# Homogeneous Hyper-Hermitian Metrics Which Are Conformally HyperKähler

María Laura Barberis

Vienna, Preprint ESI 925 (2000)

August 9, 2000

Supported by Federal Ministry of Science and Transport, Austria Available via http://www.esi.ac.at

## HOMOGENEOUS HYPER-HERMITIAN METRICS WHICH ARE CONFORMALLY HYPER-KÄHLER

#### MARÍA LAURA BARBERIS

ABSTRACT. Let g be a hyper-Hermitian metric on a simply connected hypercomplex four-manifold  $(M, \mathcal{H})$ . We show that when the isometry group I(M, g)contains a subgroup acting simply transitively on M by hypercomplex isometries then the metric g is conformal to a hyper-Kähler metric. We describe explicitely the corresponding hyper-Kähler metrics and it follows that, in four dimensions, these are the only hyper-Kähler metrics containing a homogeneous metric in its conformal class.

#### 1. Preliminaries

A hypercomplex structure on a 4n-dimensional manifold M is a family  $\mathcal{H} =$  $\{J_{\alpha}\}_{\alpha=1,2,3}$  of fibrewise endomorphisms of the tangent bundle TM of M satisfying:

(1.1) 
$$J_{\alpha}^{2} = -I, \quad \alpha = 1, 2, 3 \qquad J_{1}J_{2} = -J_{2}J_{1} = J_{3},$$
(1.2) 
$$N_{\alpha} \equiv 0, \quad \alpha = 1, 2, 3$$

$$(1.2) N_{\alpha} \equiv 0, \alpha = 1, 2, 3$$

where I is the identity on the tangent space  $T_pM$  of M at p for all p in M and  $N_\alpha$ is the Nijenhuis tensor corresponding to  $J_{\alpha}$ :

$$N_{\alpha}(X,Y) = [J_{\alpha}X, J_{\alpha}Y] - [X,Y] - J_{\alpha}([X,J_{\alpha}Y] + [J_{\alpha}X,Y])$$

for all X, Y vector fields on M. A differentiable map  $f: M \to M$  is said to be hypercomplex if it is holomorphic with respect to  $J_{\alpha}$ ,  $\alpha = 1, 2, 3$ . The group of hypercomplex diffeomorphisms on  $(M, \mathcal{H})$  will be denoted by  $\operatorname{Aut}(\mathcal{H})$ .

A riemannian metric g on a hypercomplex manifold  $(M, \mathcal{H})$  is called hyper-Hermitian when  $g(J_{\alpha}X, J_{\alpha}Y) = g(X, Y)$  for all vectors fields X, Y on  $M, \alpha =$ 1, 2, 3.

Given a manifold M with a hypercomplex structure  $\mathcal{H} = \{J_{\alpha}\}_{\alpha=1,2,3}$  and a hyper-Hermitian metric g consider the 2-forms  $\omega_{\alpha}$ ,  $\alpha = 1, 2, 3$ , defined by

(1.3) 
$$\omega_{\alpha}(X,Y) = g(X,J_{\alpha}Y).$$

The metric g is said to be hyper-Kähler when  $d\omega_{\alpha} = 0$  for  $\alpha = 1, 2, 3$ .

It is well known that a hyper-Hermitian metric g is conformal to a hyper-Kähler metric  $\tilde{g}$  if and only if there exists an exact 1-form  $\theta \in \Lambda^1 M$  such that

$$(1.4) d\omega_{\alpha} = \theta \wedge \omega_{\alpha}, \quad \alpha = 1, 2, 3$$

where, if  $g = e^f \tilde{g}$  for some  $f \in C^{\infty}(M)$ , then  $\theta = df$ .

We prove the following result:

<sup>1991</sup> Mathematics Subject Classification. Primary 53C15, 53C25, 53C30.

Key words and phrases. hyper-Hermitian metric, hypercomplex manifold, conformally hyper-Kähler metric.

The author was partially supported by CONICET, ESI (Viena) and FOMEC (Argentina).

**Theorem 1.1.** Let  $(M, \mathcal{H}, g)$  be a simply connected hyper-Hermitian 4-manifold. Assume that there exists a Lie group  $G \subset I(M,g) \cap Aut(\mathcal{H})$  acting simply transitively on M. Then g is conformally hyper-Kähler.

We conclude that one of the hyper-Kähler metrics constructed by the Gibbons-Hawking ansatz [2] contains a homogeneous hyper-Hermitian metric in its conformal class. This hyper-Hermitian metric is not symmetric and has negative sectional curvature [1].

