Blobbed topological recursion of the quartic analogue of the Kontsevich model

Raimar Wulkenhaar

Mathematisches Institut der Westfälischen Wilhelms-Universität Münster



arXiv:2008.12201 with Johannes Branahl & Alex Hock building on previous collaborations with Harald Grosse, Erik Panzer & Jörg Schürmann

Raimar Wulkenhaar (Münster)

Blobbed topological recursion of the quartic analogue of the Kontsevich model

Planar 2-point function	Residue formulae for Dyson-Schwing

Discussion

er equations

Introduction

Introduction

00000

This project started in 1998 as an attempt to understand quantum field theories on noncommutative geometries.

- Many nice results were obtained on renormalisation, β-function, extension to tensor models.
- We aimed at an analytic construction of these models (of course combined with combinatorics).

During the last two years we understood: the constructive path is misleading!

- It is about miracles in algebra, about amazing solutions of non-linear problems.
- It is another example for topological recursion, thus links to structures in complex algebraic geometry.

We just entered a new world and admire the structures seen so far. But much more remains to explore.

Interlude: Topological recursion [Eynard-Orantin 07]

Universal structure that governs, e.g.

one- and two-matrix models, Kontsevich model, Weil-Petersson volumes, Hurwitz numbers, Gromov-Witten numbers, ...

Starting from a spectral curve consisting of

- a branched covering $x : \Sigma \to \Sigma_0$ of Riemann surfaces,
- meromorphic differentials $\omega_{0,1}$ on Σ and $\omega_{0,2}$ on $\Sigma \times \Sigma$,

recursively construct family $\omega_{g,n}$ of meromorphic differentials on Σ^n , with poles at zeros of dx (ramification points), by

$$\omega_{g,n}(z_1,...,z_n) = \sum_{a} \underset{q \to a}{\operatorname{Res}} K(z_1, q, \sigma_a(q)) dz \Big(\omega_{g-1,n+1}(q, \sigma_a(q), z_2, ..., z_n) \\ + \sum_{\substack{q_1+q_2=g\\l_1 \uplus l_2 = \{z_2,...,z_n\}}} ' \omega_{g_1,|l_1|+1}(q, l_1) \omega_{g_2,|l_2|+1}(\sigma_a(q), l_2) \Big)$$

sum over ramification points *a* of *x*; local involution $x(q) = x(\sigma_a(q))$ near *a*; recursion kernel $K(z_1, z_2, z_3) = \frac{\frac{1}{2} \int_{z'=z_3}^{z_2} \omega_{0,2}(z_1, z')}{\omega_{0,1}(z_2) - \omega_{0,1}(z_3)}$

Free Euclidean fields on noncommutative geometries

Let X_N be the real vector space of self-adjoint $N \times N$ -matrices, and (E_1, \ldots, E_N) be (increasing) positive real numbers.

Theorem [Bochner 1933, Schur 1911]

For any inner product \langle , \rangle on X_N there exists a unique probability measure $d\mu_0$ on the dual space X'_N with

$$\exp\left(-\frac{1}{2}\langle M,M\rangle\right) = \int_{X'_N} d\mu_0(\Phi) \ e^{i\Phi(M)} \quad \forall M = (M_{kl}) \in X .$$

Choose $\langle M,M'\rangle_E = \frac{1}{N} \sum_{k,l=1}^N \frac{M_{kl}M'_{lk}}{E_k + E_l}$ and corresponding $d\mu_{E,0}$

- Defines the free Euclidean scalar field on *N*-dimensional approximation of a noncommutative geometry.
- (E_1, \ldots, E_N) is truncated spectrum of the Laplacian.
- All moments can be described explicitly.

Two deformations

The Kontsevich model
$$d\mu_{E,\lambda}(\Phi) = \frac{e^{-\frac{\lambda N}{3} \operatorname{Tr}(\Phi^3)} d\mu_{E,0}(\Phi)}{\int_{X'_N} e^{-\frac{\lambda N}{3} \operatorname{Tr}(\Phi^3)} d\mu_{E,0}(\Phi)}$$

- Computes intersection numbers of tautological characteristic classes on the moduli space $\overline{\mathcal{M}}_{g,n}$ of stable complex curves.
- It is integrable via a relation (suggested by Witten) to the KdV hierarchy. Its moments satisfy topological recursion.

