\operatorname{metriz}_{\mathbf{E}}$ dibtic $\mathfrak{A}$ to  $\mathfrak{A}$ 2 coordinates  $x_k$  on the fiber  $\mathbb{P}_{F_k}$  are in general non-trivial sections of Pine, bur  $\mathbb{P}_{F_e}$ rivia pondinate bise Pri and softadan et sittletet et alirana en et utilies fibraties  $\operatorname{ext}$  we imposible a first the property of the part of the property of the  $\operatorname{E}_{F_i}$ oefficients  $a_a$  have the color displace in the Fig. Then, to Fibratisectiepends only on the anti-caronical configuration,  $\overline{\mathcal{K}}$  hich is the Calabi-Value  $\frac{8}{8}$  two additional  $\frac{1}{8}$ , and  $\frac{1}{8}$ , divisor classes. Fix the calabi-Value  $\frac{1}{8}$  in Figure 1 in  $\frac{1}{8}$ , the calabi-Value  $\frac{1}{8}$  in  $\frac{1}{8}$ 

work are n = 0, 4. We refer to [01, 09] for inote details on the following discussing

#### iii. Chiral index for D=4 matter:

Standard Model with three families of quarks and leptons

we have to know the handlogy classes of the associated matter

# iii. Chiral index for D=4 matter: $\Sigma_{\mathbf{R}}$

$$\chi(\mathbf{R}) = \int_{\mathcal{C}_{\mathbf{R}}^w} G_4$$

$$\Sigma_{\mathbf{R}}$$

$$c_{mat} - c_{mat} \longrightarrow \mathcal{C}_{\mathbf{R}}$$

a) construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the discretion of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the discretion of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construct  $G_4$  (=d $C_3$ ) flux by computing  $H_V^{(2,2)}(\hat{X})$  [so-Which vertical stick of the construction of the constr

In this section we determine the matter curves  $\Sigma_{\mathbf{R}}$  and the

iv. Global consistency on paragraphs can be a that three of the six determination is complicated by the fact that three of the six the base  $B_{2}$  where the elliptic file  $\mathcal{E}$  becomes reducible are to Their irreducible components are multiple different matter

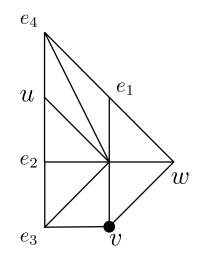
- a) satisfied for integerland to be complete intersection a b) constraint on line for integer malured flux for ever, the same for integer and the same street in the same and the same street in the same st
  - from the two equations of the original reducible codimensio

# C.f., Standard Model building blocks (via toric techniques) initiated in [Lin, Weigand'14]; SM x U(1) [1604.04292]

#### Standard Model

[M.C., Klevers, Peña, Oehlmann, Reuter, 1503.02068]

F<sub>11</sub> polytope



 $P_{F11}$ 

 $\mathbb{P}^2$  [u:v:w] with four non-generic blow-ups [e<sub>1</sub>:e<sub>2</sub>:e<sub>3</sub>:e<sub>4</sub>]  $\mathrm{Bl}_1\mathbb{P}^2(1,1,2)$ 

#### Elliptic curve:

$$p_{F_{11}} = s_1 e_1^2 e_2^2 e_3 e_4^4 u^3 + s_2 e_1 e_2^2 e_3^2 e_4^2 u^2 v + s_3 e_2^2 e_3^2 u v^2 + s_5 e_1^2 e_2 e_4^3 u^2 w + s_6 e_1 e_2 e_3 e_4 u v w + s_9 e_1 v w^2$$

Hypersurface constraint in P<sub>F11</sub>

 $Bl_3\mathbb{P}^3$ 



 $\mathbb{P}^4$ 

Construction of Calabi-Yau four-fold

Construction of Salabi-Yau four-fold denotes the total space of this fibration

parametrized by two divisors in B,  $dE_{F_{11}} \longrightarrow X_{F_{11}} \longrightarrow X_{F_{11}}$ .