As a consequence of Theorem 1.1 and the results in [1] we obtain that the following symmetric riemannian metrics are conformally hyper-Kähler:

- the riemannian product of the canonical metrics on  $\mathbb{R} \times S^3$ ;
- the riemannian product of the canonical metrics on  $\mathbb{R} \times \mathbb{R}H^3$ , where  $\mathbb{R}H^3$ denotes the real hyperbolic space;
- the canonical metric on the real hyperbolic space  $\mathbb{R}H^4$ .

Acknowledgements. I would like to thank the organizers of the program Holonomy Groups in Differential Geometry for giving me the opportunity to visit the Erwin Schrödinger Institute, Vienna. I am also grateful to D. Alekseevsky, I. Dotti Miatello, L. Ornea and S. Salamon for useful conversations.

#### 2. Proof of the main theorem

Proof of Theorem 1.1. Since G acts simply transitively on M then M is diffeomorphic to G and therefore the hypercomplex structure and hyper-Hermitian metric can be transferred to G and will also be denoted by  $\{J_{\alpha}\}_{\alpha=1,2,3}$  and g, respectively. Since G acts by hypercomplex isometries it follows that both  $\{J_{\alpha}\}_{\alpha=1,2,3}$ and g are left invariant on G. All such simply connected Lie groups were classified in [1], where it is shown that the Lie algebra  $\mathfrak{g}$  of G is either abelian or isomorphic to one of the following Lie algebras (we fix an orthonormal basis  $\{e_j\}_{j=1,\dots,4}$  of  $\mathfrak{g}$ ):

- 1.  $[e_2, e_3] = e_4$ ,  $[e_3, e_4] = e_2$ ,  $[e_4, e_2] = e_3$ ,  $e_1$  central; 2.  $[e_1, e_3] = e_1$ ,  $[e_2, e_3] = e_2$ ,  $[e_1, e_4] = e_2$ ,  $[e_2, e_4] = -e_1$ ; 3.  $[e_1, e_j] = e_j$ , j = 2, 3, 4;

- 4.  $[e_3, e_4] = \frac{1}{2}e_2$ ,  $[e_1, e_2] = e_2$ ,  $[e_1, e_j] = \frac{1}{2}e_j$ , j = 3, 4.

Observe that in case 1 above M is diffeomorphic to  $\mathbb{R} \times S^3$  while in the remaining cases it is diffeomorphic to  $\mathbb{R}^4$ , therefore in all cases any closed form on M is exact. We now proceed by finding in each case a closed form  $\theta \in \Lambda^1 \mathfrak{g}^*$  satisfying (1.4). Note that we work on the Lie algebra level since g and  $\omega_{\alpha}$  are all left invariant on G. Let  $\{e^j\}_{j=1,\dots,4} \subset \Lambda^1 \mathfrak{g}^*$  be the dual basis of  $\{e_j\}_{j=1,\dots,4}$ . From now on we will write  $e^{ij}$  to denote  $e^i \wedge e^j \wedge \cdots$ . In all the cases below the 2-forms  $\omega_\alpha$  are determined from (1.3) in terms of the hypercomplex structures constructed in [1].

Case 1. The 2-forms  $\omega_{\alpha}$  are given as follows:

$$\omega_1 = -e^{12} - e^{34}, \qquad \omega_2 = -e^{13} + e^{24}, \qquad \omega_3 = -e^{14} - e^{23}.$$

To calculate  $d\omega_{\alpha}$  we obtain first  $de^{j}$  (recall that  $d\sigma(x,y) = -\sigma[x,y]$  for  $\sigma \in \Lambda^{1}\mathfrak{g}^{*}$ ):

(2.1) 
$$de^1 = 0$$
,  $de^2 = -e^{34}$ ,  $de^3 = e^{24}$ ,  $de^4 = -e^{23}$ .

These equations and the fact that  $d(\sigma \wedge \tau) = d\sigma \wedge \tau + (-1)^r \sigma \wedge d\tau$  for all  $\sigma \in \Lambda^r \mathfrak{g}^*$ give the following formulas:

$$d\omega_1 = -e^{134}, \qquad d\omega_2 = e^{124}, \qquad d\omega_3 = -e^{123}$$

from which we conclude that (1.4) holds for  $\theta = e^1$ , which is closed and therefore exact since G is diffeomorphic to  $\mathbb{R} \times S^3$ . We conclude that this hyper-Hermitian metric, which, as shown in [1], is homothetic to the riemannian product of the canonical metrics on  $\mathbb{R} \times S^3$ , is conformal to a hyper-Kähler metric.