A quartic analogue
$$d\mu_{E,\lambda}(\Phi) = rac{e^{-rac{\lambda N}{4} \operatorname{Tr}(\Phi^4)} d\mu_{E,0}(\Phi)}{\int_{X'_N} e^{-rac{\lambda N}{4} \operatorname{Tr}(\Phi^4)} d\mu_{E,0}(\Phi)}$$

 Although perturbatively far apart, we find very similar algebraic geometrical structures. Our solutions are exact in λ.

ntroduction	Planar 2-point function	Residue formulae for Dyson-Schwinger equations	Discussion
Quarviau	,		

Aim: Say something about moments $\int d\mu_{E,\lambda}(\Phi) \Phi_{k_1l_1} \cdots \Phi_{k_nl_n}$

- As in all matrix models, best approach is 1/N-expansion.
- After 10 years of work, we found in [Grosse-Hock-W 19] the 1/N-leading contribution to the 2-point function as exact function of λ, E₁,..., E_N.
- This involved the solution of a non-linear problem (identified in [Grosse-W 09], main step is [Panzer-W 18]).

Exact solutions of non-linear problems are only possibly in presence of a profound algebraic structure

- By analogy with the Kontsevich model we expected that this structure consists in a topological recursion which governs the entirerty of moments of $d\mu_{E,\lambda}$.
- Recent discovery [Alex Hock]: The structure is present, but at an unexpected place!

The objects for blobbed topological recursion

Recall that $d\mu_{E,\lambda}$ depends on given family E_1, \ldots, E_N . Introduce

$$\sum_{g=0}^{\infty} N^{2-2g-n} \Omega_{a_1,\ldots,a_n}^{(g)} := \frac{\partial^{n-1} \left(N \sum_{k=1}^{N} \int d\mu_{E,\lambda}(\Phi) \Phi_{a_1k} \Phi_{ka_1} \right)}{\partial E_{a_1} \cdots \partial E_{a_{n-1}}} + \frac{\delta_{n,2}}{(E_{a_1} - E_{a_2})^2}$$

- As substitute for unavailable direct *E*-derivatives we derive Dyson-Schwinger equations for the Ω.
- They extend naturally (not uniquely) to (very complicated) equations for meromorphic functions in several complex variables.
- Their solutions are strikingly simple and structured: The Ω extend to meromorphic forms $\omega_{g,n}$ which exactly follow the rules of blobbed topological recursion [Borot-Shadrin 15].

Introduction

00000

Planar 2-point function

Residue formulae for Dyson-Schwinger equations

Discussion

Equations of motion

Fourier transform
$$\mathcal{Z}(M):=\int_{X_N'}d\mu_{E,\lambda}(\Phi)\;e^{\mathrm{i}\Phi(M)}$$
 satisfies

$$= -N(E_p - E_q) \sum_{k=1}^{N} \frac{\partial^2 \mathcal{Z}(M)}{\partial M_{pk} \partial M_{kq}} = \sum_{k=1}^{N} \left(M_{kp} \frac{\partial \mathcal{Z}(M)}{\partial M_{kq}} - M_{qk} \frac{\partial \mathcal{Z}(M)}{\partial M_{pk}} \right)$$

$$= \frac{1}{N} \frac{\partial \mathcal{Z}(M)}{\partial E_p} = \sum_{k=1}^{N} \frac{\partial^2 \mathcal{Z}(M)}{\partial M_{pk} \partial M_{kp}} + \mathcal{Z}(M) \int d\mu_{E,\lambda}(\Phi) \frac{1}{N} \sum_{k=1}^{N} \Phi_{pk} \Phi_{kp}$$

They allow to express $\sum_{k=1}^{N} \frac{\mathcal{Z}(M)}{\partial M_{pk} \partial M_{kq}}$ in Dyson-Schwinger equations by fewer derivatives, i.e. of same or lower order.

Eq. (essentially due [Disertori-Gurau-Magnen-Rivasseau 06]) can be used for $p \neq q$, whereas p = q requires (2).