Coordinates  $x_k$  on the fiber  $\mathbb{P}_F$  of the base B [Toric techniques via Stanley-Reisner ideal]  $\to$  Next we impose equatic section Line Bundle Consist

Next we impose equatic section						
Section	Line Bundle	have to t	$s_1$			
$\boldsymbol{u}$	$\mathcal{O}(H - E_1 - E_2 - E_4 + S_9 + [K_B])$		60			
v	$\mathcal{O}(H - E_2 - E_3 + \mathcal{S}_9 - \mathcal{S}_7)$	tion of tl	$s_2$			
$\boldsymbol{w}$	$\mathcal{O}(H-E_1)$		$s_3$			
$e_1$	$\mathcal{O}(E_1-E_4)$	(0.00	$s_4$			
$e_2$	$\mathcal{O}(E_2-E_3)$	1  ion  (2.23)	- 1			
$e_3$	$\mathcal{O}(E_3)$	, , , , , , , , , , , , , , , , , , ,	$s_5$			
$e_4$	$\mathcal{O}(E_4)$	generic <sub>I</sub>	$s_6$			

have to t	$s_1$
tion of tl	$s_2$
01011 01 01	$s_3$
ion (2.22	$s_4$

$$\begin{array}{c} \text{ion } (2.23) \\ \text{generic I} \\ s_6 \end{array}$$

 $S_7$ 

 $s_8$ 

 $s_9$ 

 $s_{10}$ 

ın 
$$\mathbb{F}_{F_i}$$
. The total Calabi-Ya denoted by  $X_{F_i}$  in the follo

$$K_{B}^{-1}$$
 - anti-canonical divisor  $\overline{\mathcal{K}}$ 

Diffe Duffdie	
$\mathcal{O}_B(3[K_B^{-1}] - \mathcal{S}_7 - \mathcal{S}_9)$	he two
$\mathcal{O}_B(2[K_B^{-1}] - \mathcal{S}_9)$ $\mathcal{O}_B([K_B^{-1}] + \mathcal{S}_7 - \mathcal{S}_9)$	$\in K_{\mathbb{P}^B_{F_i}}^{-1},$
$\mathcal{O}_B([\Lambda_B] + \mathcal{O}_7 - \mathcal{O}_9)$ $\mathcal{O}_B(2\mathcal{S}_7 - \mathcal{S}_9)$	$\operatorname{clearly}^{\operatorname{\scriptscriptstyle{\mathrm{II}}}}$
$\mathcal{O}_B(2[K_B^{-1}] - \mathcal{S}_7)$	rface (2
$K_B^{-1} \ \mathcal{O}_B(\mathcal{S}_7)$	the fibr
$\mathcal{O}_B([K_B^{-1}] + \mathcal{S}_9 - \mathcal{S}_7)$	tion st
$\mathcal{O}_B(\mathcal{S}_9)$	
$\mathcal{O}_B(2\mathcal{S}_9 - \frac{\mathcal{S}_7)}{\mathcal{C}_{F_i}}$	$\mathbf{\hat{\lambda}}F_{i}$ .

Fibration depends only on additional S<sub>7</sub> and S<sub>9</sub> divisor classes.

$$S_7$$
  $S_9$ 

 $\overset{\forall}{B}$ 

Construction of Calabi-Yau four-fold  $\rightarrow$  Divisors. Here  $\mathbb{P}_{F_i}^B(D, D)$  denotes the total space of this fibration

parametrized by two divisors in B,  $dE_{F_{11}} \subset \mathbb{P}_{F_{11}} \longrightarrow X_{F_{11}} \subset \mathbb{P}_{F_{11}} \longrightarrow X_{F_{11}} \subset \mathbb{P}_{F_{11}} \subset$ 

Over the laction  $s_3$  filter the entire of the settent  $s_4$  filter the period of the settent  $s_4$  filter the settent

Over the lobestrigiable phenological distribution of the phenolog

Cartan divisors of these improse groups on (2.23) in  $\mathbb{P}_{F_i}(D,D)$ . Consist the coefficients  $a_{\tilde{e}}$  have to take values in  $F_i$  the two

 $\mathsf{E}_1^{\mathrm{SU}(2)} = [e_1] \,, \quad \mathsf{E}_1^{\mathrm{SU}(3)} = [e_2] \quad \mathsf{E}_2^{\mathrm{SU}(3)} = [u] \text{ the anti-canonical bundle } K_{\mathbb{P}_{F_i}}^{-1} \,,$ 

Two rational exertions: equation (2.23) imposed in  $\mathbb{P}^B_{F_i}(D, D)$  clearly  $B_{F_i}(D, D)$  clearly on  $B_{F_i}(D, D)$  the hypersurface (2.23)

 $\hat{s}_0 = X_{F_{11}} \cap \{v = 0\}$ :  $[1:0:s_1:1:1:-s_5:1]_{0}$ 

 $\hat{s}_1 = X_{F_{11}} \cap \{e_4 = 0\}: \quad [s_9:1:1:-s_3:1:1:0] - \text{section-associated without the points of the section of the sect$ 



$$\mathcal{C}_{F_i} \longrightarrow X_{F_i}$$
.

