

# Resonance based schemes for dispersive equations via decorated trees

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"Higher Structures Emerging from Renormalisation",  
ESI Vienna, 15 October 2020

# Dispersive PDEs

We consider nonlinear dispersive equations of the form

$$\begin{aligned}i\partial_t u(t, x) + \mathcal{L}u(t, x) &= p(u(t, x), \bar{u}(t, x)) \\ u(0, x) &= v(x), \quad (t, x) \in \mathbb{R}_+ \times \mathbf{T}^d\end{aligned}$$

where  $\mathcal{L}$  is a differential operator and  $p$  is a polynomial nonlinearity.

- Assume local wellposedness of the problem on the finite time interval  $]0, T]$ ,  $T < \infty$  for  $v \in H^n$ .
- Aim: give a numerical approximation of  $u$  at low regularity when  $n$  is small.

NLS:  $\mathcal{L} = \Delta$  and  $p(u, \bar{u}) = |u|^2 u$ .

KdV:  $\mathcal{L} = i\partial_x^3$  and  $p(u, \bar{u}) = i\partial_x(u^2)$ .

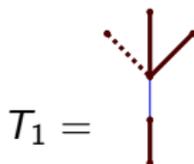
# Decorated trees approach

Mild solution given by Duhamel's formula:

$$u(t) = \underbrace{e^{it\mathcal{L}}}_{\text{edge } |} v + \underbrace{e^{it\mathcal{L}}}_{\text{edge } |} \underbrace{\left(-i \int_0^t e^{-i\xi\mathcal{L}} p(u(\xi), \bar{u}(\xi)) d\xi\right)}_{\text{edge } |}$$

Definition of a character  $\Pi : \text{Decorated trees} \rightarrow \text{Iterated integrals}$

- ①  $e^{it\mathcal{L}} v = (\Pi T_0)(t, v), \quad T_0 = |$
- ②  $-ie^{it\mathcal{L}} \int_0^t e^{-i\xi\mathcal{L}} p(e^{i\xi\mathcal{L}} v, e^{-i\xi\mathcal{L}} \bar{v}) d\xi = (\Pi T_1)(t, v),$



## B-series type expansion

Solution  $U^r$  up to order  $r$  can be represented by a series:

$$U^r(t, v) = \sum_{T \in \mathcal{V}^r} \frac{\Upsilon^p(T)}{S(T)} (\Pi T)(t, v),$$

- $\mathcal{V}^r$ : decorated trees of order  $r$ .
- $S(T)$ : symmetry factor.
- $\Upsilon^p$ : elementary differentials.
- Error of order

$$\mathcal{O}\left(t^{r+1}q(v)\right)$$

for some polynomial  $q$ .

## Treatment of oscillations

The principal oscillatory integral takes the form

$$\mathcal{I}_1(t, \mathcal{L}, \nu, \rho) = \int_0^t \text{Osc}(\xi, \mathcal{L}, \nu, \rho) d\xi$$

with the central oscillations  $\text{Osc}$  given by

$$\text{Osc}(\xi, \mathcal{L}, \nu, \rho) = e^{-i\xi\mathcal{L}} \rho \left( e^{i\xi\mathcal{L}} \nu, e^{-i\xi\mathcal{L}} \bar{\nu} \right).$$

In general it will be

$$\text{Osc}(\xi, \mathcal{L}, \nu, \rho) = e^{-i\xi\mathcal{L}} \rho \left( e^{i\xi(\mathcal{L}+\mathcal{L}_1)} q_1(\nu), e^{-i\xi(\mathcal{L}+\mathcal{L}_2)} q_2(\nu) \right).$$

# Various approaches

Classical Methods:

- exponential method:  $\text{Osc}(\xi, \mathcal{L}, v, p) \approx e^{-i\xi\mathcal{L}}p(v, \bar{v})$
- splitting method:  $\text{Osc}(\xi, \mathcal{L}, v, p) \approx p(v, \bar{v})$

Resonance as a computational tool:

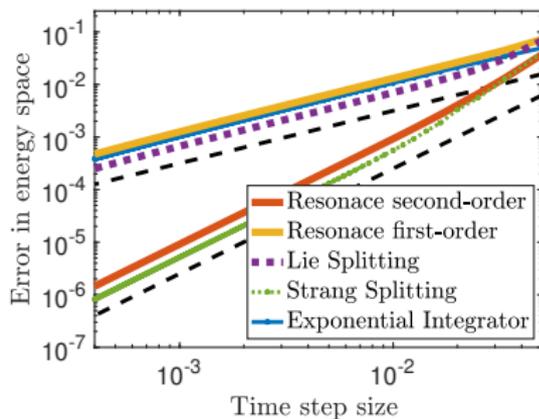
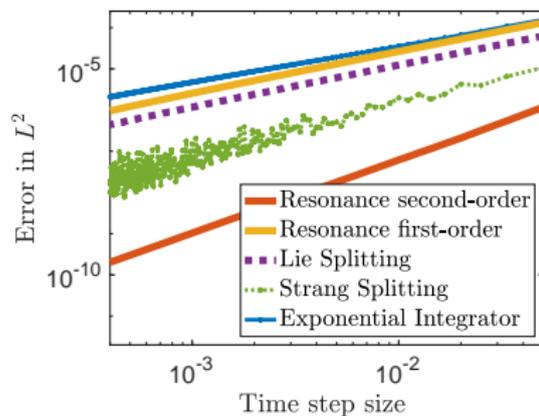
$$\text{Osc}(\xi, \mathcal{L}, v, p) = \left[ e^{i\xi\mathcal{L}_{\text{dom}}} p_{\text{dom}}(v, \bar{v}) \right] p_{\text{low}}(v, \bar{v}) + \mathcal{O}(\xi\mathcal{L}_{\text{low}}v).$$