Introduction
000000

Planar 2-point function

Residue formulae for Dyson-Schwinger equations

Discussion

Meromorphic continuation

These Dyson-Schwinger equations complexify to equations for meromorphic functions in several complex variables in which we can admit multiplicities $(E_1, \ldots, E_N) = (\underbrace{e_1, \ldots, e_1}_{r_1}, \ldots, \underbrace{e_d, \ldots, e_d}_{r_d})$

Example: The two-point function

For
$$p \neq q$$
, set $\sum_{g=0}^{\infty} N^{-2g} G_{|pq|}^{(g)} = N \int d\mu_{E,\lambda}(\Phi) \Phi_{pq} \Phi_{qp}$.
Then $G_{|pq|}^{(g)} = G^{(g)}(e_p, e_q)$ with initial equation [Grosse-W 09]
 $(\xi + \eta)G^{(0)}(\zeta, \eta)$
 $= 1 - \lambda G^{(0)}(\zeta, \eta)\Omega_1^{(0)}(\zeta) + \frac{1}{N} \sum_{k=1}^d r_k \frac{G^{(0)}(e_k, \eta) - G^{(0)}(\zeta, \eta)}{e_k - \zeta}$,
 $\Omega_1^{(0)}(\zeta) = \frac{1}{N} \sum_{k=1}^d r_k G^{(0)}(\zeta, e_k)$

Planar 2-point function

Residue formulae for Dyson-Schwinger equations

Discussion

Main achievement: We can guess $\Omega_1^{(0)}$!

$$\left(\eta+\zeta+\frac{\lambda}{N}\sum_{k=1}^{d}r_{k}G^{(0)}(\zeta,\boldsymbol{e}_{k})+\frac{\lambda}{N}\sum_{k=1}^{d}\frac{r_{k}}{\boldsymbol{e}_{k}-\zeta}\right)G^{(0)}(\zeta,\eta)=1+\frac{\lambda}{N}\sum_{k=1}^{d}r_{k}\frac{G^{(0)}(\boldsymbol{e}_{k},\eta)}{\boldsymbol{e}_{k}-\zeta}$$

Ansatz [Schürmann-W 19, building on Grosse-Hock-W 19]

Assume there is a branched cover $\boldsymbol{R}: \hat{\mathbb{C}} \to \hat{\mathbb{C}}$ with

R has degree d + 1 and maps (for λ > 0) some domain U bijectively to neighbourhood of e₁,..., e_d

$$(2) \quad \zeta = R(z), \ \eta = R(w), \ e_k = R(\varepsilon_k), \ G^{(0)}(\zeta, \eta) = \mathcal{G}^{(0)}(z, w)$$

$$R(z) + \frac{\lambda}{N} \sum_{k=1}^{d} r_k \mathcal{G}^{(0)}(z, \varepsilon_k) + \frac{\lambda}{N} \sum_{k=1}^{d} \frac{r_k}{R(\varepsilon_k) - R(z)} = -R(-z)$$

Gives
$$(R(w) - R(-z))\mathcal{G}^{(0)}(z, w) = 1 + \frac{\lambda}{N} \sum_{k=1}^{a} r_k \frac{\mathcal{G}^{(0)}(\varepsilon_k, w)}{R(\varepsilon_k) - R(z)}$$

Introduction

Rationality

With techniques for Lagrange interpolation polynomials [Schechter 59] it is *easy* to establish:

Theorem [Schürmann-W 19]

•
$$R(z) = z - \frac{\lambda}{N} \sum_{k=1}^{d} \frac{\varrho_k}{z + \varepsilon_k}$$
 where $e_k = R(\varepsilon_k)$, $r_k = R'(\varepsilon_k)\varrho_k$
• $\mathcal{G}^{(0)}(z, w) = \frac{1 - \frac{\lambda}{N} \sum_{k=1}^{d} \frac{r_k \prod_{j=1}^{d} \frac{R(w) - R(-\varepsilon_k^j)}{R(w) - R(\varepsilon_j)}}{(R(z) - R(\varepsilon_k))(R(\varepsilon_k) - R(-w))}}{R(w) - R(-z)}$
where $u \in \{z, \hat{z}^1, \dots, \hat{z}^d\}$ are all solutions of $R(u) = R(z)$.
• $\mathcal{G}^{(0)}(z, w) = \mathcal{G}^{(0)}(w, z)$

The ansatz ③ is identically satisfied!