Standard Model gauge symmetry: SU(3) x SU(2) x U(1)

## Global Standard Model Gauge Symmetry & Matter Reps.

gauge algebra  $\mathfrak{su}(3) \oplus \mathfrak{su}(2) \oplus \mathfrak{u}(1)$ 

#### Shioda map:

[M.C., Lin, 1706.08521]

C-central element

$$\sigma(\hat{s}_1) = S_1 - S_0 + \frac{1}{2} E_1^{\mathfrak{su}(2)} + \frac{1}{3} (2 E_1^{\mathfrak{su}(3)} + E_2^{\mathfrak{su}(3)}) \Rightarrow C^6 = 1,$$



$$G_{\mathsf{global}} = [SU(3) \times SU(2) \times U(1)]/\langle C \rangle \cong [SU(3) \times SU(2) \times U(1)]/\mathbb{Z}_6.$$

Matter (at co-dim 2 singularities):

$$(\mathbf{3},\mathbf{2})_{\frac{1}{6}}, \ (\mathbf{1},\mathbf{2})_{-\frac{1}{2}}, \ (\overline{\mathbf{3}},\mathbf{1})_{-\frac{2}{3}}, \ (\overline{\mathbf{3}},\mathbf{1})_{\frac{1}{3}}, \ (\mathbf{1},\mathbf{1})_{1}$$

Compatible with the Z<sub>6</sub> global constraint



Construct G<sub>4</sub> for chiral index & D3-tadpole constraint

### **Standard Model:**

Hyperplane divisor class

$$H=4\overline{\mathcal{K}}$$

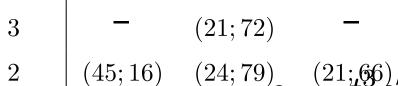
Base  $B = \mathbb{P}^3$  Divisors in the base:

$$S_7 = n_7 H$$



$n_7 \setminus n_9$	1		$y = \mathfrak{F}_7 H_1$	
7	_	$(27;16)^{-5}$	$\underline{n} = \underline{n}_9 H_1$	P <sup>3</sup> _
6	_	(12; 81)	(21;42)	_
5	_	-	(12; 57)	(30; 8)

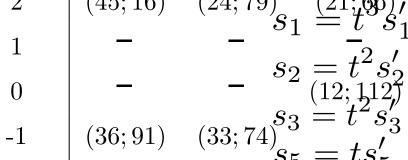
(30; 32)

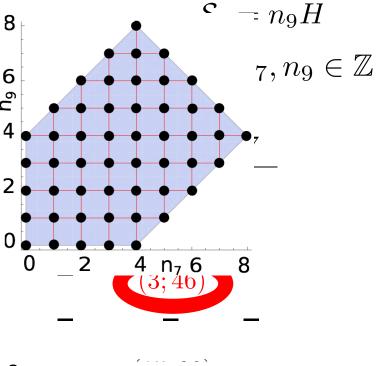


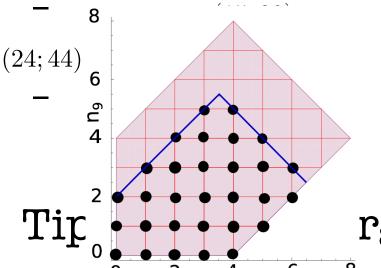
(42;4)

4

-2







# II. Landscape of Standard Models Toric analysis

[M.C., J. Halverson, L. Lin, M. Liu and J. Tian, 1903.0009]

a) Take the same toric elliptic fibration as before:

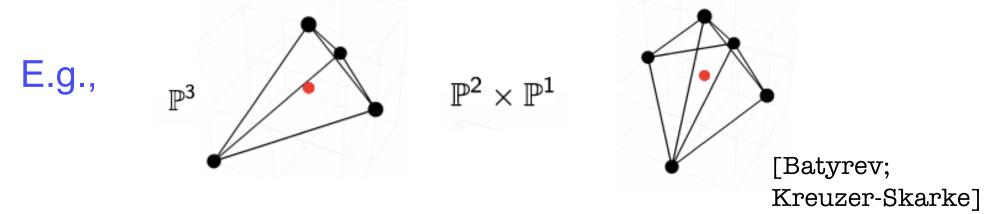
hyperplane constraint in 2D reflexive polytope F<sub>11</sub>

Gauge symmetry: 
$$\frac{\mathrm{SU}(3) \times \mathrm{SU}(2) \times \mathrm{U}(1)}{\mathbb{Z}_6}$$