Here,  $\mathcal{L}_{\text{dom}}$  denotes a suitable dominant part of the high frequency interactions and

$$\mathcal{L}_{\text{low}} = \mathcal{L} - \mathcal{L}_{\text{dom}}$$

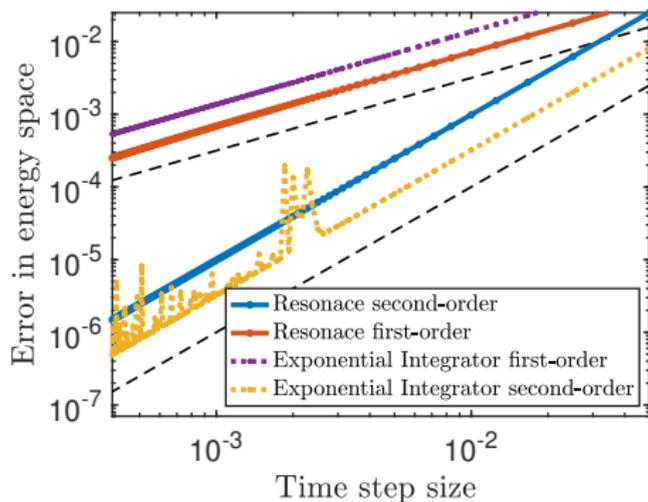
# Experiments

Comparison of classical and resonance based schemes for the Schrödinger equation for smooth ( $C^\infty$  data) and non-smooth ( $H^2$  data) solutions.



# Experiments

Comparison of classical and resonance based schemes for the KdV equation with smooth data in  $C^\infty$ .



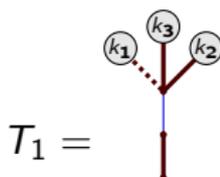
# Fourier iterated integrals

Mild solution given by Duhamel's formula in Fourier ( $P(k) \leftrightarrow \mathcal{L}$ ):

$$\hat{u}_k(t) = \underbrace{e^{itP(k)}}_{\text{edge } |} \hat{v}_k + \underbrace{e^{itP(k)}}_{\text{edge } |} \underbrace{\left(-i \int_0^t e^{-i\xi P(k)} p_k(u(\xi), \bar{u}(\xi)) d\xi\right)}_{\text{edge } |}$$

Definition of a character  $\hat{\Pi} : \text{Decorated trees} \rightarrow \text{Iterated integrals}$

- 1  $e^{itP(k)} = (\hat{\Pi} T_0)(t), \quad T_0 = \text{⌞}^{\textcircled{k}}$
- 2  $-ie^{itP(k)} \int_0^t e^{-i\xi(P(k)-P(-k_1)+P(k_2)+P(k_3))} d\xi = (\hat{\Pi} T_1)(t),$



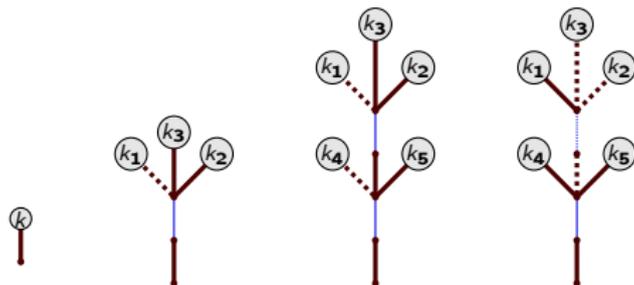
# Fourier B-series type expansion

Resonance scheme  $U_k^r$  of order  $r$  with regularity  $n$  (initial data):

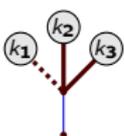
$$U_k^r(\tau, \nu) = \sum_{T \in \mathcal{V}_k^r} \frac{\Upsilon^P(T)(\nu)}{S(T)} \left( \hat{\Pi}_n^r T \right) (\tau)$$

- $\mathcal{V}_k^r$ : decorated trees of order  $r$  with frequency  $k$ .
- Character  $\hat{\Pi}_n^r$  resonance approximation of  $\hat{\Pi}$ .