Thus, planar 2-point function solved by the composition of a rational function $\mathcal{G}^{(0)}$ with inverse of another rational function *R*.

Raimar Wulkenhaar (Münster)

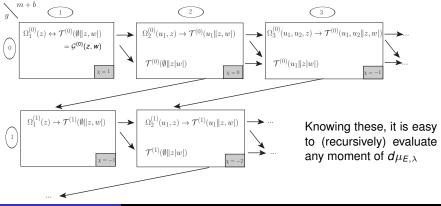
Blobbed topological recursion of the quartic analogue of the Kontsevich model

Introduction

Solution procedure

Three types of functions involved:

- $\Omega_m^{(g)}(u_1, ..., u_m)$ objects of BTR, most difficult to compute
- *T*^(g)(*u*₁,..., *u_m*||*z*, *w*|) auxiliary functions, easy to compute
 T^(g)(*u*₁,..., *u_m*||*z*|*w*|) auxiliary functions, easy to compute



 $\Omega_{2}^{(0)}(u,z)$

$$\Omega_{2}^{(0)}(u,z)R'(z)\mathfrak{G}_{0}(z) - \frac{\lambda}{N^{2}}\sum_{n,k=1}^{d}\frac{r_{k}r_{n}\mathcal{T}^{(0)}(u||\varepsilon_{k},\varepsilon_{n}|)}{(R(\varepsilon_{k}) - R(z))(R(\varepsilon_{n}) - R(-z))}$$
$$= -\frac{\partial}{\partial R(u)}\left(\mathcal{G}^{(0)}(u,z) + \mathcal{G}^{(0)}(u,-z)\right)$$
where $\mathfrak{G}_{0}(z) = \operatorname{Res}_{W \to -z} \mathcal{G}^{(0)}(z,W).$

- Seems to need $\mathcal{T}^{(0)}(u \| \varepsilon_k, \varepsilon_n|)$ which itself needs $\Omega_2^{(0)}$.
 - But poles separate by partial fraction decomposition $\mathcal{G}^{(0)}(z,u) = \frac{\mathfrak{G}_0(z)}{u+z} + \frac{\lambda^2}{N^2} \sum_{k,l,m,n=1}^d \frac{C_{k,l}^{m,n}}{(z+\widehat{\varepsilon_l}^n)(z-\widehat{\varepsilon_k}^m)(u-\widehat{\varepsilon_l}^n)}$

Proposition

$$\Omega_2^{(0)}(u,z) = \frac{1}{R'(u)R'(z)} \Big(\frac{1}{(u-z)^2} + \frac{1}{(u+z)^2}\Big)$$

One recognises the Bergman kernel of topological recursion!

Blobbed topological recursion of the quartic analogue of the Kontsevich model

Planar 2-point function

Residue formulae for Dyson-Schwinger equations

Discussion

13

Proposition: $\mathcal{T}^{(g)}(u_1, \ldots, u_m || z, w|)$

The solution is easier than the equation itself:

$$\begin{aligned} \mathcal{T}^{(g)}(u_{1},...,u_{m}||z,w|) \\ &= \lambda \mathcal{G}^{(0)}(z,w) \underset{t \to z, -\hat{w}^{j}}{\operatorname{Res}} \frac{R'(t) dt}{(R(z) - R(t))(R(w) - R(-t))\mathcal{G}^{(0)}(t,w)} \\ &\times \left[\sum_{\substack{l_{1} \uplus l_{2} = \{u_{1},...,u_{m}\}\\g_{1} + g_{2} = g\\(l_{1},g_{1}) \neq (\emptyset,0)} \Omega_{ll_{1}|+1}^{(g_{1})}(l_{1},t)\mathcal{T}^{(g_{2})}(l_{2}||t,w|) \right. \\ &+ \sum_{i=1}^{m} \frac{\partial}{\partial R(u_{i})} \frac{\mathcal{T}^{(g)}(u_{1},...\check{u}_{i}...u_{m}||u_{i},w|)}{R(u_{i}) - R(t)} \\ &+ \frac{\mathcal{T}^{(g-1)}(u_{1},...,u_{m}||t|w|) - \mathcal{T}^{(g-1)}(u_{1},...,u_{m}||w|w|)}{R(w) - R(t)} \\ &+ \mathcal{T}^{(g-1)}(u_{1},...,u_{m},t||t,w|) \right] \end{aligned}$$