Case 2. In this case we have the following equations for  $\omega_{\alpha}$ :

$$\omega_1 = e^{14} - e^{23}, \qquad \omega_2 = -e^{12} + e^{34}, \qquad \omega_3 = -e^{13} - e^{24}.$$

and we calculate

$$(2.2) de^1 = -e^{13} + e^{24}, de^2 = -e^{23} - e^{14}, de^3 = 0, de^4 = 0,$$

(2.3) 
$$d\omega_1 = -2e^{134}, \qquad d\omega_2 = -2e^{123}, \qquad d\omega_3 = 2e^{234}$$

so that (1.4) is satisfied for  $\theta = 2e^3$ , which again is closed, so this hyper-Hermitian metric is also conformal to a hyper-Kähler metric. In this case the hyper-Hermitian metric is homothetic to the riemannian product of the canonical metrics on  $\mathbb{R} \times \mathbb{R}H^3$ , where  $\mathbb{R}H^3$  denotes the real hyperbolic space.

Case 3. In this case the 2-forms  $\omega_{\alpha}$  are given as follows:

$$\omega_1 = -e^{12} - e^{34}, \qquad \omega_2 = -e^{13} + e^{24}, \qquad \omega_3 = -e^{14} - e^{23}$$

and a calculation of exterior derivatives gives:

(2.4) 
$$de^1 = 0, de^j = -e^{1j}, j = 2, 3, 4$$

(2.5) 
$$d\omega_1 = 2e^{134}, \qquad d\omega_2 = -2e^{124}, \qquad d\omega_3 = -2e^{123}$$

so that (1.4) is satisfied for  $\theta = -2e^1$ . This hyper-Hermitian metric is homothetic to the canonical metric on the real hyperbolic space  $\mathbb{R}H^4$ .

Case 4. In this case we have the following equations for  $\omega_{\alpha}$ :

$$\omega_1 = -e^{12} + e^{34}, \qquad \omega_2 = -e^{13} - e^{24}, \qquad \omega_3 = e^{14} - e^{23}$$

and we calculate

$$(2.6) de^1 = 0, de^2 = -e^{12} - \frac{1}{2}e^{34}, de^j = -\frac{1}{2}e^{1j}, j = 3, 4$$

(2.7) 
$$d\omega_1 = -\frac{3}{2}e^{134}, \qquad d\omega_2 = \frac{3}{2}e^{124}, \qquad d\omega_3 = \frac{3}{2}e^{123}$$

so that (1.4) is satisfied for  $\theta = -\frac{3}{2}e^{1}$ . This hyper-Hermitian metric is not symmetric and has negative sectional curvature (cf. [1]).

Remark 2.1. All the hyper-Hermitian manifolds  $(M, \mathcal{H}, g)$  considered above admit a connection  $\nabla$  such that:

$$\nabla q = 0$$
,  $\nabla J_{\alpha} = 0$ ,  $\alpha = 1, 2, 3$ 

and the (3,0) tensor c(X,Y,Z) = g(X,T(Y,Z)) is totally skew-symmetric, where T is the torsion of  $\nabla$ . Such a connection is called an HKT connection (cf. [3]). In case M is diffeomorphic to  $\mathbb{R} \times S^3$  it can be shown that, moreover, the corresponding 3-form c is closed.

### 3. COORDINATE DESCRIPTION OF THE HYPER-KÄHLER METRICS

In this section we will use global coordinates on each of the Lie groups considered in the previous section to describe the corresponding hyper-Kähler metrics. This will allow us to identify the hyper-Kähler metric in §2, Case 4, with one constructed by the Gibbons-Hawking ansatz [2].

Case 1. 
$$G = \mathbb{H}^* = GL(1, \mathbb{H}) = \left\{ \begin{pmatrix} x & -y & -z & -t \\ y & x & -t & z \\ z & t & x & -y \\ t & -z & y & x \end{pmatrix} : (x, y, z, t) \in \mathbb{R}^4 - \{0\} \right\}.$$