Global gauge symmetry

[M.C., Lin, 1706.08521]

b) Take bases B, associated with 3D reflexive polytopes



For each reflexive polytope, different bases B are associated with different fine-star-regular triangulations of a chosen polytope [Triangulations determine intersections of divisors] Triangulations grow exponentially with the complexity of a polytope

# c) Specific choice of divisors: $S_{7,9} = \overline{\mathcal{K}}$

[anti-canonical divisor of the base B – fixed by the polytope]

SU(3) an SU (2) divisors 
$$S_9$$
 and  $S_3$  with class  $\overline{\mathcal{K}}$   $g_{3,2}^2 = 2/\mathrm{vol}(\overline{\mathcal{K}})$ 

U(1) - (height-pairing) divisor volume  $5\overline{\mathcal{K}}/6 \rightarrow$ 

$$\frac{5}{3}g_Y^2 = \frac{2}{\operatorname{vol}(\overline{\mathcal{K}})}$$



Standard Model with gauge coupling unification!

$$g_3^2 = g_2^2 = 5/3g_Y^2$$

# Connected torically to Pati-Salam Model SU(4)<sub>C</sub>xSU(2)<sub>L</sub>xSU(2)<sub>R</sub>

c.f., [M.C., Klevers, Peña, Oehlmann, Reuter, 1503.02068]

#### Non-torically connected to SU(5) GUT

[Taylor, Turner 1906.11092; & Ranghuram 1912.10991]

# d) Remaining conditions:

- iii. 3-families of quarks and leptons (chiral index)
- iv. D3-tadpole constraints

Technical, no time

- Construct G<sub>4</sub> flux in terms of (1,1)-forms, Poincaré dual to divisor classes

  c.f., [Lin, Mayrhofer, Till, Weigand, 1508.00162]

  [M.C., Grassi, Klevers, Piragua, 1306.3987]
- Chirality, D3 tadpole and G<sub>4</sub> integrality expressed in terms of intersection numbers of divisors in the base B → Geometric conditions!
- In the case  $S_{7.9} = \overline{\mathcal{K}}$  and  $n_F$  families, the D3 tadopole:

$$n_{D3}(n_F, \overline{K}^3) = 12 + \frac{5\overline{K}^3}{8} - \frac{5n_F^2}{2\overline{K}^3} \in \mathbb{Z}_{\geq 0}$$

Geometrized D3-tadpole condition

Depends only on the polytope and not on triangulation 

Universality of the Standard Model

# Landscape count for n<sub>F</sub>=3 families:

toric

$$12 + \frac{5}{8}\overline{\mathcal{K}}^3 - \frac{45}{2\overline{\mathcal{K}}^3} \in \mathbb{Z}_{\geq 0} \quad \text{satisfied for } \overline{\mathcal{K}}^3 \in \{2, 6, 10, 18, 30, 90\}$$

- Out of 4319 3D reflective polytopes → 708 satisfy the constraint (many of them with a large number of lattice points).
- Triangulation of polytopes can be handled combinatorially (each corresponds to a different basis B).
   It can be implemented on computer, e.g., in SageMath:
  - i) for 237 polytopes w/ < 15 lattice points  $\rightarrow$  414310 MSSM models.
- ii) for 471 polytopes w/ ≥ 15 lattice points exp. growing comp. time → counting via fine-regular triangulation of facets & estimate regular fine-star triang. c.f., [Halverson, Tian, 1610.08864]



• Provide a bound:  $7.6 \times 10^{13} \lesssim N_{\rm SM}^{\rm toric} \lesssim 1.6 \times 10^{16}$ 

# Summary

Globally consistent F-theory Standard Models (Toric techniques w/elliptic fibration: hypersurface in F<sub>11</sub>)



First three family Standard Models

Anticipated: tip of the iceberg



Indeed, geometric advances

Landscape of globally consistent Standard Models w/ exact chiral spectrum of three-families of quarks &leptons & gauge coupling unification > quadrillion models

# III. Further Analysis

#### III.a Moduli stabilization

Related to issues of supersymmetry breaking, cosmological implications, dark matter candidates...