Planar 2-point function

Residue formulae for Dyson-Schwinger equations

Discussion 00

Proposition: $\overline{\mathcal{T}^{(g)}(u_1,\ldots,u_m||z|w|)}$

$$\begin{split} \mathcal{T}^{(g)}(u_{1},...,u_{m}||z|w|) \\ &= \frac{\lambda \prod_{j=1}^{d} \frac{R(z) - R(\alpha_{j})}{R(z) - R(-z))}}{(R(z) - R(-z))} \underset{t \to z, \alpha_{j}}{\operatorname{Res}} \frac{R'(t) dt}{(R(z) - R(t)) \prod_{j=1}^{d} \frac{R(t) - R(\alpha_{j})}{R(t) - R(\varepsilon_{j})}} \\ &\times \bigg[\sum_{\substack{l_{1} \uplus l_{2} = \{u_{1},...,u_{m}\}\\g_{1} + g_{2} = g\\(l_{1},g_{1}) \neq (\emptyset, 0)}} \Omega_{(l_{1}|+1}^{(g_{1})}(l_{1},t) \mathcal{T}^{(g_{2})}(l_{2}||t|w|) + \mathcal{T}^{(g-1)}(u_{1},...,u_{m},t||t|w|)} \\ &+ \sum_{i=1}^{m} \frac{\partial}{\partial R(u_{i})} \frac{\mathcal{T}^{(g)}(u_{1},...\check{u}_{i}...,u_{m} \setminus u_{i}||u_{i}|w|)}{R(u_{i}) - R(t)} \\ &+ \frac{\mathcal{T}^{(g)}(u_{1},...,u_{m}||t,w|) - \mathcal{T}^{(g)}(u_{1},...,u_{m}||w,w|)}{R(w) - R(t)} \bigg], \end{split}$$

where $\{0, \pm \alpha_1, .. \pm \alpha_d\}$ are all solutions of R(z) - R(-z) = 0.

Planar 2-point function

Residue formulae for Dyson-Schwinger equations

Discussion 00

$\Omega^{(g)}(u_1,\ldots,u_m)$ for 2-2g-m<0

Proposition (g = 0) / Conjecture ($g \ge 1$)

$$\begin{aligned} R'(z)\Omega_{m+1}^{(g)}(u_{1},...,u_{m},z) &= \underset{q \to 0,-u_{l},\beta_{l}}{\operatorname{Res}} \frac{dq}{(q-z)\mathfrak{G}_{0}(q)} \left[\\ &\sum_{\substack{l_{1} \uplus l_{2} = \{u_{1},...,u_{m}\}\\g_{1}+g_{2}=g}(l_{1},...,u_{m})}{\sum_{q=1}^{g+g_{2}=g}(l_{1},g_{1})\neq(\emptyset,0)\neq(l_{2},g_{2})} &+ \sum_{i=1}^{m} \frac{\partial}{\partial R(u_{i})}\mathcal{T}^{(g)}(u_{1},...\check{u}_{i}...,u_{m}||u_{i},q|) \\ &+ \frac{\lambda}{N}\sum_{n=1}^{d} \frac{r_{n}\mathcal{T}^{(g-1)}(u_{1},...,u_{m},q||q,\varepsilon_{n}|)}{R(\varepsilon_{n})-R(-q)} \\ &+ \frac{\lambda}{N}\sum_{n=1}^{d} \frac{r_{n}\mathcal{T}^{(g-1)}(u_{1},...,u_{m}||q||g_{n}|)}{R(\varepsilon_{n})-R(-q)} - \mathcal{T}^{(g-1)}(u_{1},...,u_{m}||q||q|) \right] \end{aligned}$$