We obtain a basis of left invariant 1-forms on G as follows. Set  $r^2 = x^2 + y^2 + z^2 + t^2$ , r > 0, and  $\Omega = g^{-1}dg$  for  $g \in G$ , that is,

if 
$$g = \begin{pmatrix} x & -y & -z & -t \\ y & x & -t & z \\ z & t & x & -y \\ t & -z & y & x \end{pmatrix}$$
 then  $\Omega = \begin{pmatrix} \sigma_1 & -\sigma_2 & -\sigma_3 & -\sigma_4 \\ \sigma_2 & \sigma_1 & -\sigma_4 & \sigma_3 \\ \sigma_3 & \sigma_4 & \sigma_1 & -\sigma_2 \\ \sigma_4 & -\sigma_3 & \sigma_2 & \sigma_1 \end{pmatrix}$ 

where

$$\begin{pmatrix} \sigma_1 \\ \sigma_2 \\ \sigma_3 \\ \sigma_4 \end{pmatrix} = \frac{1}{r^2} \begin{pmatrix} x & y & z & t \\ -y & x & t & -z \\ -z & -t & x & y \\ -t & z & -y & x \end{pmatrix} \begin{pmatrix} dx \\ dy \\ dz \\ dt \end{pmatrix}.$$

Then  $\sigma_j$ ,  $1 \leq j \leq 4$ , is a basis of left invariant 1-forms on G and it follows from  $d\Omega + \Omega \wedge \Omega = 0$  that

$$d\sigma_1 = 0$$
,  $d\sigma_2 = -2\sigma_3 \wedge \sigma_4$ ,  $d\sigma_3 = 2\sigma_2 \wedge \sigma_4$ ,  $d\sigma_4 = -2\sigma_2 \wedge \sigma_3$ 

Setting

$$e^1 = 2\sigma_1$$
,  $e^2 = 2\sigma_2$ ,  $e^3 = 2\sigma_3$ ,  $e^4 = 2\sigma_4$ 

so that  $\{e^j\}_{1\leq j\leq 4}$  satisfy (2.1), the left-invariant hyper-Hermitian metric is

(3.1) 
$$g = (e^1)^2 + (e^2)^2 + (e^3)^2 + (e^4)^2 = \frac{4}{r^2}(dx^2 + dy^2 + dz^2 + dt^2)$$

and since the Lee form is  $\theta = e^1 = d(2 \log r)$  the corresponding hyper-Kähler metric is  $\tilde{g} = e^{-2 \log r} g$ , that is,

(3.2) 
$$\tilde{g} = \frac{4}{r^2} \left( \frac{(dr)^2}{r^2} + (\sigma_2)^2 + (\sigma_3)^2 + (\sigma_4)^2 \right) = \frac{4}{r^4} (dx^2 + dy^2 + dz^2 + dt^2)$$

Case 2. Define a product on  $\mathbb{R}^4$  as follows:

$$(x, y, z, t)(x', y', z', t') = (x + e^{z}(x'\cos t - y'\sin t), y + e^{z}(x'\sin t + y'\cos t), z + z', t + t').$$

This defines a Lie group structure on  $\mathbb{R}^4$  that makes it isomorphic to the Lie group considered in  $\S 2$ , Case 2. The following 1-forms are left-invariant with respect to the above product:

(3.3) 
$$e^1 = e^{-z} \cos t dx + e^{-z} \sin t dy, \qquad e^3 = -dz,$$

(3.4) 
$$e^2 = -e^{-z} \sin t dx + e^{-z} \cos t dy, \qquad e^4 = -dt$$

These forms satisfy relations (2.2). The hyper-Hermitian metric is therefore given as follows:

$$g = (e^1)^2 + (e^2)^2 + (e^3)^2 + (e^4)^2 = e^{-2z}(dx^2 + dy^2) + dz^2 + dt^2$$

and the Lee form is  $\theta = 2e^3 = -2dz$ , so that the hyper-Kähler metric becomes

$$\tilde{g} = e^{2z}g = (dx^2 + dy^2) + e^{2z}(dz^2 + dt^2).$$

Observe that the change of coordinates  $s=e^z$  gives the following simple form for  $\tilde{q}$  on  $\mathbb{R}^+ \times \mathbb{R}^3$ :

$$\tilde{g} = dx^2 + dy^2 + ds^2 + s^2 dt^2.$$

Case 3. We endow  $\mathbb{R}^4$  with the following product:

$$(x, y, z, t)(x', y', z', t') = (x + e^t x', y + e^t y', z + e^t z', t + t')$$

thereby obtaining the Lie group structure considered in §2, Case 3, with corresponding left-invariant 1-forms:

$$e^{1} = dt$$
,  $e^{2} = e^{-t}dx$ ,  $e^{3} = e^{-t}dy$ ,  $e^{4} = e^{-t}dz$ .