Examples of decorated trees for NLS ( $r = 2$ ):



## A practical example

The iterated integral associated to  $T =$   is given by:

$$(\hat{\Pi} T)(t) = \int_0^t e^{i\xi(-k^2 - k_1^2 + k_2^2 + k_3^2)} d\xi, \quad k = -k_1 + k_2 + k_3$$

One has

$$-k^2 - k_1^2 + k_2^2 + k_3^2 = \underbrace{-2k_1^2}_{\mathcal{L}_{\text{dom}}} + \underbrace{\quad}_{\mathcal{L}_{\text{low}} \text{ order one}}$$

Taylor expansion of  $\mathcal{L}_{\text{low}}$ :

$$(\hat{\Pi} T)(t) = \underbrace{\frac{e^{-2itk_1^2} - 1}{-2ik_1^2}}_{(\hat{\Pi}_n^r T)(t)} + \mathcal{O}(t\mathcal{L}_{\text{low}})$$

# Local Error Analysis

Main idea is to single out oscillations:

$$\int_0^t e^{i\xi P(k)} d\xi = \frac{e^{itP(k)} - 1}{iP(k)}$$

Butcher-Connes-Kreimer coproduct  $\Delta$

$$\Delta = \text{tree}(k_1, k_2, k_3) \otimes \text{tree}(k_4, k_5, l, k_3, k_2) + \dots, \quad l = -k_1 + k_2 + k_3$$

Integrals  $\int_0^t \xi^\ell e^{i\xi P(k)} d\xi \rightarrow$  deformed BCK coproduct  $\hat{\Delta}$  (SPDEs).

# Main result

Birkhoff type factoriation:

$$\tilde{\Pi}_n^r = \left( \hat{\Pi}_n^r \otimes (\mathcal{Q} \circ \hat{\Pi}_n^r \mathcal{A} \cdot)(0) \right) \hat{\Delta}.$$

where  $\mathcal{A}$  is an antipode and  $\mathcal{Q}$  is a projector.

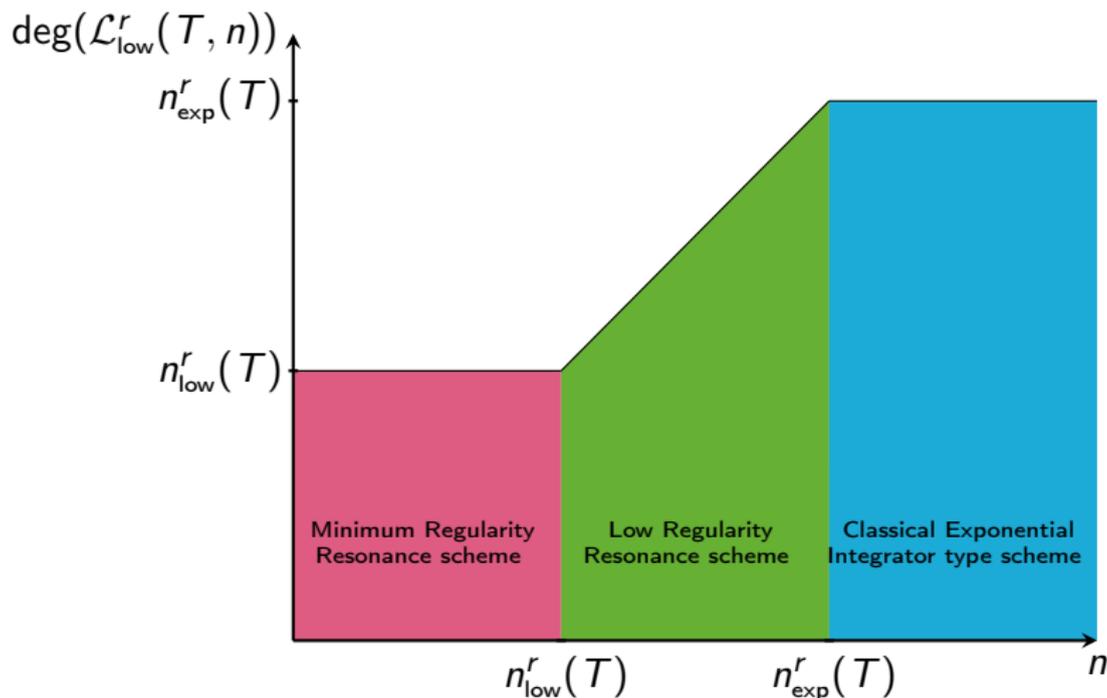
Theorem (B., Schratz 2020)

For every  $T \in \mathcal{V}_k^r$

$$\left( \hat{\Pi} T - \hat{\Pi}_n^r T \right) (\tau) = \mathcal{O} \left( \tau^{r+1} \mathcal{L}_{low}^r(T, n) \right).$$

where  $\mathcal{L}_{low}^r(T, n)$  involves all lower order frequency interactions.

# Family of schemes $\hat{\Pi}_n^r$



- B-series  $\rightarrow$  Regularity Structures  $\rightarrow$  PDEs Numerical Schemes
- New example of a deformation of the BCK coproduct.
- Birkhoff type factorisation as in SPDEs see (B., Ebrahimi-Fard 2020).
- Backward error analysis for the scheme.
- Structure preservation.
- Generalisation to more general domains not only  $\mathbf{T}^d$  and wave equations.
- Potential connection with the study of dispersive (S)PDEs.