 $\beta_1, \ldots, \beta_{2d}$ – ramification points, $R'(\beta_i) = 0$

Introduction Planar 2-point function 0000 Pla

Differential forms $\omega_{g,m}(z_1,...,z_m) = \lambda^{2-2g-m} \Omega_m^{(g)}(z_1,...,z_m) \prod^m dR(z_i)$

There is not much to improve for the $\mathcal{T}^{(g)}$. In contrast, the $\Omega_m^{(g)}$ simplify tremendously (by a factor > 100) compared with the formula. Their structure is only visible after these simplifications.

$$\begin{split} \omega_{0,3}(u_1, u_2, z) &= -\sum_{i=1}^{2d} \frac{\left(\frac{1}{(u_1 - \beta_i)^2} + \frac{1}{(u_1 + \beta_i)^2}\right) \left(\frac{1}{(u_2 - \beta_i)^2} + \frac{1}{(u_2 + \beta_i)^2}\right) du_1 \, du_2 \, dz}{R'(-\beta_i) R''(\beta_i) (z - \beta_i)^2} \\ &+ \left[d_{R(u_1)} \left(\frac{\omega_{0,2}(u_2, -u_1)}{(dR)(-u_1)} \frac{dz}{R'(u_1)(z + u_1)^2}\right) + u_1 \leftrightarrow u_2 \right] \\ \omega_{1,1}(z) &= \sum_{i=1}^{2d} \frac{dz}{R'(-\beta_i) R''(\beta_i)} \left\{ -\frac{1}{8(z - \beta_i)^4} + \frac{R'''(\beta_i)}{24R''(\beta_i)(z - \beta_i)^3} \right. \\ &+ \frac{\frac{R'''(\beta_i)}{48R''(\beta_i)} - \frac{(R'''(\beta_i))^2}{48(R''(\beta_i))^2} + \frac{R''(-\beta_i)R'''(\beta_i)}{48R'(-\beta_i)R''(\beta_i)} + \frac{(R''(-\beta_i))^2}{48(R'(-\beta_i))^2} - \frac{1}{8\beta_i^2}}{(z - \beta_i)^2} \right\} \\ &- \frac{dz}{8(R'(0))^2 z^3} + \frac{R''(0) dz}{16(R'(0))^3 z^2} \end{split}$$

Raimar Wulkenhaar (Münster)

16

Abstract loop equations [Borot-Eynard-Orantin 13]

Proposition $(g, m) \in \{(0, 2), (0, 3), (0, 4), (1, 1)\}$ / Conjecture

Let $\mathbf{R}: \hat{\mathbb{C}} \to \hat{\mathbb{C}}$ be the ramified cover identified in the solution of $\mathcal{G}^{(0)}(z, w)$.

Let $\beta_1, ..., \beta_{2d}$ be the ramification points of *R* and σ_i be the corresponding local Galois involution in the vicinity of β_i .



2

Define $\omega_{0,1}(z) = -R(-z)R'(z)dz$ and for $2 - 2g - m \le 0$ the $\omega_{g,m}$ as before. Then:

 linear loop equation: ω_{g,m}(u₁,..., u_{m-1}, z) + ω_{g,m}(u₁,..., u_{m-1}, σ_i(z)) = O(z-β_i)dz

 quadratic loop equation:

$$\begin{split} & \omega_{g-1,m+1}(u_1,...,u_{m-1},z,\sigma_i(z)) \\ & + \sum_{\substack{l_1 \uplus l_2 = \{u_1,...,u_{m-1}\}\\g_1+g_2=g}} \omega_{g_1,|l_1|+1}(l_1,z)\omega_{g_2,|l_2|+1}(l_2,\sigma_i(z)) \\ & = \mathcal{O}((z-\beta_i)^2)(dz) \end{split}$$

Blobbed topological recursion [Borot-Shadrin 15]