The hyper-Hermitian metric is therefore

$$g = (e^{1})^{2} + (e^{2})^{2} + (e^{3})^{2} + (e^{4})^{2} = e^{-2t}(dx^{2} + dy^{2} + dz^{2}) + dt^{2}$$

with corresponding Lee form  $\theta = -2e^1 = -2dt$ , yielding the following hyper-Kähler metric:

$$\tilde{g} = e^{2t}g = dx^2 + dy^2 + dz^2 + e^{2t}dt^2.$$

Setting  $s = e^t$ ,  $\tilde{g}$  is the euclidean metric  $ds^2 + dx^2 + dy^2 + dz^2$  on  $\mathbb{R}^+ \times \mathbb{R}^3$  Case 4. Consider the following product on  $\mathbb{R}^4$ :

$$(x, y, z, t)(x', y', z', t') = (x + e^{\frac{t}{2}}x', y + e^{\frac{t}{2}}y', z + e^{t}z' + \frac{e^{\frac{t}{2}}}{4}(xy' - yx'), t + t')$$

which yields the Lie group structure considered in §2, Case 4. It is easily checked that the following left-invariant 1-forms satisfy (2.6):

$$e^{1} = dt$$
,  $e^{2} = e^{-t}(dz - \frac{1}{4}xdy + \frac{1}{4}ydx)$ ,  $e^{3} = e^{-\frac{t}{2}}dx$ ,  $e^{4} = e^{-\frac{t}{2}}dy$ .

The hyper-Hermitian metric is now obtained as in the above cases:

$$\begin{split} g &= (e^1)^2 + (e^2)^2 + (e^3)^2 + (e^4)^2 \\ &= dt^2 + e^{-t}(dx^2 + dy^2) + e^{-2t}(dz - \frac{1}{4}(xdy - ydx))^2 \end{split}$$

and the Lee form is  $\theta = -\frac{3}{2}dt$ , from which we obtain the hyper-Kähler metric as usual:

$$\tilde{g} = e^{-\frac{3}{2}t}dt^2 + e^{-\frac{t}{2}}(dx^2 + dy^2) + e^{-\frac{t}{2}}(dz - \frac{1}{4}(xdy - ydx))^2.$$

Setting  $s = e^{\frac{t}{2}}$ ,  $\tilde{g}$  becomes

$$\tilde{g} = s(ds^2 + dx^2 + dy^2) + \frac{1}{s}(dz - \frac{1}{4}(xdy - ydx))^2$$

on  $\mathbb{R}^+ \times \mathbb{R}^3$ , which allows us to identify  $\tilde{g}$  with one of the hyper-Kähler metrics constructed by the Gibbons-Hawking ansatz [2]. The identification is easily obtained from [4], Proposition 1.

We can now rephrase Theorem 1.1 as follows, where [h] denotes the conformal class of h:

Corollary 3.1. Let h be a hyper-Kähler metric on a simply connected hypercomplex 4-manifold  $(M,\mathcal{H})$  such that there exist  $g \in [h]$  and a Lie group  $G \subset I(M,g) \cap$  $Aut(\mathcal{H})$  acting simply transitively on M. Then (M,h) is homothetic to either  $\mathbb{R}^4$ with the euclidean metric or one of the following riemannian manifolds:

- 1.  $M = \mathbb{R}^4 \{0\}$ ,  $h = r^{-4}(dx^2 + dy^2 + dz^2 + dt^2)$ , 2.  $M = \mathbb{R}^+ \times \mathbb{R}^3$ ,  $h = ds^2 + dx^2 + dy^2 + s^2 dt^2$ , 3.  $M = \mathbb{R}^+ \times \mathbb{R}^3$ ,  $h = ds^2 + dx^2 + dy^2 + dz^2$ , 4.  $M = \mathbb{R}^+ \times \mathbb{R}^3$ ,  $h = s(ds^2 + dx^2 + dy^2) + s^{-1}(dz \frac{1}{4}(xdy ydx))^2$ .

#### References

- 1. M. L. Barberis, Hypercomplex structures on 4-dimensional Lie groups, Proc. Amer. Math. Soc. 125(4) (1997), 1043-1054.
- 2. G. W. Gibbons, S. W. Hawking, Gravitational multi-instantons, Phys. Lett. B 78 (1978), 430 - 432.
- 3. G. Grantcharov, Y. S. Poon, Geometry of hyper-Kähler connections with torsion, preprint, available in http://xxx.lanl.gov/, math.DG 9908015
- 4. C. LeBrun, Explicit self-dual metrics on @P 2# ... #@P 2, J. Differential Geometry 34(1) (1991), 223-253.

FAMAF, UNIVERSIDAD NACIONAL DE CÓRDOBA, CIUDAD UNIVERSITARIA, 5000 - CÓRDOBA, ARGENTINA

E-mail address: barberis@mate.uncor.edu