# Moduli Stabilization for quadrillion Standard Models

[M.C., Long, Halverson, Lin, 2004.00630]

Under which conditions moduli stabilization can be pursued via effective field theory techniques w/ g<sub>s</sub> perturbative (à la KKLT or Large Volume Scenario):

- i) gauge coupling constraint:  $\alpha_{1,2,3} = \alpha_{GUT} = (g_s \ell_s^4)/\mathrm{Vol}(\bar{K}) \sim 1/25$  $\rightarrow \mathrm{Vol}(\overline{\mathcal{K}}) \lesssim \mathcal{O}(100)$
- ii) all divisor and curves  $w/vol(C_a) \ge 1$  (in string units) in order to suppress world-sheet and ED3 instanton contributions, c.f.,  $e^{-2\pi n vol(c)} \ll 1$ .

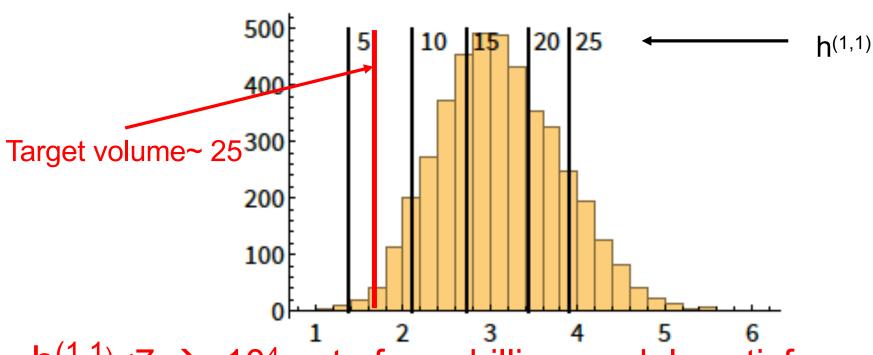
#### → Stretched Kähler cone

- where all divisors & curves w/volumes >1
- Since  $\overline{\mathcal{K}} = -\sum D_i \rightarrow \text{Vol}(\overline{\mathcal{K}})$  expected to be typically large

Distribution of  $Min(Vol(\mathbb{K}))$ 

[one triangulation per each of 4319 polytopes]

$$\log_{10}(\min(\operatorname{vol}(\overline{K}_B)))$$



h<sup>(1,1)</sup><7 →~10<sup>4</sup> out of quadrillion models satisfy constraints

#### Comments:

- Moduli stabilization scenarios, based on effective theory & perturbative g<sub>s</sub> (KKLT,LVS), significantly reduces the number of viable Standard Models with gauge coupling unification.
- Moduli stabilization could take place in a regime where effective theory & perturbative string theory approaches fail.
   poorly explored/difficult to explore.
- Could abandon to have only Standard Model and/or gauge coupling unification → typically leads to additional D7-(p,q) sectors w/ interesting dark gauge sector implications.

c.f., Halverson, Long, Nelson, Salinas 1909.05257



Further exploration of other Standard Model constructions

# III.a Counting of vector matter pairs

[...Donagi, Wijnholt '08,...,Bies, Mayrhofer, Pehle, Weigand '14,'17]

Depends on C<sub>3</sub> potential, encoded in intermediate the Jacobian of Y<sub>4</sub>.

When restricted to the matter curve C,  $C_3$  defines a line bundle  $\mathcal{L}$  w/ massless chiral modes  $\subset H^0(C; \mathcal{L})$  massless anti-chiral modes  $\subset H^1(C; \mathcal{L})$ 

[chiral index  $\chi = h^0 - h^1$  topological invariant (depends on  $G_4 = dC_3$ )]

H<sup>i</sup>(*C,L*) – computation via algorithm implemented in computer algebra system CAP [Bies '17; Bies, Posur '19]

# Counting of vector matter pairs

For the quadrillion Standard Models the analysis difficult due to the complexity of the construction and high genera of matter curves.

$$g = 1 + 9/2\overline{\mathcal{K}}^3$$

Goal: to determine the range of complex structure moduli of the F-theory compactification, for which we have the Minimal Supersymmetric Standard Model (one Higgs doublet pair, no other vector pair exotics).

Making progress...

[Martin Bies, M.C., Ron. Donagi, L. Lin, M. Liu, Fabian Rühle, 2007.0009] [Bies, M.C. Donagi, Liu, 2102.10115, 2104.08297] [Bies, M.C. Donagi, Marielle Ong, 2205.00008, 2307.02535]

#### Outlook

Particle physics models in F-theory compactifications have come a long way, but there is much more to go.

Technical advances, to be pursued:

Exact matter spectrum for quadrillion Standard Models
Yukawa couplings (some progress for a toy model)
[M.C., Lin, Liu, Zoccarato, Zhang 1906.10119]

Systematic exploration of other particle physics models (possibly beyond toric techniques) →

c.f., W. Taylor's et al. talk(s)

Thank you!