Theorem

Let $\{\omega_{g,m}\}_{a>0,m>0}$ be a family of meromorphic differential forms which satisfy the abstract loop equations. Then their parts $\mathcal{P}\omega_{q,m}$ containing the poles at ramification points are given by $\mathcal{P}_{z}\omega_{q,m}(u_1,...,u_{m-1},z)$ $=\sum_{i=1}^{2d} \operatorname{Res}_{q \to \beta_{i}} \frac{\frac{1}{2} \int_{q'=\sigma(q)}^{q'=q} B(z,q')}{\omega_{0,1}(q) - \omega_{0,1}(\sigma_{i}(q))} \left(\omega_{g-1,m+1}(u_{1},...,u_{m-1},q,\sigma_{i}(q)) + \sum_{\substack{l_{1} \uplus l_{2} = \{u_{1},...,u_{m-1}\}}} \omega_{g_{1},|l_{1}|+1}(l_{1},q) \omega_{g_{2},|l_{2}|+1}(l_{2},\sigma_{i}(q)) \right)$ $g_1 + g_2 = g$ $(l_1, q_1) \neq (\emptyset, \bar{0}) \neq (l_2, q_2)$

where $B(u, z) = \frac{du dz}{(u-z)^2}$ is the Bergman kernel (for $x : \hat{\mathbb{C}} \to \hat{\mathbb{C}}$).

$$\mathcal{H}_{z}\omega_{g,m}(...,z) := \omega_{g,m}(...,z) - \mathcal{P}_{z}\omega_{g,m}(...,z)$$
 is made of blobs.



K

• The universal formula of topological recursion produces the parts $\mathcal{P}\omega_{g,m}$ from the entire $\omega_{g',m'}$ of smaller degree.

 $\omega_{g,m}$ = meromorphic forms on space of compactified complex lines through the marked points on a genus-*g* Riemann surface.

 $\mathcal{H}_{z_1}\omega_{q,m}(z_1,\ldots,z_m)$

 The part Hω_{g,m} is an additional input at every recursion step. We have a residue formula for them.

The quartic analogue of the Kontsevich model distinguishes a unique such form $\omega_{g,m}$ for every (g, m). What is its significance?

 $\omega_{q,m}(z_1, ..., z_m)$

K

 $\mathcal{P}_{z_1}\omega_{q,m}(z_1,\ldots,z_m)$

 $\sigma(c$

 $z_j, j \in I \setminus J$

19

Intersection numbers and integrability

Fact [Borot-Shadrin 15]

- Forms ω_{g,m} which satisfy BTR encode intersection numbers on the moduli space M_{g,m} of stable complex curves.
- These are several copies of the same intersections of ψ, κ -classes as in the Kontsevich model, coupled via blobs.

We do not know whether this coupling function encoded in the quartic Kontsevich model is interesting or not.

Integrability

- Inderstanding better our recursion should give access to the partition function itself, a function of λ and (*E_i*).
- Is it a *τ*-function for a Hirota equation, i.e. is it integrable? [not known in general BTR]

21

Outlook: $\lambda \Phi^4$ on noncommutative Moyal space

- **2** Moyal in 2*D* and $\theta \to \infty$ described by a Lambert curve $x(z) = R(z) = z + \lambda \log(1+z), y(z) = -R(-z)$ [Panzer-W 18]
 - It has a single ramification point of infinite order.
 - According to the [Bouchard-Mariño 07] conjecture, *another* Lambert curve generates simple Hurwitz numbers.
 - Is this more than a formal analogy? See [Borot-Eynard-Mulase-Safnuk 09]
- Moyal in 4D and $\theta \to \infty$ according to [Grosse-Hock-W 20] described by $(arcsin(2\pi)) = arcsin(2\pi)$

$$R(z) = z \cdot {}_{2}F_{1} \left({}^{\alpha_{\lambda}, 1-\alpha_{\lambda}}_{2} \middle| -z \right), \ \alpha_{\lambda} = \begin{cases} \frac{\arcsin(\lambda \pi)}{\pi} & \text{for } |\lambda| \leq \frac{1}{\pi} \\ \frac{1}{2} + i\frac{\operatorname{arcsn}(\lambda \pi)}{\pi} & \text{for } \lambda \geq \frac{1}{\pi} \end{cases}$$

- Ramification points accumulate to a curve. Can one turn the TR-formula into an integral representation?
- Can one generalise this to finite θ ?

Raimar Wulkenhaar (Münster)

Introduction